

keV Neutrino Model Building



Alexander Merle

University of Southampton, U.K.

E-mail: A.Merle (AT) soton.ac.uk



Based on:

AM, Niro: JCAP **1107** (2011) 023

Lindner, **AM**, Niro: JCAP **1101** (2011) 034

King, **AM**: JCAP **1208** (2012) 016

AM: J. Phys. Conf. Ser. 375 (2012) 012047

AM: Phys. Rev. **D86** (2012) 121701(R)

AM, Niro: 1302.2032

AM: 1302.xxxx

Seminar, University of Sussex, 11-02-2013

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2. keV and/or Warm Dark Matter
3. Model building for keV neutrinos
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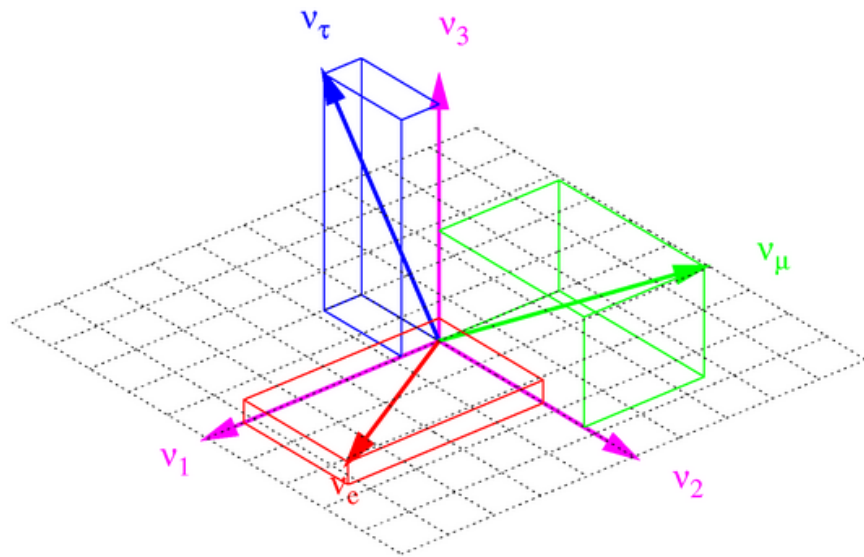
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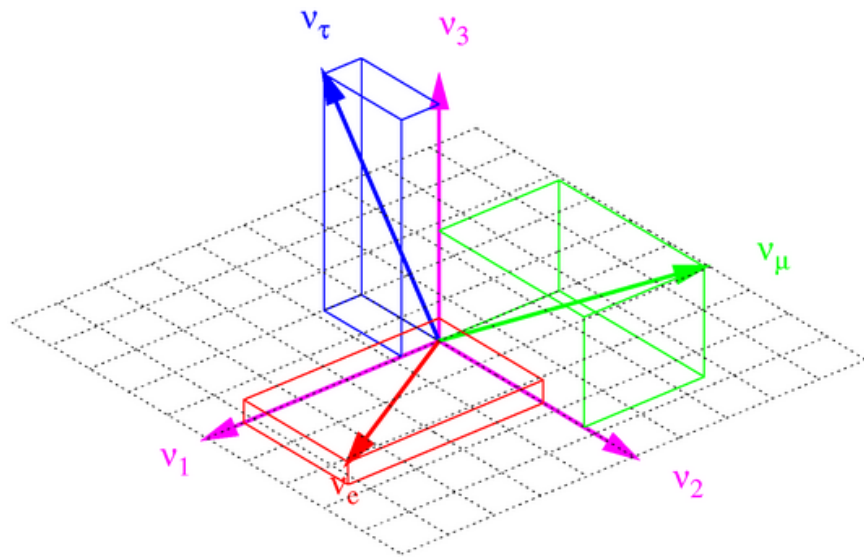


(<http://nu.phys.laurentian.ca/~fleurot/oscillations/>)

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$$\theta_{12} \approx 34.4^\circ$$

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$$\theta_{23} \approx 51.1^\circ$$

$$\Delta m_{21}^2 \approx 7.6 \times 10^{-5} \text{eV}^2$$

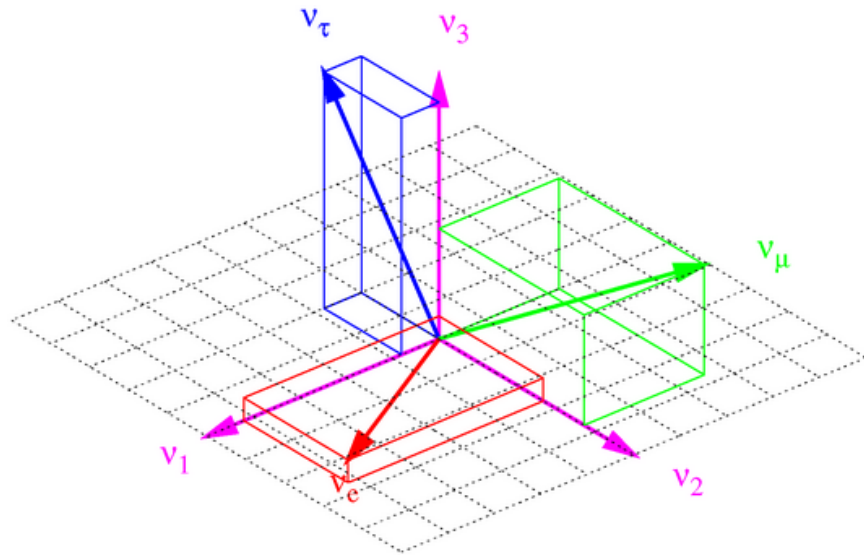
$$|\Delta m_{31}^2| \approx 2.5 \times 10^{-3} \text{eV}^2$$

Forero, Tórtola, Valle:
Phys. Rev. **D86** (2012) 073012

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BUT: We don't understand these values!!!



(<http://brainstunts.blogspot.co.uk/2011/02/angry.html>)

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$|m_{ee}| < 0.3\text{-}0.6 \text{ eV}$ [KamLAN-Zen: Phys. Rev. **C85** (2012) 045504]

$|m_{ee}| < 0.140\text{-}0.380 \text{ eV}$ [EXO-200: Phys. Rev. Lett. **109** (2012) 032505]

$|m_{ee}| < 0.300\text{-}0.710 \text{ eV}$ [CUORECINO: Astropart. Phys. **34** (2011) 822-831]

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$\Sigma < 0.58 \text{ eV}$ [WMAP, Astrophys. J. Suppl. **192** (2011) 18]

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BUT: We don't know why it is so small!!!

(<http://imprintrainingcenter.blogspot.co.uk/2010/12/understanding-and-controlling-anger.html>)

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(<http://2.bp.blogspot.com/-WTeCZueCvFI/T5fSKtzDwOI/AAAAAAAAAf8/3zpFpaUaHUI/s1600/hulk-marvel-uk.jpg>)

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(<http://www.duckipedia.de/images/e/e9/Daniel%C3%BCsentrrieb.jpg>)

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- lepton mixing: flavour symmetries, anarchy, radiative transmission, GUTs,...
- neutrino mass: seesaw(s), loop masses, R-parity violation, broken symmetries, Dark Energy connection,...
- Dark Matter: WIMPs, FIMPs, EWIPs, WIMPzillas, keVins,...
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- ☺ appeal, testability, missing links,...
- ☹ difficult, sometimes complicated,...



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- Dwarf galaxies ✓ (?)
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Still okay.

I don't wanna enter that debate... **NOBODY KNOWS IT FOR SURE!!!**
→ *As long as something is not excluded, I do not see any problem in thinking about it. Maybe we can exclude it with particle physics.*

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- Some model-independent data analysis (however,
by WDM fans...) point towards the keV scale [de Vega,
Sanchez: Mon. Not. Roy. Astron. Soc. **404** (2010) 085; de Vega,Salucci,Sanchez:
New Astron. **17** (2012) 653]

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Simple framework: **ν MSM** [Asaka,Blanchet,
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- BUT: **keV mass not explained**
 - GeV-degeneracy not explained**
 - ν -masses not explained**
 - hardly testable** ➔ **model building needed!!!**

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- **METHODS**: use what is known from light neutrino models to explain the desired structure → flavour symmetries, mass suppression mechanisms, ...
- **LINK**: be careful to accommodate the information from astrophysics/cosmology DM-production mechanism, correct abundance, X-ray bound, structure formation, ...

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- *Differences to “ordinary” model building:*

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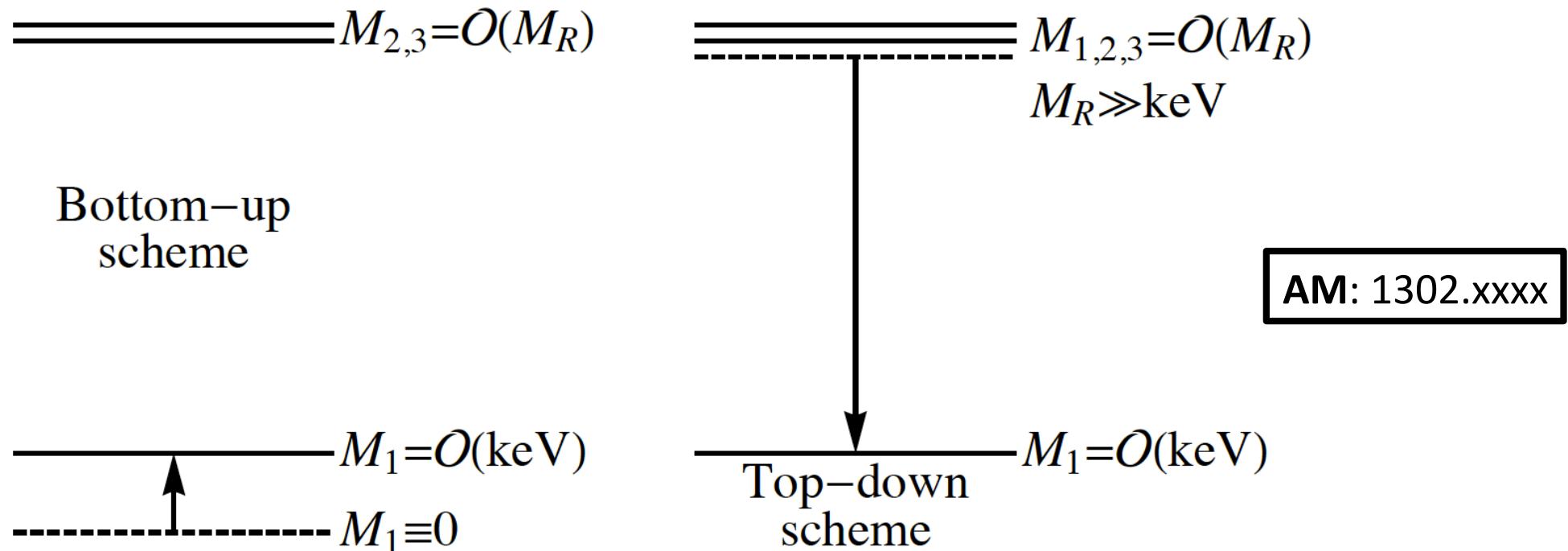
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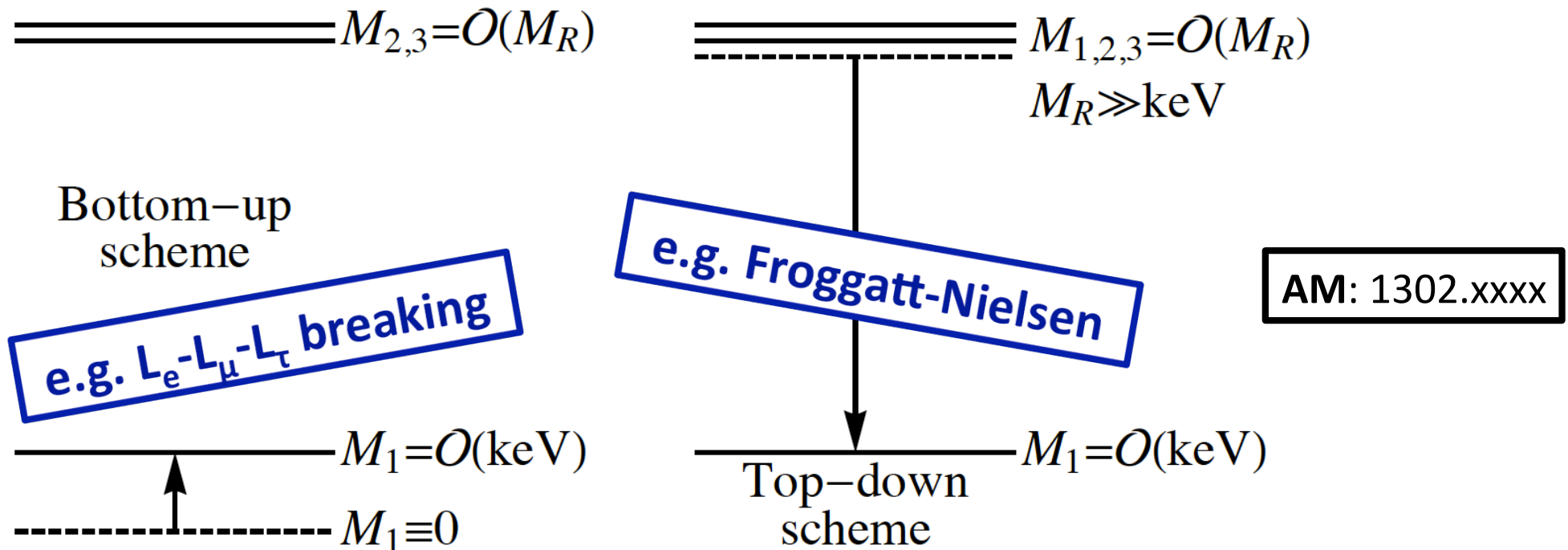
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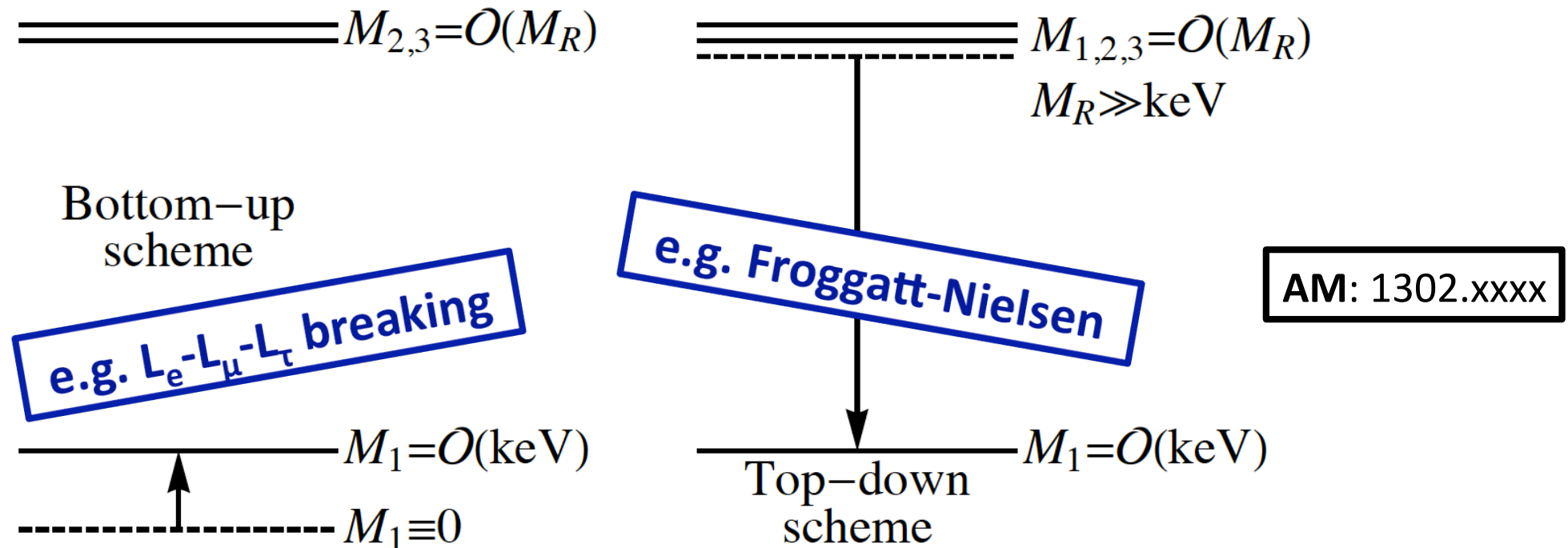
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→ Most models are in one or the other category!

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 - a sterile neutrino N_1 can decay: $N_1 \rightarrow \nu\gamma$
 - ➔ this produces a monoenergetic X-ray line with $E = M_1/2$ ($N_1 = \text{DM} \rightarrow$ many around)

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 - ➔ this produces a monoenergetic X-ray line with $E = M_1/2$ ($N_1 = \text{DM} \rightarrow$ many around)
 - HOWEVER: this line is **NOT** observed
 - ➔ strong bound on active-sterile mixing $\theta_i^2 = \sum_\alpha |\theta_{\alpha i}|^2$ with the keV neutrino:

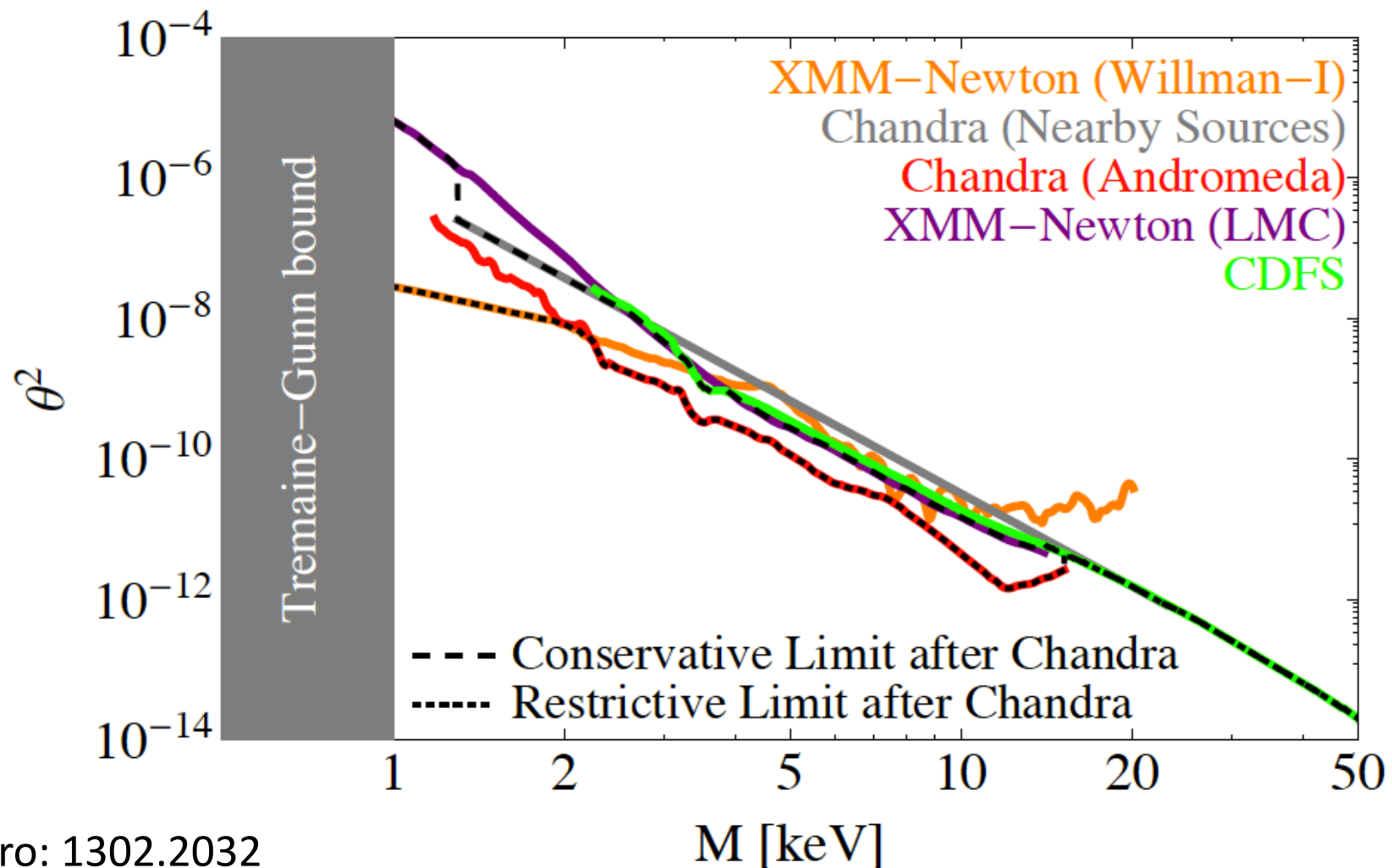
$$\theta_1^2 \lesssim 1.8 \cdot 10^{-5} \left(\frac{1 \text{ keV}}{M_1} \right)^5$$

Boyarsky, Ruchayskiy, Shaposhnikov:
Ann. Rev. Nucl. Part. Sci. **59** (2009) 191

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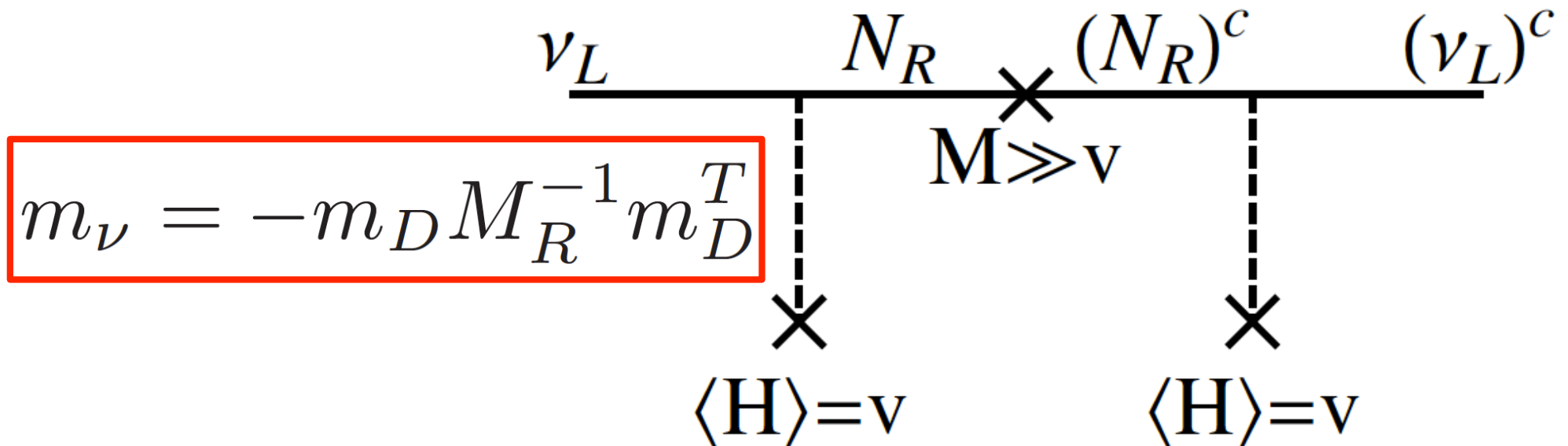
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The diagram illustrates the seesaw mechanism for neutrino mass generation. It features a horizontal line representing the fermion propagator with four external labels: ν_L on the left, N_R in the middle, $(N_R)^c$ on the right, and $(\nu_L)^c$ on the far right. A large 'X' is placed over the N_R and $(N_R)^c$ labels, with the text $M \gg v$ below it. Two vertical dashed lines extend downwards from the horizontal line, each ending in a large 'X' and labeled $\langle H \rangle = v$. To the left of the diagram, a red-bordered box contains the equation $m_\nu = -m_D M_R^{-1} m_D^T$.

$$m_\nu = -m_D M_R^{-1} m_D^T$$

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→ Does that also work when “dividing by keV mass“?!?

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- [Kusenko,Takahashi,Yanagida: Phys. Lett. **B693** (2010) 144]
[AM,Niro: JCAP **1107** (2011) 023]

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→ *Actually okay in most of the cases!*

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[Asaka,Shaposhnikov,Kusenko: Phys. Lett. **B638** (2006) 401]
[Anisimov,Bartocci,Bezrukov: Phys. Lett. **B671** (2009) 211]
[Bezrukov,Gorbunov: JHEP **1005** (2010) 010]

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[Bezrukov,Gorbunov: JHEP **1005** (2010) 010]
- **thermal overproduction with entropy dilution**
[Bezrukov,Hettmansperger,Lindner: Phys. Rev. **D81** (2010) 085032]
[Nemevsek,Senjanovic,Zhang: JCAP **1207** (2012) 006]

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- *In this part of the talk:*

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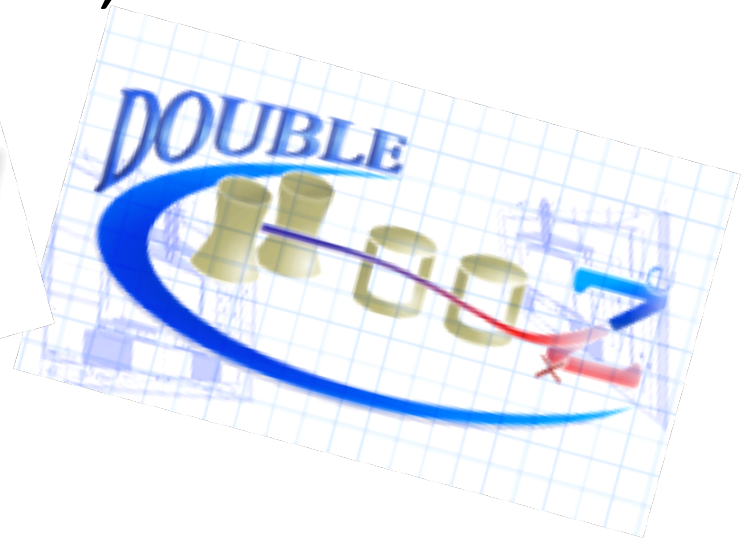
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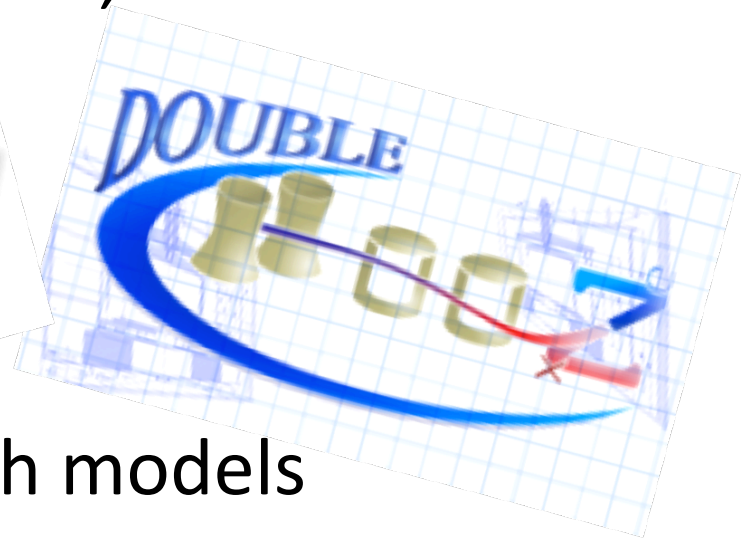
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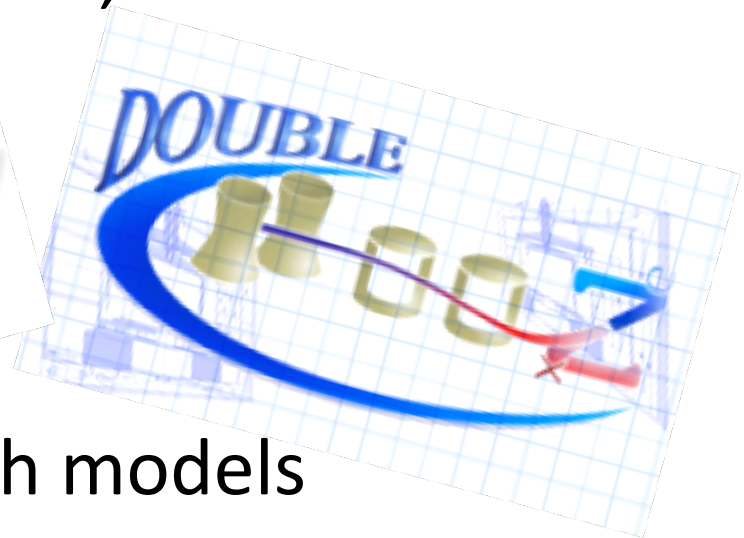


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- ... Now let's start!

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 - original: [Petcov: Phys. Lett. **B110** (1982) 245]
 - 2 RH neutrinos: [Grimus, Lavoura: JHEP **0009** (2000) 007]
 - 3 RH neutrinos:
 - [Barbieri, Hall, Tucker-Smith, Strumia, Weiner: JHEP **9812** (1998) 017]
 - [Mohapatra: Phys. Rev. **D64** (2001) 091301]

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 - [Barbieri, Hall, Tucker-Smith, Strumia, Weiner: JHEP **9812** (1998) 017]
 - [Mohapatra: Phys. Rev. **D64** (2001) 091301]
 - *application to keV sterile neutrinos*:
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 - [Lindner, **AM**, Niro: JCAP **1101** (2011) 034]

4. Example Models

- probably the most intuitive: $\mathcal{F} = L_e - L_\mu - L_\tau$
 - original: [Petcov: Phys. Lett. **B110** (1982) 245]
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 - general features:
 - *symmetry* \rightarrow patterns: $(0, m, m)$ & $(0, M, M)$
 - *broken* \rightarrow small mass, degeneracy lifted

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 - charge assignment under global U(1) [or: Z_4]:

	L_{eL}	$L_{\mu L}$	$L_{\tau L}$	e_R	μ_R	τ_R	N_{1R}	N_{2R}	N_{3R}	ϕ	Δ
\mathcal{F}	1	-1	-1	1	-1	-1	1	-1	-1	0	0

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- then, only symmetry preserving terms are allowed:

$$\mathcal{L}_{\text{mass}} = -\frac{1}{2} \overline{\Psi^C} \mathcal{M}_\nu \Psi + h.c.$$

with: $\Psi \equiv ((\nu_{eL})^C, (\nu_{\mu L})^C, (\nu_{\tau L})^C, N_{1R}, N_{2R}, N_{3R})^T$

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$$\mathcal{M}_\nu = \left(\begin{array}{ccc|ccc} 0 & m_L^{e\mu} & m_L^{e\tau} & m_D^{e1} & 0 & 0 \\ m_L^{e\mu} & 0 & 0 & 0 & m_D^{\mu2} & m_D^{\mu3} \\ m_L^{e\tau} & 0 & 0 & 0 & m_D^{\tau2} & m_D^{\tau3} \\ \hline m_D^{e1} & 0 & 0 & 0 & M_R^{12} & M_R^{13} \\ 0 & m_D^{\mu2} & m_D^{\tau2} & M_R^{12} & 0 & 0 \\ 0 & m_D^{\mu3} & m_D^{\tau3} & M_R^{13} & 0 & 0 \end{array} \right)$$

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 - eigenvalues of \mathcal{M}_ν (with μ - τ symmetry):
 - light neutrinos: $(\lambda_+, \lambda_-, 0)$
 - heavy neutrinos: $(\Lambda_+, \Lambda_-, 0)$
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 - WAY OUT: **broken symmetry**
 - \rightarrow will remedy the above issues
 - \rightarrow important: no matter how the breaking is achieved, the results will always look similar

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→ new mass matrix:

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→ new eigenvalues: $\Lambda_s = S$ $\Lambda'_\pm = S \pm \sqrt{2}M_R$

$$\lambda_s = s \quad \lambda'_\pm = s \pm \sqrt{2} \left[m_L - \frac{m_D^2}{M_R} \right] + \frac{5m_D^2 S}{4M_R^2}$$

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 - mixings also require soft breaking:

$$\mathcal{M}_l \mathcal{M}_l^\dagger \simeq \begin{pmatrix} m_e^2 + m_\mu^2 \lambda^2 & m_\mu^2 \lambda & 0 \\ m_\mu^2 \lambda & m_\mu^2 & 0 \\ 0 & 0 & m_\tau^2 \end{pmatrix}$$

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$$\lambda = \theta_{12} - \pi/4$$

$$|U_{e3}| \simeq \frac{\lambda}{\sqrt{2}} \rightarrow \theta_{13} \simeq 8^\circ,$$

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- prediction for the masses (under assumptions):

$$|m_1| = 0.0486 \text{ eV}, |m_2| = 0.0494 \text{ eV}, \text{ and } |m_3| = 0.0004 \text{ eV}$$

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○ mass shifting scheme:



$$M_3 \approx M_2$$

$$M_2 = M_3 \gtrsim \text{GeV}$$

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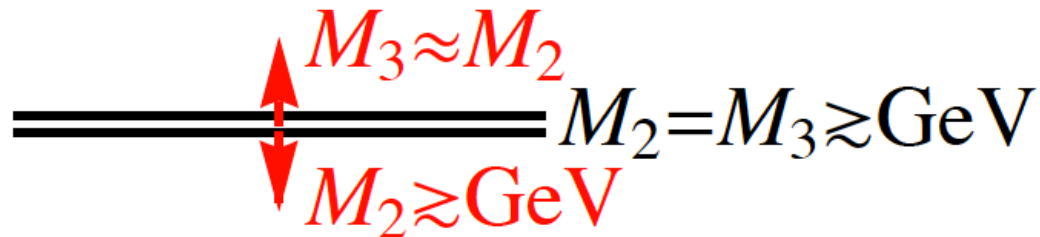
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➔ clear bottom-up type scheme

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- *probably the most effective*: **Split Seesaw**

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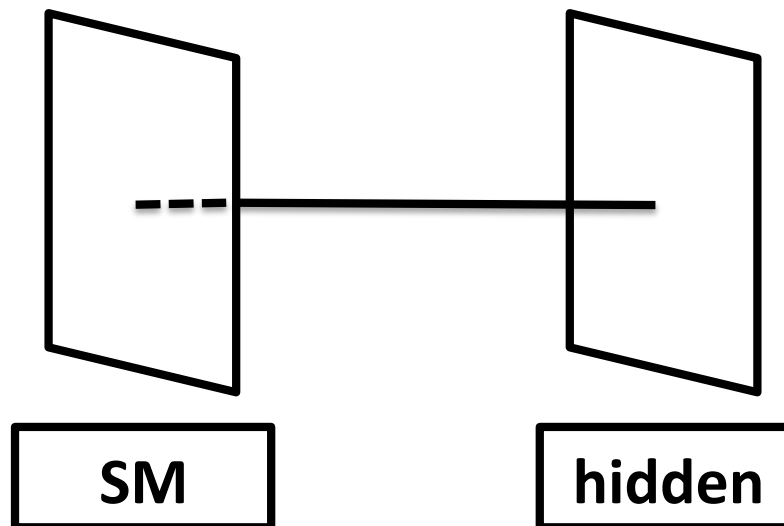
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[Kusenko,Takahashi,Yanagida: Phys. Lett. **B693** (2010) 144]

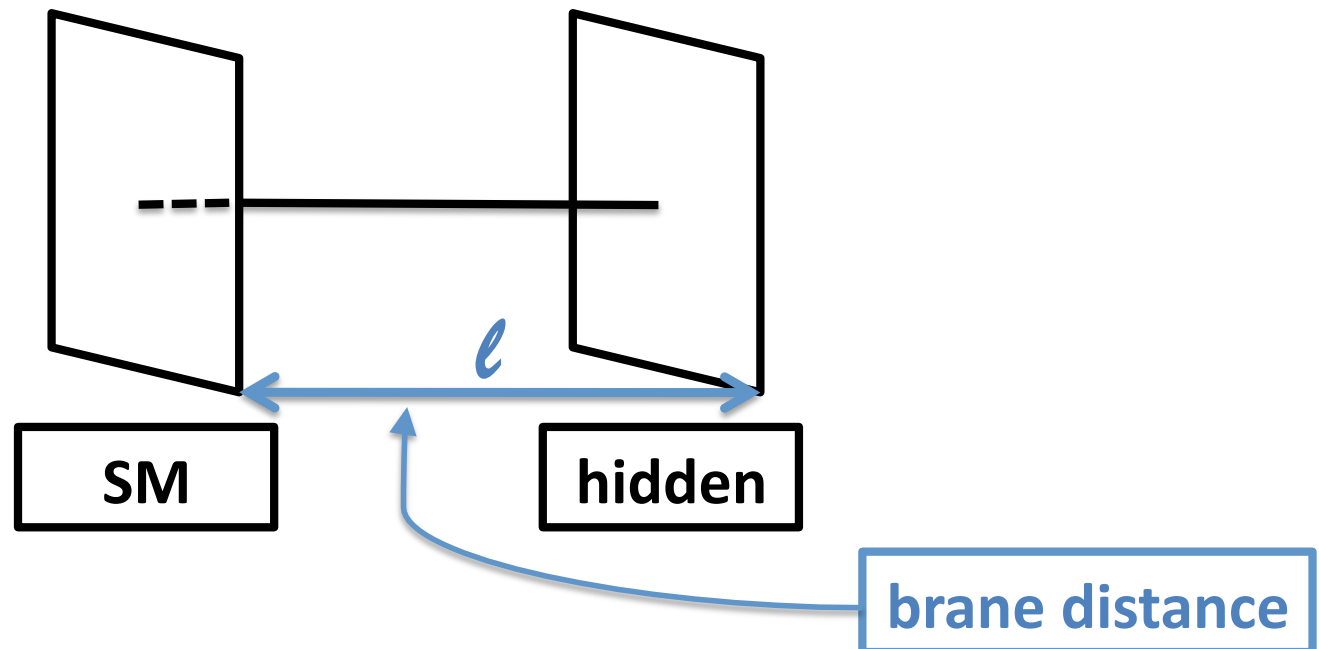
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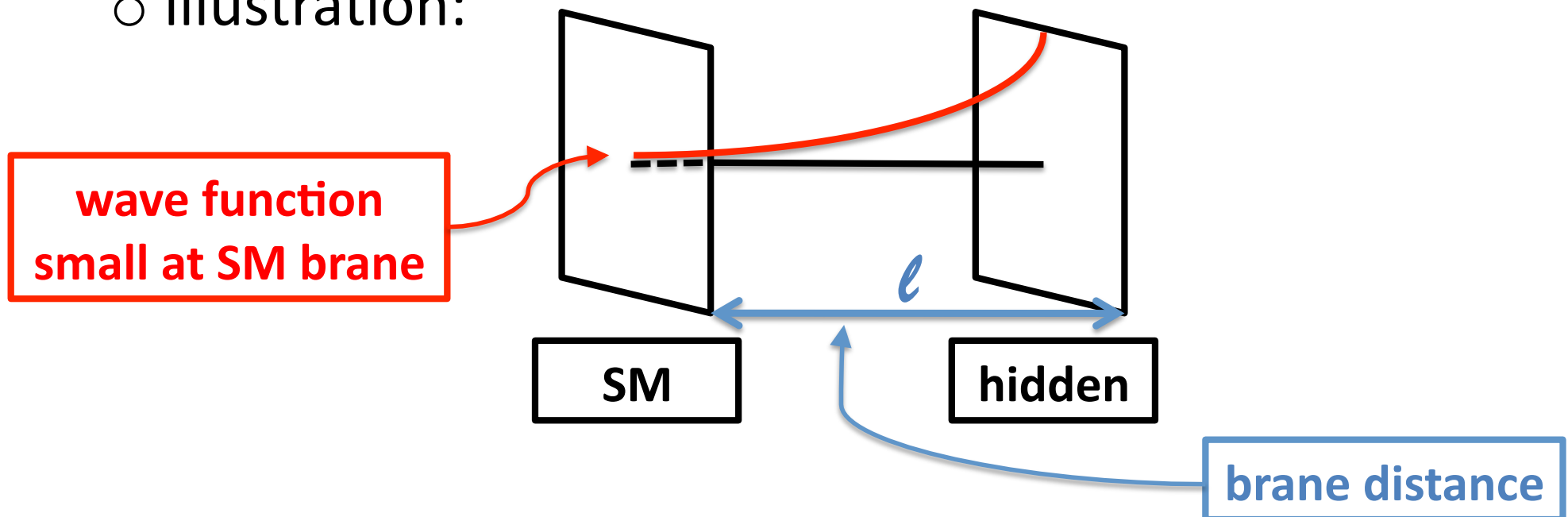


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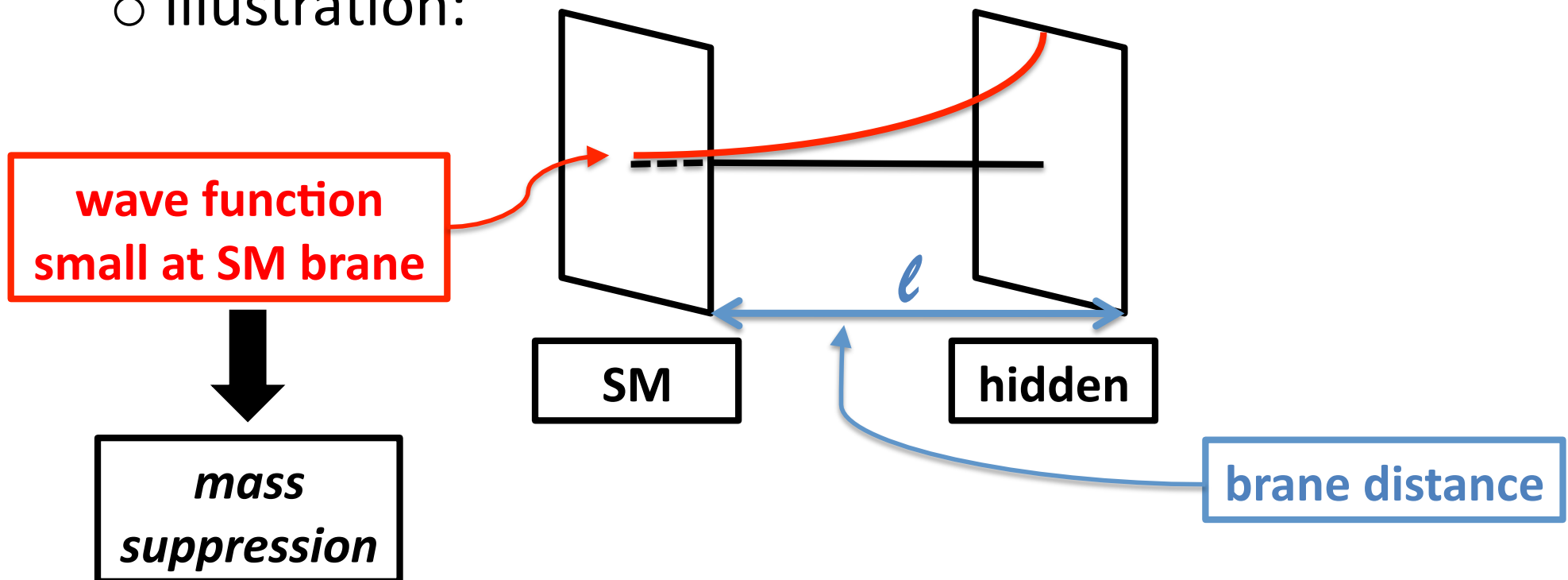


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- starting point: 5D action

$$S = \int d^4x \int_0^l dy M_0 (i\bar{\Psi}\Gamma^A\partial_A\Psi - m\bar{\Psi}\Psi)$$

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- Fourier expansion of the field:

$$\Psi_{L,R}(x^\mu, y) = \sum_n \psi_{L,R}^{(n)}(x^\mu) f_{L,R}^{(n)}(y)$$

➔ equation of motion in the Extra Dimension:

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- solution (“**bulk profile**”) for the zero mode: $m_n=0$

$$\boxed{f_{L,R}^{(0)}(y) = C e^{\mp my}} \quad C = \sqrt{\frac{2m}{e^{2ml} - 1}} \frac{1}{\sqrt{M_0}}$$

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$$S = \int d^4x \int dy \left[M_0 \left(\overline{\Psi_{iR}^{(0)}} i\Gamma^A \partial_A \Psi_{iR}^{(0)} - m_i \overline{\Psi_{iR}^{(0)}} \Psi_{iR}^{(0)} \right) \right. \\ \left. - \delta(y) \left(\frac{\kappa_i}{2} v_{B-L} (\overline{\Psi_{iR}^{(0)}})^c \Psi_{iR}^{(0)} + \tilde{\lambda}_{i\alpha} \overline{\Psi_{iR}^{(0)}} L_\alpha H \right) \right]$$

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STRONG SUPPRESSION!!!

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$$m_3 < m_2 < m_1 \rightarrow M_3 \gg M_2 \gg M_1!!!$$

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- bonus: seesaw *guaranteed* to work, due to conspiracy between the suppressions

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- **issue #1**: slight enhancement of active-sterile mixing

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→ can be cured by A_4 extension:

$$m_1 > m_2 = m_3 \rightarrow M_1 \ll M_2 = M_3$$

[Adulpravitchai, Takahashi: JHEP **1109** (2011) 127]

BUT: $\theta_{13}=0$, $\theta_{23}=\pi/4$, excluded by X-ray bound!

→ needs seesaw type II situation to work

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 - features:
 - suppression maybe as strong as for split seesaw
 - *stronger* enhancement of active-sterile mixing
 - more predictive than one would naively expect
 - seesaw guaranteed to work

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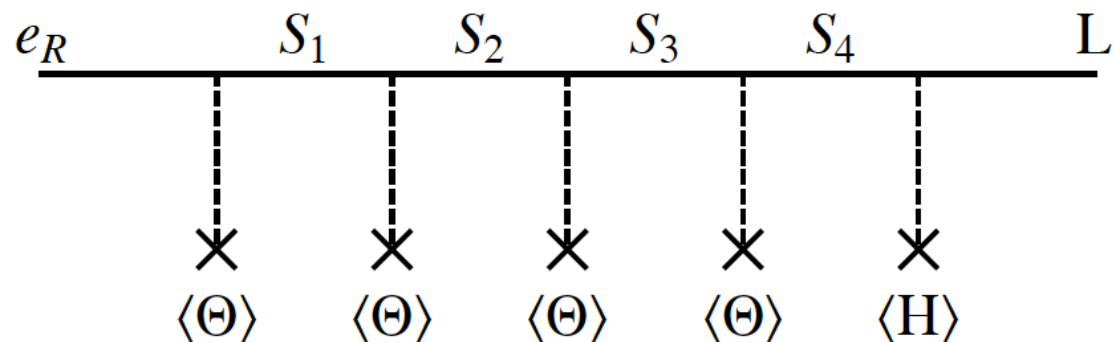
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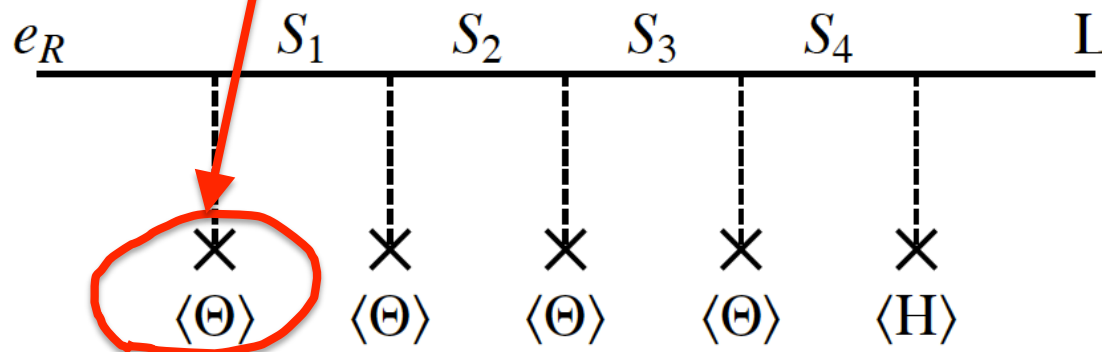
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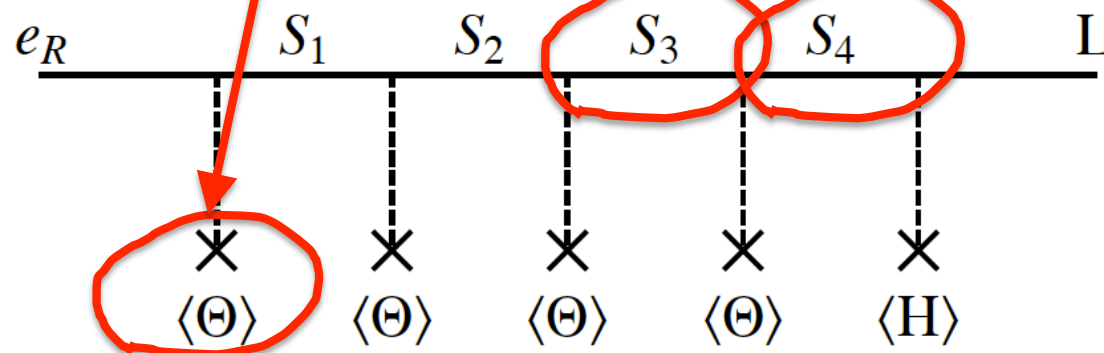
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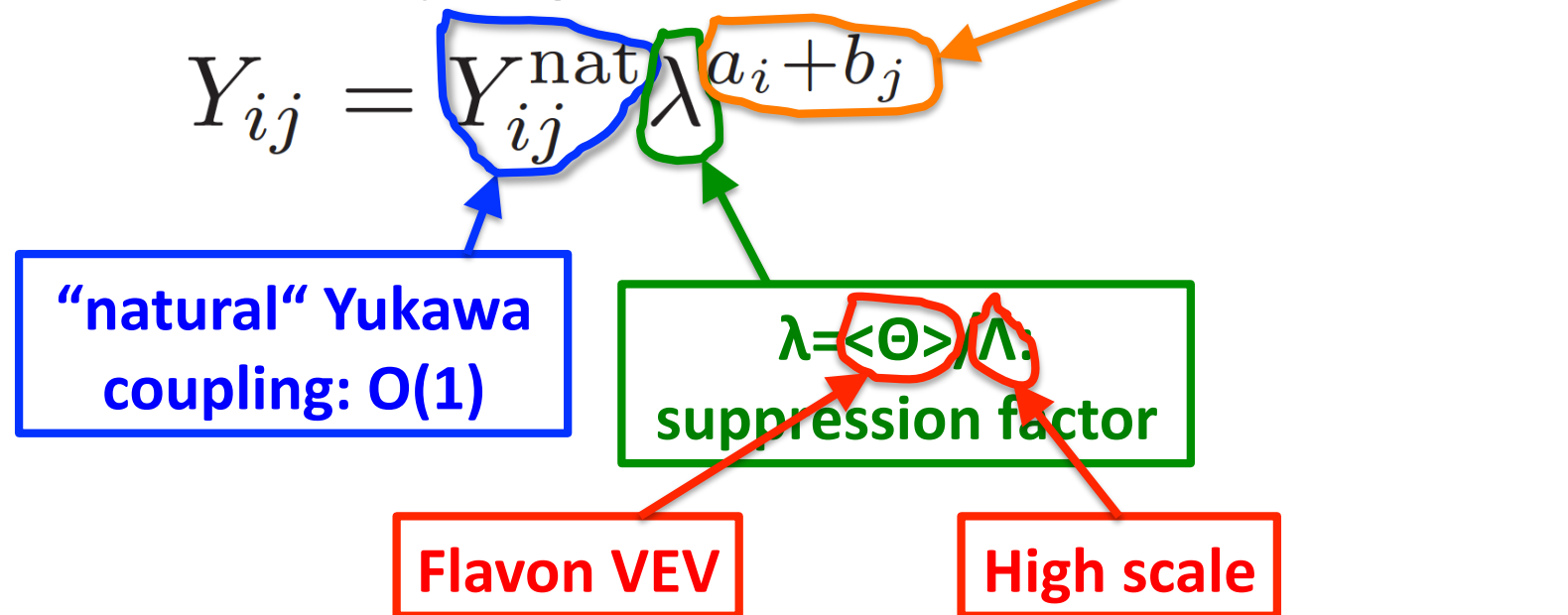
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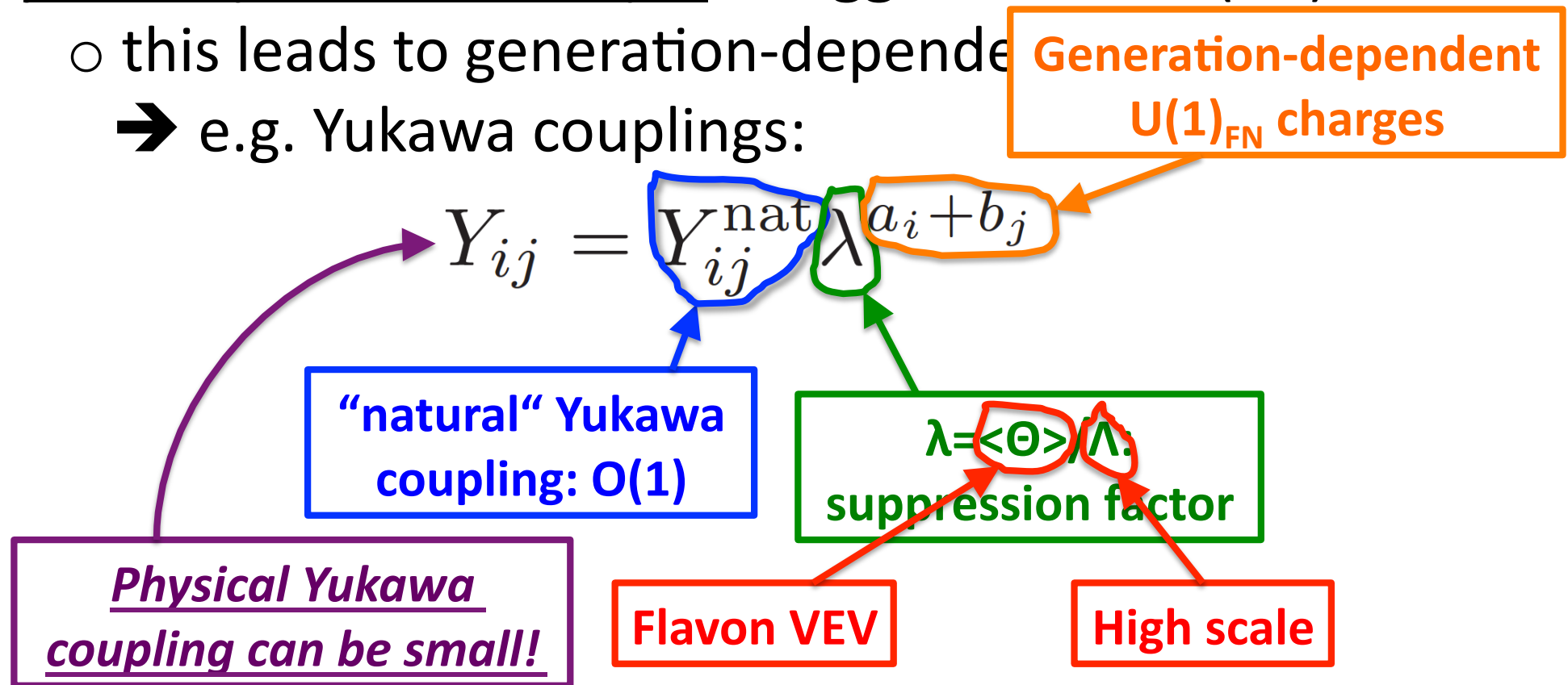


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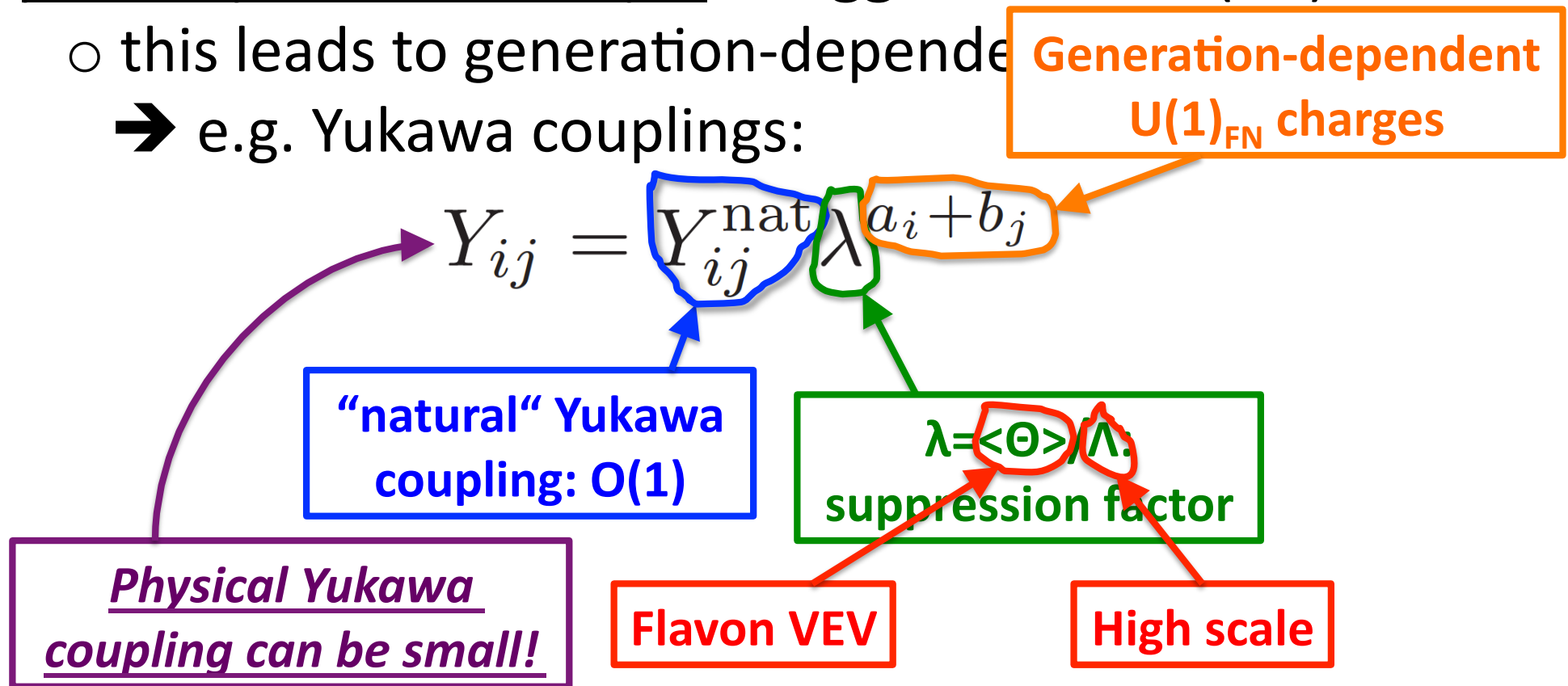


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- **HOWEVER**: several problems are swept under the carpet (UV-completion, $U(1)$ -breaking,...)

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 - application to keV sterile neutrinos: $U(1)_{\text{FN}} \times Z_{2,\text{aux}}$

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$$L_{1,2,3} : (f_1, f_2, f_3; +, +, -)$$

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Quasi-Degeneracy

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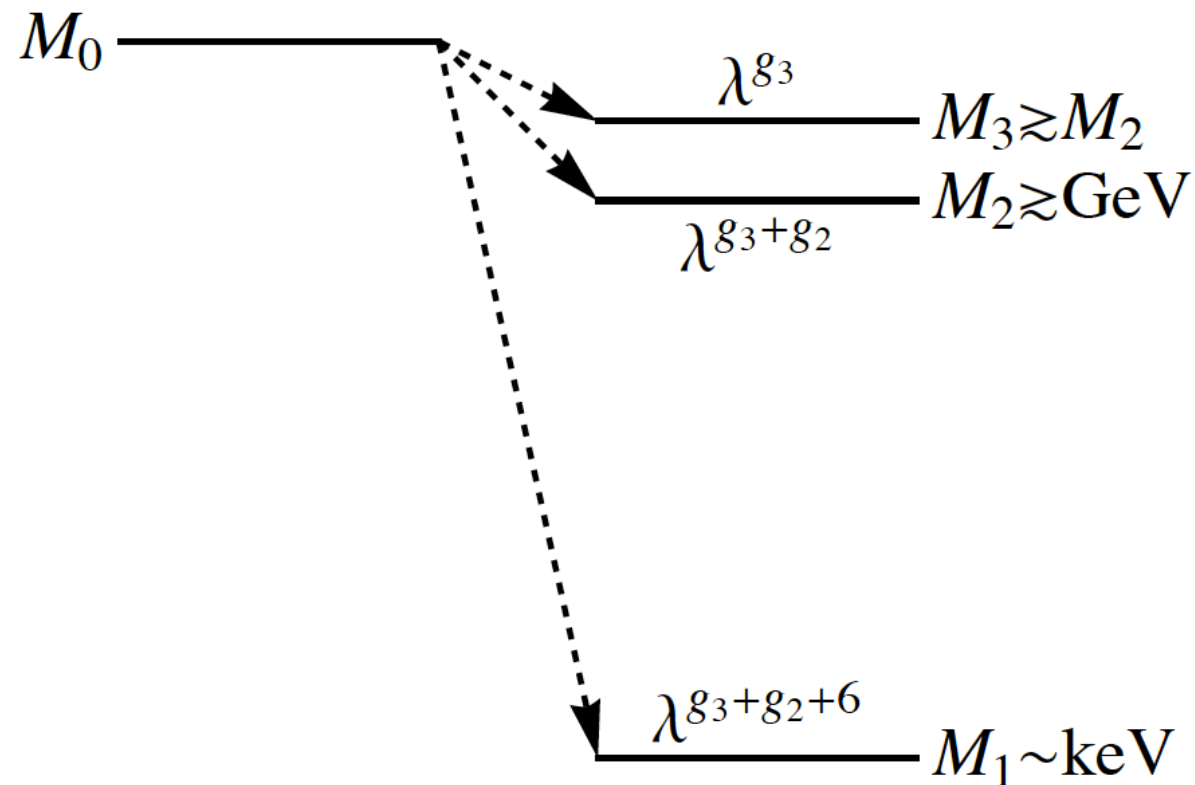
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Hierarchy

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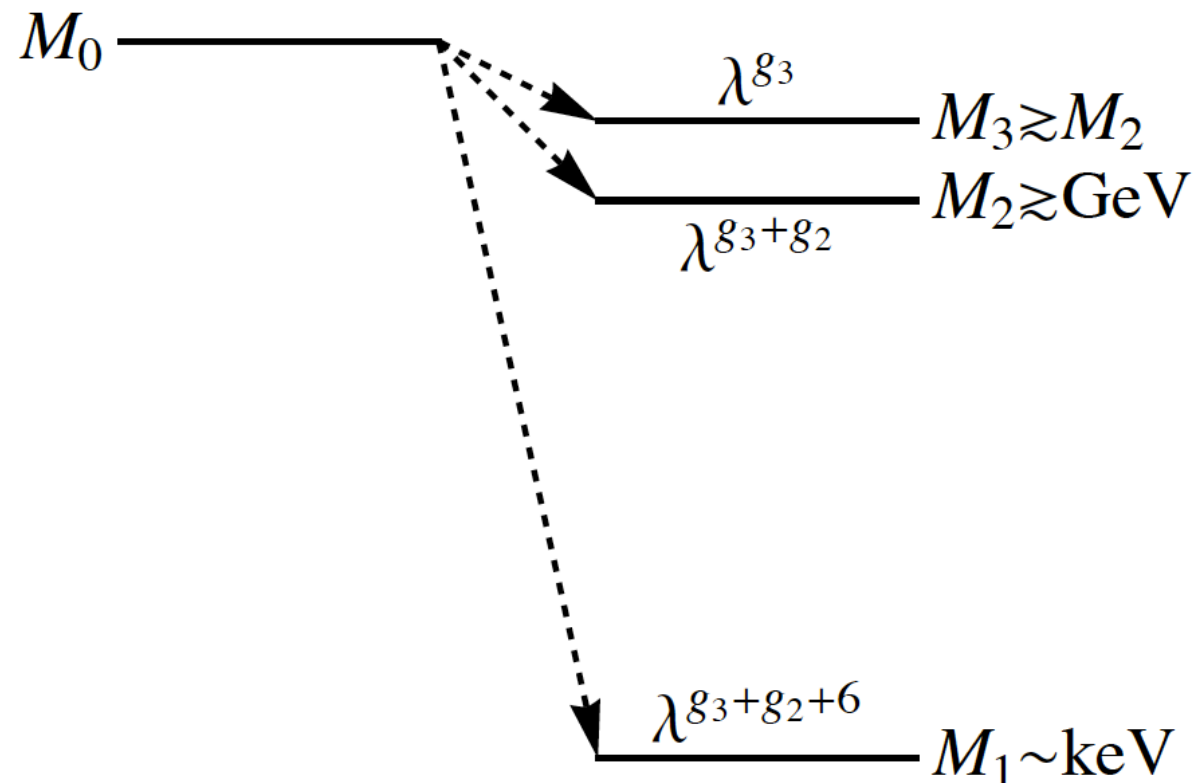
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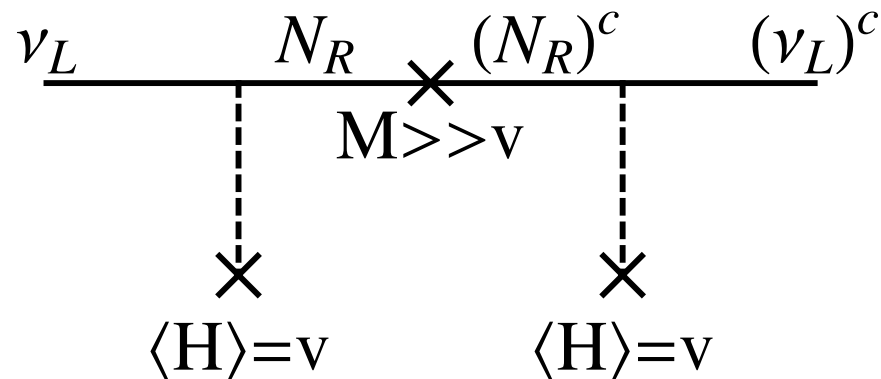


→ large mass scale gets suppressed

→ *top-down*

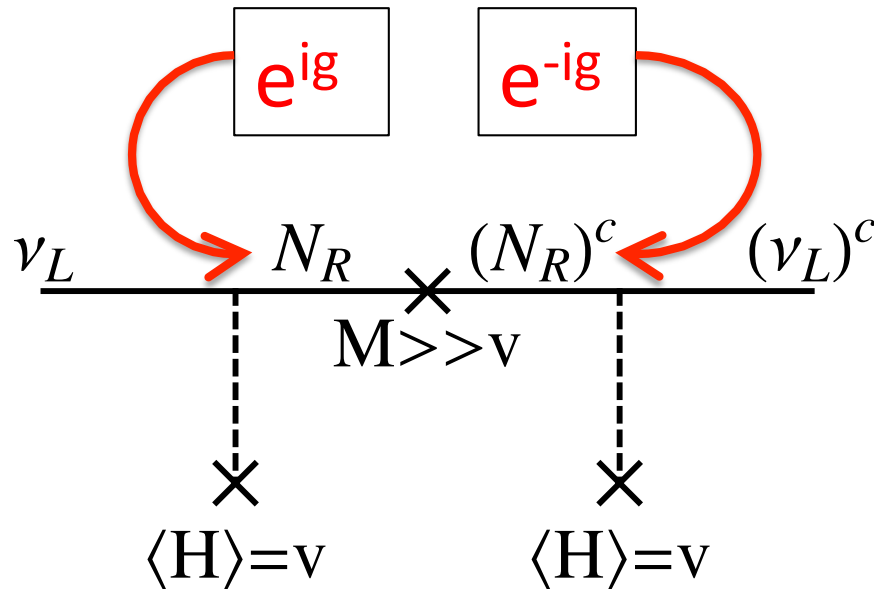
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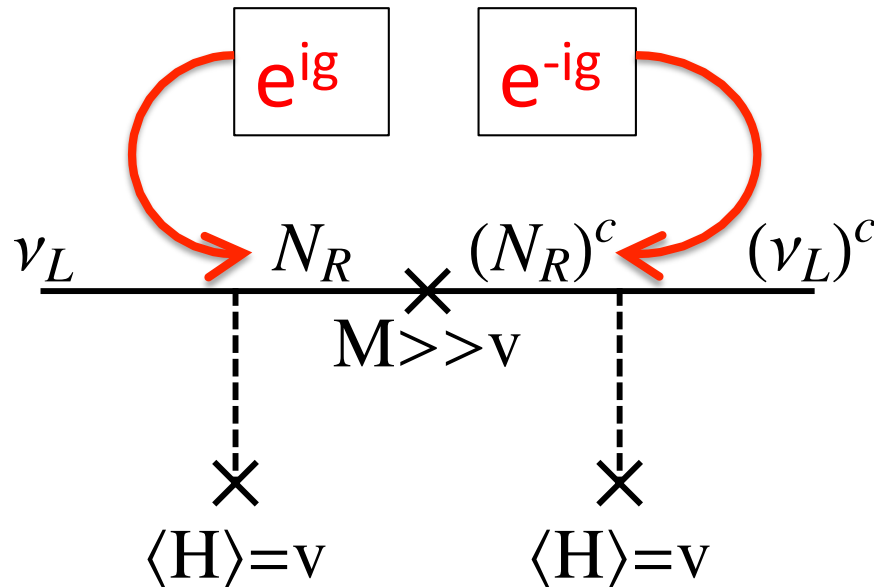
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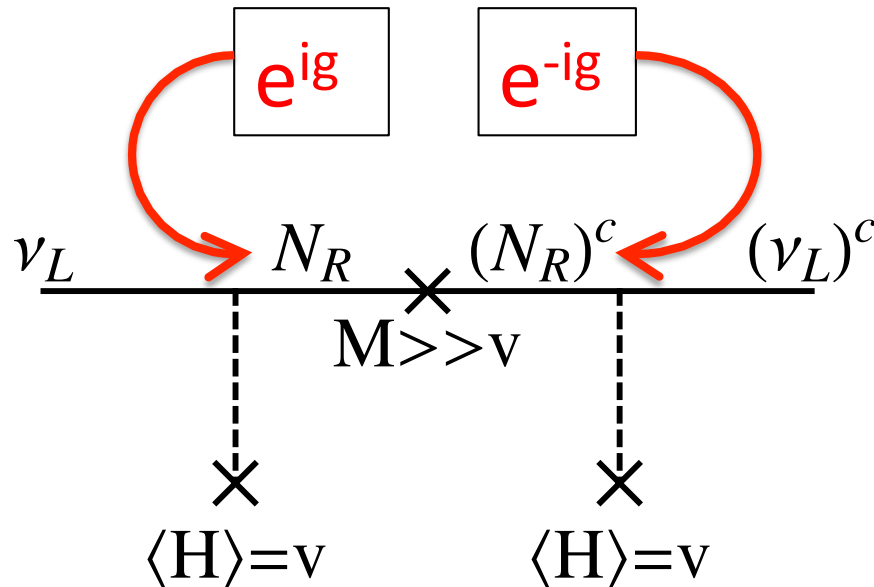
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- ➔ the charges (g_1, g_2, g_3) drop out of the light neutrino mass matrix ✓

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 - no anomalies within SU(5)

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 - anomaly-free U(1)-extension
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 - important features:
 - necessarily goes beyond 3 sterile neutrinos
 - not justified by itself → needs framework
 - structural implications (one massless ν , only possible for certain numbers of sterile ν 's)

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 - idea: introduce another singlet fermion S_R and **assume** the following Lagrangian

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- one now assumes a hierarchy: $m_D \ll M_S \ll M_R$

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$$M_{\nu}^{4 \times 4} = \begin{pmatrix} m_D M_R^{-1} m_D^T & m_D M_R^{-1} M_S^T \\ M_S (M_R^{-1})^T m_D^T & M_S M_R^{-1} M_S^T \end{pmatrix}$$

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 - general: although the mechanism cannot stand alone, it may be resembled in more concrete models

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- plus many **scenarios** (=frameworks without explanation for scale): Left-Right symmetry, scotogenic model, ...

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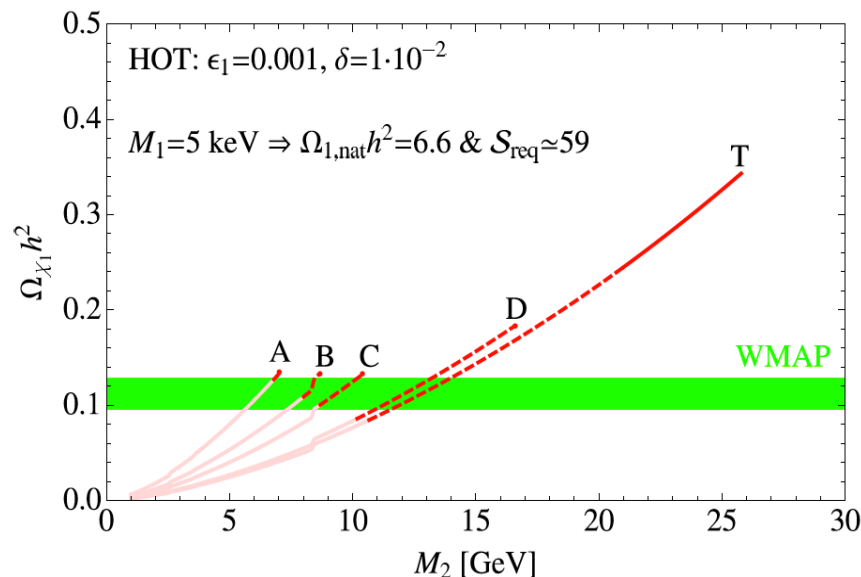
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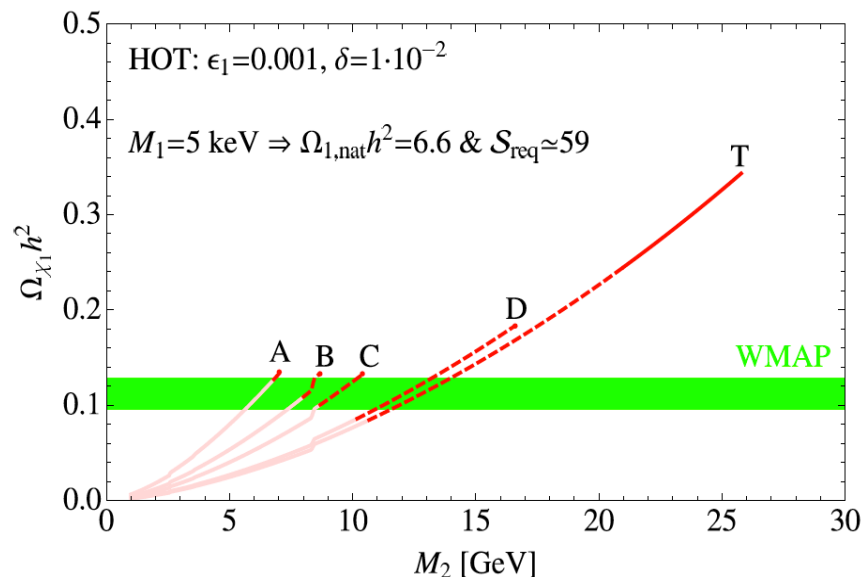
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- other production mechanisms and more or less model-independent studies possible

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**THANK
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