Symmetry in Particle Physics, Problem Sheet 1, Solutions

- 1. Consider a group G. Then show the following:
 - (a) for any $a, b, c \in G$, if ab = ac, then b = c (cancellation rule);

$$b = 1 \cdot b = a^{-1}(ab) = a^{-1}(ac) = c$$
.

(b) the unit element is unique; Suppose there exists another unit e'. Then, for any $g \in G$, we have

$$g = eg = e'g \implies e = e'$$
.

(c) the inverse of any group element is unique; Consider an element g, with two inverses g^{-1} and $(g^{-1})'$. Then

$$e = gg^{-1} = g(g^{-1})' \implies g^{-1} = (g^{-1})'.$$

(d)
$$(g^{-1})^{-1} = g$$
.
 $e = (g^{-1})^{-1}g^{-1} = gg^{-1} \implies (g^{-1})^{-1} = g$.

- 2. Consider a field $(\mathbb{K}, +, \cdot)$ and let 0 be the unit of the + operation. Then show the following:
 - (a) for any x ∈ K, we have 0x = 0;
 Using the fact that 0 is a unit, the distributive property and the cancellation rule, we get

$$0x + 0 = 0x = (0 + 0)x = 0x + 0x \implies 0x = 0.$$

(b) for any x, y ∈ K, x ≠ 0 and y ≠ 0 implies xy ≠ 0;
Suppose xy = 0. Then, since both x and y have inverses, we have

$$0 = (x^{-1}y^{-1})(xy) = 1.$$

But this is not possible because by construction $1 \neq 0$. Hence it must be $xy \neq 0$.

(c) for any $x, y \in \mathbb{K}$, x(-y) = (-x)y = -(xy). We need to show that x(-y) + (xy) = 0. In fact

$$x(-y) + (xy) = x(-y + y) = x0 = 0.$$

Similarly

$$(-x)y + (xy) = (-x + x)y = 0y = 0.$$

(d) (-x)(-y) = xy. Using the result of part (c), we find

$$(-x)(-y) = -[x(-y)] = -[-(xy)] = xy.$$

3. Show that D is a representation of a group G into a vector space V if and only if $D(g_1g_2^{-1}) = D(g_1)D(g_2)^{-1}$.

If D is a representation of a group G, then $D(g^{-1}) = D(g)^{-1}$ for any $g \in G$. Also, for any $g_1, g_2 \in G$

$$D(g_1g_2^{-1}) = D(g_1)D(g_2^{-1}) = D(g_1)D(g_2)^{-1}.$$

Now we deal with the reversed implication. For any $g \in G$, we have

$$D(e) = D(gg^{-1}) = D(g)D(g)^{-1} = 1$$
.

Also, for any $g_1, g_2 \in G$, we have

$$D(g_1g_2) = D\left(g_1(g_2^{-1})^{-1}\right) = D(g_1)D(g_2^{-1})^{-1} = D(g_1)\left(D(g_2)^{-1}\right)^{-1} = D(g_1)D(g_2).$$

4. Consider the following map $D : \mathbb{Z}_3 \to GL(3, \mathbb{C})$ given by

$$D(e) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \qquad D(a) = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}, \qquad D(b) = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}.$$

Show that D is a representation of \mathbb{Z}_3 .

By construction D(e) = 1. An explicit calculation gives D(a)D(b) = 1, which gives the multiplication table of \mathbb{Z}_3 .

5. Let V be a real vector space and g a scalar product. Let $\{e_i\}_{i=1,2,\dots,n}$ be an orthornomal basis and let us define the matrix $g_{ij} = g(e_i, e_j) = \pm \delta_{ij}$. By construction, the matrix g is its own inverse, in fact $g^2 = \mathbb{1}$.

(a) Consider a vector $u = u_i e_i \in V$. Show that $u_i = g_{ij}g(e_j, u)$.

$$u = u_i e_i \implies g(e_j, u) = u_k g_{kj} \implies g_{ij} g(e_j, u) = u_k g_{ij} g_{jk} = u_k \delta_{ik} = u_i$$

(b) Let M be an orthogonal operator, and let us define the matrix $M_{ij} = g(e_i, Me_j)$. Show that

$$M^T g M = g$$
.

where g is the matrix whose components are the g_{ij} . The statement to prove is better understood by writing indexes separately.

$$g_{kl}M_{ki}M_{lj} = g_{ij} \,,$$

Since M is orthogonal, $g(Me_i, Me_j) = g(e_i, e_j) = g_{ij}$. Using the result of part (a), we have

$$Me_i = (Me_i)_l e_l = g_{kl}g(e_k, Me_i)e_l$$
.

This gives, using the bilinearity of the scalar product,

$$g(Me_i, Me_j) = g_{kl}g(e_k, Me_i)g(e_l, Me_j) = g_{kl}M_{ki}M_{lj} = g_{ij},$$

which is what we needed to show.

6. Let $(A^{\dagger})_{ij}$ be the matrix associated to the adjoint of the operator A. Show that $(A^{\dagger})_{ij} = A^*_{ij}$.

$$\langle A^{\dagger} \rangle_{ij} \equiv \langle e_i | A^{\dagger} e_j \rangle = \langle A e_i | e_j \rangle = \langle e_j | A e_i \rangle^* = A^*_{ij}.$$

7. Let U be an anti-unitary operator. Show that $U^{\dagger} = U^{-1}$. For any v, w in a Hilbert space, the unitarity of U implies

$$\langle Uv|Uw\rangle = \langle v|w\rangle^*$$
.

From the definition of the adjoint of an anti-unitary operator, we have

$$\langle v|w\rangle^* = \langle Uv|Uw\rangle = \langle U^{\dagger}Uv|w\rangle^* \implies U^{\dagger}U = \mathbb{1}.$$