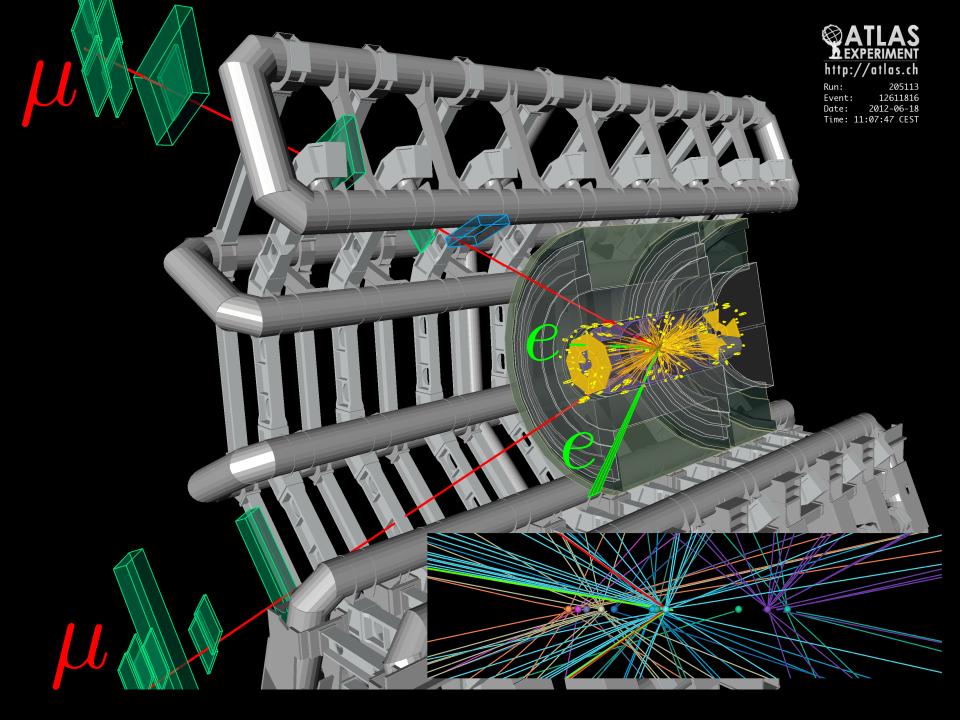
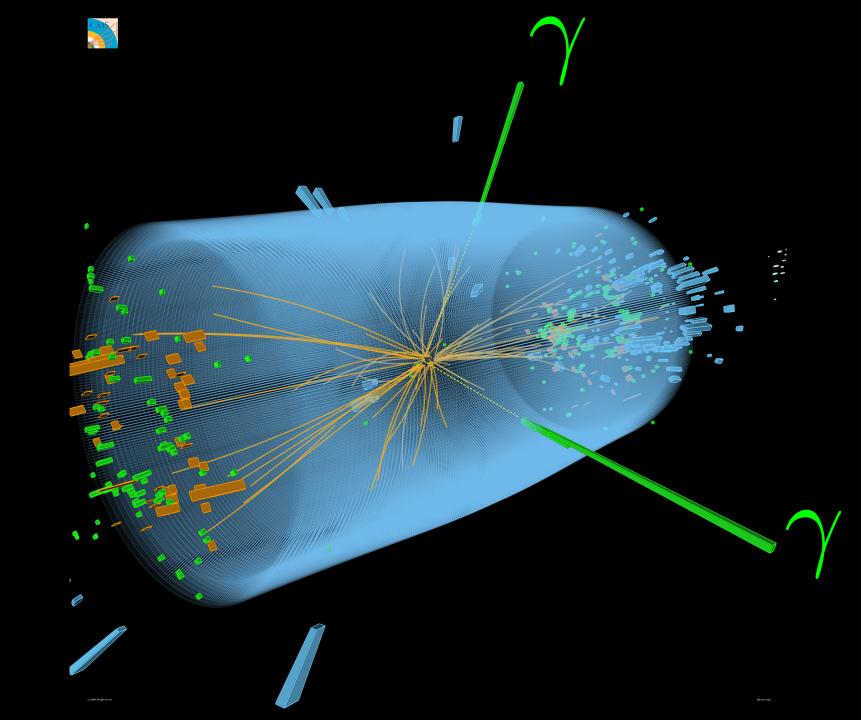


# A Higgs Near 125 GeV Beyond the MSSM

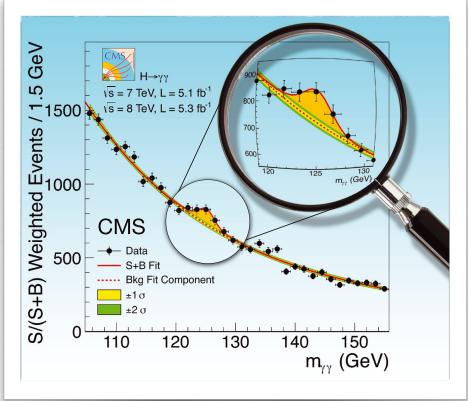
Part 1 Higgs in SM and NMSSM Part 2 Higgs and Gluinos in E6SSM

Steve King Sussex, Mon 22nd Oct 2012

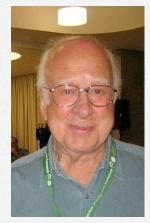




#### LHC has discovered a new particle



Congratulations to both attas and CMB Collaborations and to the builders of the LHC on a magnificent achievement!



Peter Stigge 30 August 2012

<sup>&</sup>quot; ... The decay to two photons indicates that the new particle is a boson with spin different from one. The results presented here are consistent, ... with expectations for a standard model Higgs boson."

Not only does the discovery yield the missing link to the present Standard Model theory of elementary particles, but a detailed analysis of the decays, in particular of the decay of the Scalar to two photons which is sensitive to loops of intermediated charged particles, will possibly yield information about the spectrum beyond the Standard Model.

Prof. François Englert





It is great to know that the famous boson almost certainly exists, and we are eagerly waiting for detailed measurement of its properties.

Prof. Tom Kibble

Prof. Carl R. Hagen

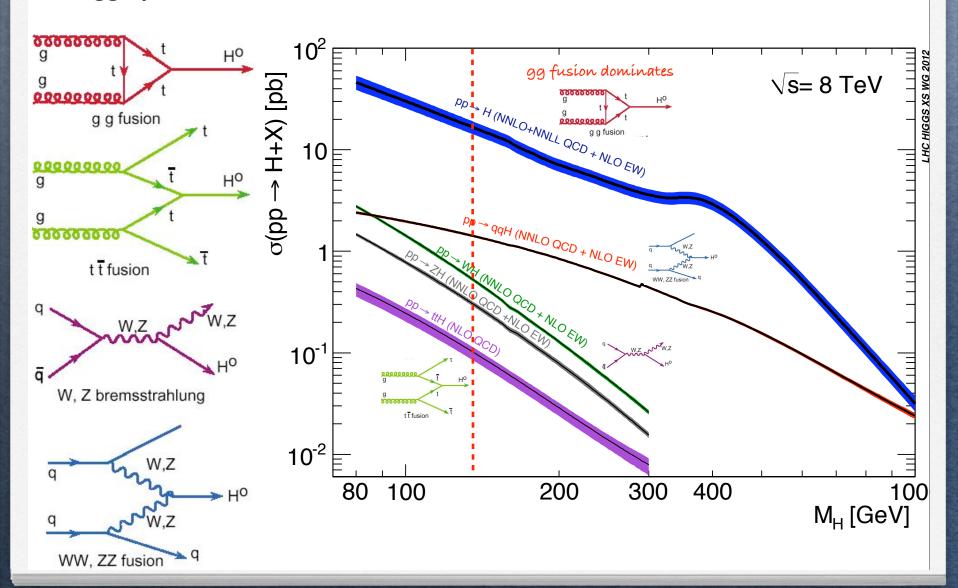
Prof. Gerald Guralnik

Tom Kilble

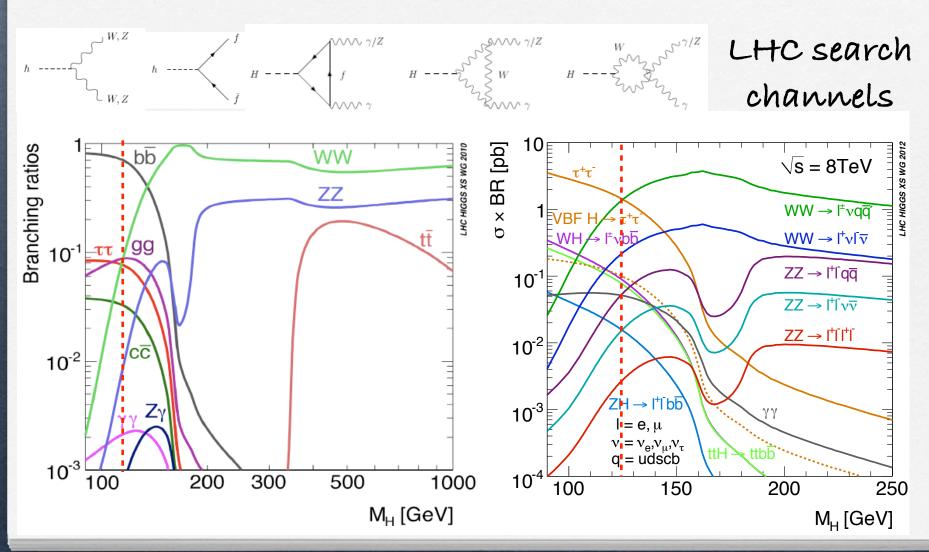
Corl RHogen

S. D. Meralniz

• Higgs production mechanisms and cross sections

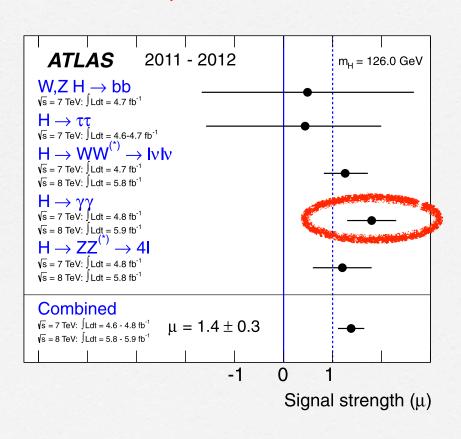


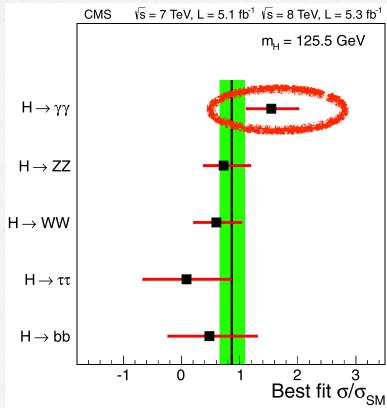
## **Higgs Decays**



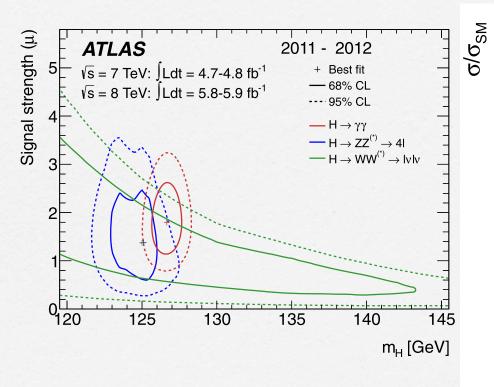
## Higgs decay LHC signal strengths

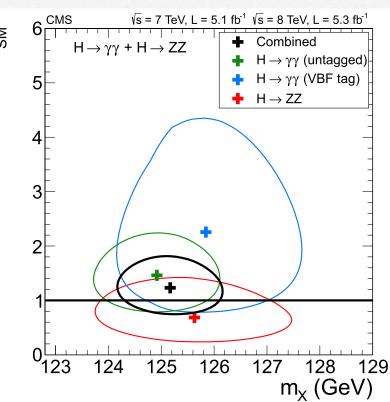
Two photon rate too high in both experiments





## Higgs mass 124-127 GeV





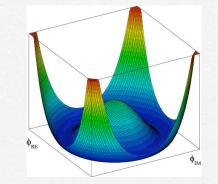
 $126.0 \pm 0.4 \text{ (stat)} \pm 0.4 \text{ (sys) GeV}$ 

 $125.3 \pm 0.4$ (stat.)  $\pm 0.5$ (syst.) GeV.

## Higgs Theory in SM

Higgs potential

$$V = m_H^2 \left| H \right|^2 + \frac{1}{2} \lambda \left| H \right|^4$$



Tree-level min cond

$$m_H^2 = -\lambda v^2 = -\lambda \left(246 \, GeV\right)^2$$

Including rad corr

$$m_H^2 + \delta m_H^2 = -\lambda \left(246 \, GeV\right)^2$$

$$H - - - \underbrace{\underbrace{\frac{Q_{3L}}{t_R}}_{Q_{3L}} - - - - H$$

$$\delta m_H^2(top \, loop) = -\frac{3}{\sqrt{2\pi^2}} G_F m_t^2 \Lambda^2 = -\left(100 \, GeV\right)^2 \left(\frac{\Lambda}{1 \, TeV}\right)^2$$

Fine-tuning is required if the cut-off  $\Lambda\gg 1\,TeV$ 

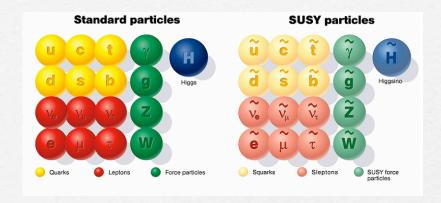
Motivates new physics at TeV scale e.g. SUSY

#### Minimal SUSY SM (MSSM)

Table 1: The MSSM Particle Spectrum

Superfield	Bosons	Fermions	
Gauge			
$\widehat{G}$	g	$\widetilde{g}$	
$\widehat{V}^a$	$W^a$	$\dfrac{\widetilde{g}}{\widetilde{W}^a}$	
$\widehat{V}'$	B	$\widetilde{B}$	

# $\frac{V'}{\text{Matter}} \qquad B \qquad B \\ \underline{\widehat{H}}_{\text{atter}} \qquad \text{leptons} \begin{cases} \widetilde{L} = (\widetilde{\nu}, \widetilde{e}^-)_L & (\nu, e^-)_L \\ \widetilde{E} = \widetilde{e}_R^+ & e_L^c \end{cases}$ $\frac{\widehat{Q}}{\widehat{Q}} \qquad \text{quarks} \begin{cases} \widetilde{Q} = (\widetilde{u}_L, \widetilde{d}_L) & (u, d)_L \\ \widetilde{U}^c = \widetilde{u}_R^* & u_L^c \\ \widetilde{D}^c = \widetilde{d}_R^* & d_L^c \end{cases}$ $\frac{\widehat{H}_d}{\widehat{H}_u} \qquad \text{Higgs} \begin{cases} H_d^i & (\widetilde{H}_d^0, \widetilde{H}_d^-)_L \\ H_u^i & (\widetilde{H}_u^+, \widetilde{H}_u^0)_L \end{cases}$



Higgs/Higgsino mass parameter

$$\mu \tilde{H}_u \tilde{H}_d$$

$$|\mu| \lesssim 200 \text{ GeV}.$$

to avoid tree-level tuning since

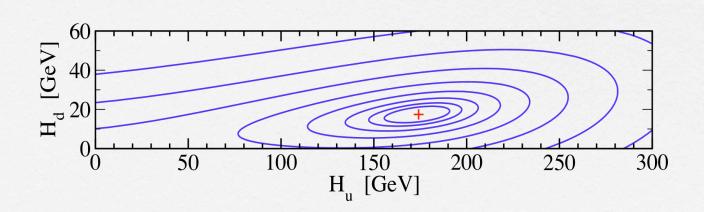
$$m_H^2 = \mu^2 + m_0^2$$

## Higgs Theory in MSSM $W = \mu \hat{H}_u \hat{H}_d$

$$H_u = (H_u^+, H_u^0)$$
  $H_d = (H_d^0, H_d^-)$ 

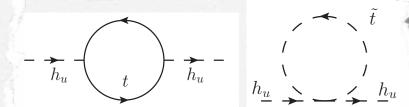
Higgs 
$$V=(|\mu|^2+m_{H_u}^2)|H_u^0|^2+(|\mu|^2+m_{H_d}^2)|H_d^0|^2-(b\,H_u^0H_d^0+\mathrm{c.c.})$$
 potential 
$$+\frac{1}{8}(g^2+g'^2)(|H_u^0|^2-|H_d^0|^2)^2.$$

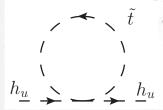
$$\tan \beta = \frac{v_u}{v_d}$$



#### In SUSY, stop loops dominate Higgs mass parameter correction

 $\delta m_H^2(stop\ loop)$ 





#### Leading quadratic divergence cancels

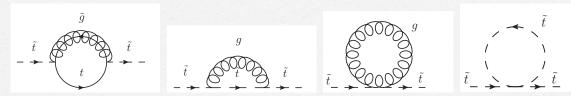
$$\delta m_{h_u}^2 = -\frac{3y_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln\left(\frac{\Lambda_{UV}}{m_{\tilde{t}}}\right)$$

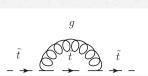
To avoid tuning need

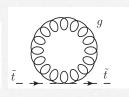
$$m_{\tilde{t}} \lesssim 400 {
m GeV}$$
. 500 GeV OK

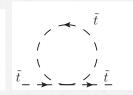
LHC should find stops SOON

## Gluino corrections to stop









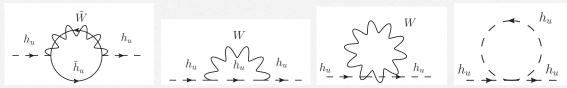
$$\delta m_{\tilde{t}}^2 = \frac{2g_s^2}{3\pi^2} m_{\tilde{g}}^2 \ln \frac{\Lambda_{UV}}{m_{\tilde{g}}}.$$

To avoid tuning need

$$m_{\tilde{g}} \lesssim 2 m_{\tilde{t}}.$$
 1 TeV OK

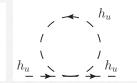
LHC should find gluinos SOON

## Other important loops









$$\delta m_{h_u}^2 = \frac{3g^2}{8\pi^2} (m_{\tilde{W}}^2 + m_{\tilde{h}}^2) \ln \frac{\Lambda_{UV}}{m_{\tilde{W}}}.$$

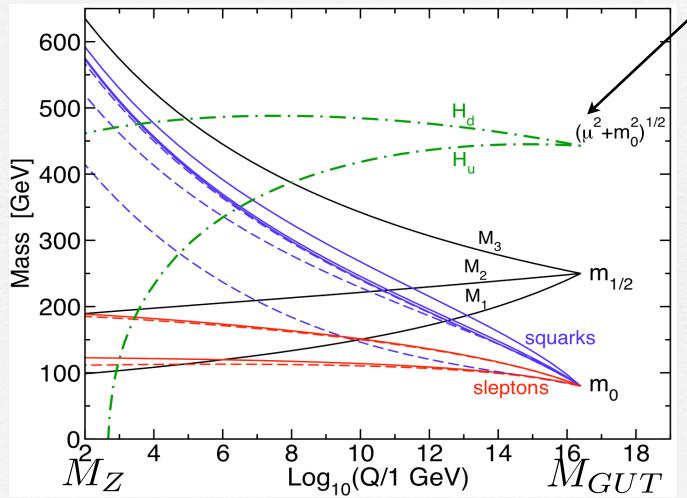
To avoid tuning need

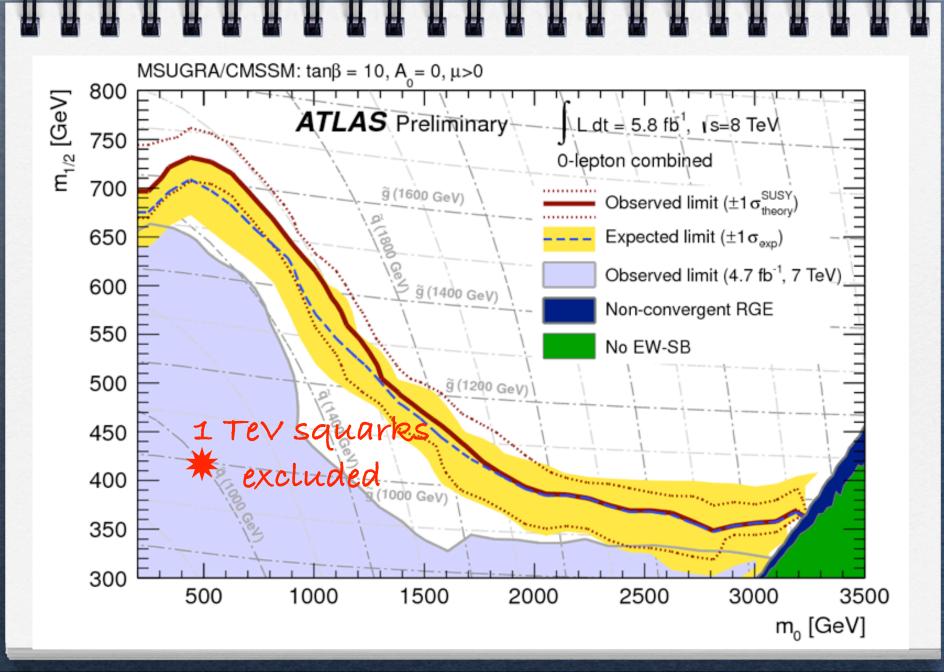
$$m_{\tilde{W}} \lesssim {
m TeV}$$
. 1 TeV OK

Difficult to find colour singlets at LHC

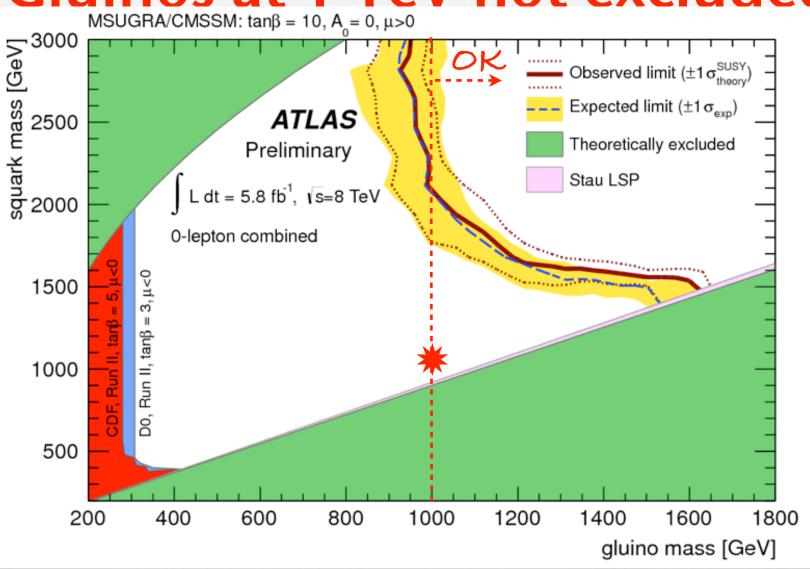
#### **Constrained MSSM**

 $m_H^2 = \mu^2 + m_0^2$ 

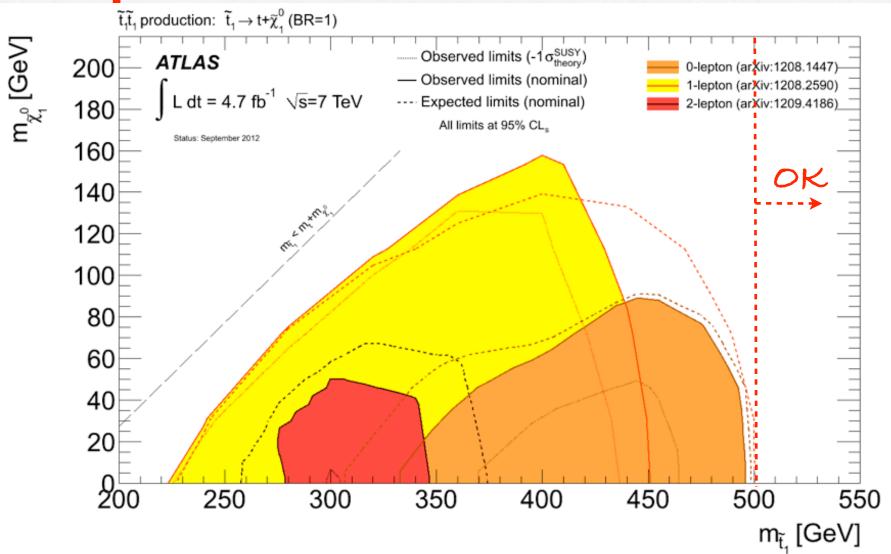




#### Gluinos at 1 TeV not excluded



#### Stops at 500 GeV not excluded



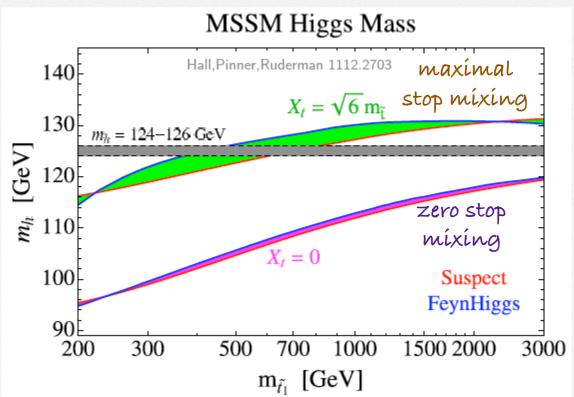
## Mass Spectrum in MSSM CHC Higgs

Names	Spin	$P_R$	Gauge Eigenstates Mass Eigensta			
Higgs bosons	0	+1	$H_u^0 \ H_d^0 \ H_u^+ \ H_d^-$	$h^0$ $H^0$ $A^0$ $H^\pm$		
			$\widetilde{u}_L \ \widetilde{u}_R \ \widetilde{d}_L \ \widetilde{d}_R$	(same)		
squarks	0	-1	$\widetilde{s}_L \ \widetilde{s}_R \ \widetilde{c}_L \ \widetilde{c}_R$	(same)		
			$\widetilde{t}_L \ \widetilde{t}_R \ \widetilde{b}_L \ \widetilde{b}_R$	$\widetilde{t}_1$ $\widetilde{t}_2$ $\widetilde{b}_1$ $\widetilde{b}_2$		
			$\widetilde{e}_L \ \widetilde{e}_R \ \widetilde{ u}_e$	(same)		
sleptons	0	-1	$\widetilde{\mu}_L \ \widetilde{\mu}_R \ \widetilde{ u}_\mu$	(same)		
			$\widetilde{ au}_L \ \widetilde{ au}_R \ \widetilde{ u}_{ au}$	$\widetilde{ au}_1 \ \widetilde{ au}_2 \ \widetilde{ u}_{ au}$		
neutralinos	1/2	-1	$\widetilde{B}^0$ $\widetilde{W}^0$ $\widetilde{H}_u^0$ $\widetilde{H}_d^0$	$\widetilde{N}_1$ $\widetilde{N}_2$ $\widetilde{N}_3$ $\widetilde{N}_4$		
charginos	1/2	-1	$\widetilde{W}^{\pm}$ $\widetilde{H}_{u}^{+}$ $\widetilde{H}_{d}^{-}$	$\widetilde{C}_1^{\pm}$ $\widetilde{C}_2^{\pm}$		
gluino	1/2	-1	$\widetilde{g}$	(same)		
goldstino (gravitino)	$\frac{1/2}{(3/2)}$	-1	$\widetilde{G}$	(same)		

### Higgs h Mass in MSSM

$$\mathcal{W} = \mu \widehat{H}_u \widehat{H}_d$$

$$m_h^2 \approx M_Z^2 \cos^2 2\beta + \Delta m_h^2$$



$$\Delta m_h^2 \approx \frac{3}{(4\pi)^2} \frac{m_t^4}{v^2} \left[ \ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left( 1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$

$$m_{\tilde{t}}^2 = m_{\tilde{Q}_3} m_{\tilde{t}_R} \qquad X_t = A_t - \mu \cot \beta.$$

Must have at least one stop mass heavier than 500 GeV -->
Fine Tuning!

#### Next-to-Minimal SUSY SM (NMSSM)

Model gives dynamical origin of  $\mu$  term via complex singlet S:

SHuHd where singlet <S>  $\sim \mu \sim$  TeV

Danger from weak scale axion due to global U(1) symmetry

Need to avoid axion somehow

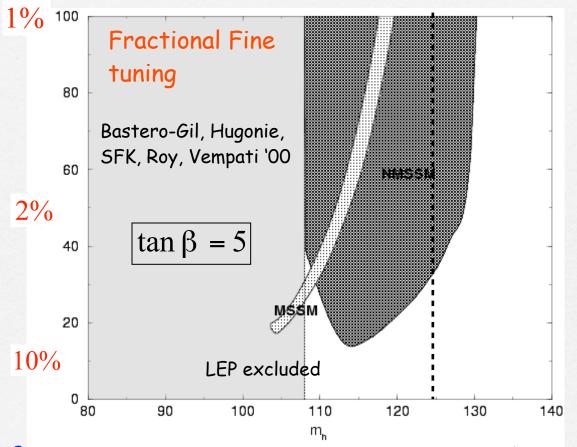
In NMSSM we add  $S^3$  to break U(1) to  $Z_3$ 

Extra tree-level contribution to

$$\mathcal{W}=\lambda \widehat{S}\widehat{H}_u\widehat{H}_d+rac{\kappa}{3}\,\widehat{S}^3$$
 Higgs mass reduces fine-tuning

$$m_h^2 \approx M_Z^2 \cos^2 2\beta + \lambda^2 v^2 \sin^2 2\beta + \Delta m_h^2$$

# Fine Tuning vs. Higgs Mass



For 125 GeV
Higgs the
MSSM fine
tuning is
much worse
than in
NMSSM

LEP favours NMSSM over MSSM (12 years ago) LHC with Higgs @ 125 GeV strengthens conclusion

## **NMSSM Higgs Theory**

Spectrum has an extra CP even S plus extra CP odd A (both singlets) compared to MSSM

CP even mass eigenstates

$$H_1 = S_{1,d} H_d + S_{1,u} H_u + S_{1,s} S,$$
  
 $H_2 = S_{2,d} H_d + S_{2,u} H_u + S_{2,s} S,$ 

$$H_3 = S_{3,d} H_d + S_{3,u} H_u + S_{3,s} S$$
.

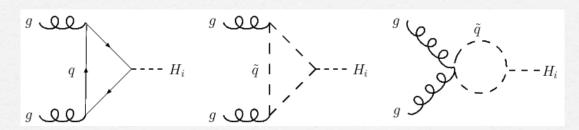
HI or H2 have reduced couplings due to the singlet component

 $h^{125\,\mathrm{GeV}}$  can be  $H_1,H_2$ 

# NMSSM Higgs Phenomenology

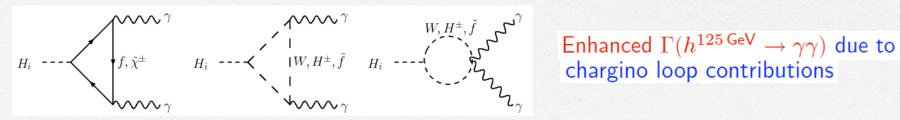
#### **Enhanced gluon fusion production**

Stop and sbottom loop contributions in  $gg o H_i$ 



$$BR(h^{125\,\mathrm{GeV}} \to \gamma\gamma) = \frac{\Gamma(h^{125\,\mathrm{GeV}} \to \gamma\gamma)}{(\Gamma_{b\bar{b}} + \Gamma_{WW} + \Gamma_{ZZ} + \ldots)[h^{125\,\mathrm{GeV}}]}$$

Suppression of  $\Gamma(h^{125\,{\rm GeV}} \to b\bar{b})$  due to strong singlet-doublet mixing



#### NMSSM Higgs Benchmarks Near 125 GeV

Point	NMP1	NMP2	NMP3		
$\tan \beta$	3	2	2		
$\mu_{\rm eff} \ [{ m GeV}]$	200	200	200		
λ	0.64	0.6	0.57		
κ	0.25	0.18	0.2		
$A_{\lambda} [{ m GeV}]$	560	405	395		
$A_{\kappa} [\mathrm{GeV}]$	-10	-10	-80		
$M_{Q_{3L}}$ [GeV]	650	700	530		
$M_{t_R}$ [GeV]	650	700	530		
$M_1$ [GeV]	106	91	115		
$M_2$ [GeV]	200	200	200		
$M_3$ [GeV]	600	600	600		
SM-like Hig	gs boson	11			
$M_{H_1}$ [GeV]	124.5	126.5	124.6		
$R_{\gamma\gamma}(H_1)$	1.06	1.24	1.47		
$R_{WW}(H_1)$	0.85	0.93	1.02		
$R_{ZZ}(H_1)$	0.76	0.85	0.90		
$R_{b\bar{b}}(H_1)$	1.12	1.09	1.04		
$R_{\Gamma_{tot}}(H_1)$	1.02	0.93	0.76		
$R_{\sigma_{gg}}(H_1)$	0.97	0.96	0.77		
$R_{\sigma_{tot}}(H_1)$	0.84	0.91	0.82		
$m_{\tilde{t}_1}$ [GeV]	548	587	358		
$m_{\tilde{t}_2} \; [{ m GeV}]$	782	838	686		
$X_t/m_{\tilde{t}}$	1.74	1.86	2.26		
Relic density					
$\Omega h^2$	0.9819	0.1170	0.1100		

Point	NMP4	NMP5	NMP6		
$\tan eta$	3	3	2		
$\mu_{\rm eff} \ [{ m GeV}]$	200	200	140		
λ	0.67	0.66	0.55		
κ	0.1	0.12	0.31		
$A_{\lambda} \; [{ m GeV}]$	650	650	210		
$A_{\kappa}$ [GeV]	-10	-10	-210		
$M_{Q_{3L}}$ [GeV]	600	600	800		
$M_{t_R}$ [GeV]	600	600	600		
$M_1$ [GeV]	200	200	145		
$M_2$ [GeV]	400	400	300		
$M_3$ [GeV]	600	600	800		
SM-like Hig	gs boson	112			
$M_{H_2}$ [GeV]	123.8	126.5	124.5		
$R_{\gamma\gamma}(H_2)$	1.09	1.19	1.431		
$R_{WW}(H_2)$	0.91	0.98	1.00		
$R_{ZZ}(H_2)$	0.80	0.89	0.89		
$R_{bar{b}}(H_2)$	1.08	1.06	1.04		
$R_{\Gamma_{tot}}(H_2)$	0.96	0.90	0.78		
$R_{\sigma_{gg}}(H_2)$	1.00	0.96	0.91		
$R_{\sigma_{tot}}(H_2)$	0.92	0.95	0.93		
$m_{\tilde{t}_1} \; [{ m GeV}]$	517	483	549		
$m_{\tilde{t}_2} \; [{ m GeV}]$	724	741	892		
$X_t/m_{\tilde{t}}$	1.56	1.89	-1.83		
Relic density					
$\Omega h^2$	0.0999	0.1352	0.1258		

Kíng, Muhlleítner, Nevzorov arxív:1201.2671

#### Key features:

- Stops below 1 TeV in all cases
- Two photon Higgs rate enhanced

$$R_{\gamma\gamma}(H_i) \equiv R_{\sigma_{incl}}(H_i) R_{\gamma\gamma}^{BR}(H_i)$$

$$R_{VV}(H_i) \equiv R_{\sigma_{incl}}(H_i)R_{VV}^{BR}(H_i)$$

$$R_{b\bar{b}}(H_i) \equiv R_{\sigma_{incl}}(H_i) R_{b\bar{b}}^{BR}(H_i)$$

## **Summary of Part 1**

- D Particle discovered at LHC consistent with SM Higgs boson
- But LHC Higgs decay signal strengths have large errors (two photon rate too high)
- ☐ Higgs may be window into BSM physics
- Higgs theory fine-tuning problem solved by SUSY
- □ Natural Susy requires stops = 500 GeV, gluino = 1 TeV (or less)
- These are not excluded by LHC searches (so far!)
- □ Stops>500 GeV required for Higgs mass in MSSM
- ☐ Stops = 500 GeV possible for Higgs mass in NMSSM
- □ NMSSM can lead to large di-photon rate especially with h(125) = H2
- Other decays such as WW, ZZ, bb, tau tau can also have different rates
- ☐ We eagerly await the next LHC results!

## Exceptional SUSY SM (E<sub>6</sub>SSM)

$$E_6 \rightarrow SO(10) \times U(1)_{\psi}$$

$$SO(10) \to SU(5) \times U(1)_{\chi}$$

M<sub>string</sub> –

M<sub>3</sub> –

M<sub>2</sub> –

M<sub>1</sub> –

 $SU(3) \times SU(2) \times U(1)_{Y} \times U(1)_{N}$ 

RH neutrinos neutral under:

$$U(1)_N = \frac{\sqrt{15}}{4}U(1)_{\psi} + \frac{1}{4}U(1)_{\chi}$$

remaining matter content of 3 families of 27's of  $E_6$  survives down to the TeV scale

 $u(1)_N$  broken, Z' and exotics get mass,  $\mu$  term generated  $su(2)_L \times u(1)_Y$  broken

Energy

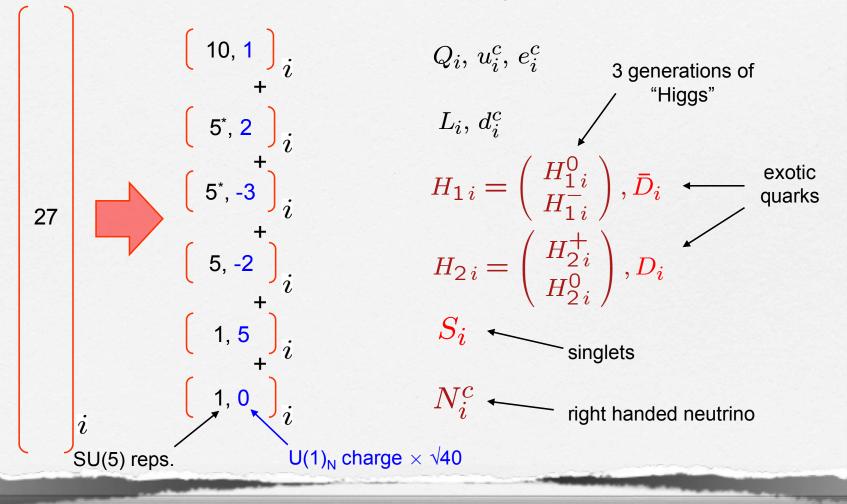
M

Tev

#### Matter Content of 27's of E<sub>6</sub>

All the SM matter fields are contained in one 27-plet of  $E_6$  per generation.

Miller



### Higgs mass bounds

 $m_{h_1}$ SFK, Moretti, Nevzorov 160  $E_6SSM$ 140120 **NMSSM** MSSM 100 80 60 2-loop Higgs mass bounds 40 20  $tan \beta$ 

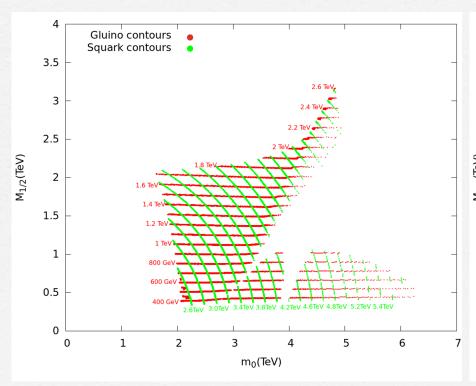
$$m_h^2 \approx \underbrace{M_Z^2 \cos^2 2\beta}_{MSSM} + \frac{\lambda^2}{2} v^2 \sin^2 2\beta + \frac{M_Z^2}{4} (1 + \frac{1}{4} \cos 2\beta)^2 + \Delta m_h^2$$

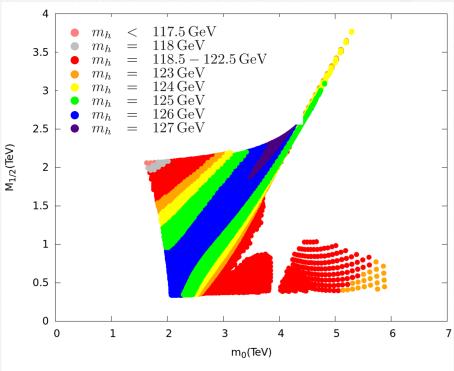
 $E_6SSM$ 

### Higgs, Squarks, Gluinos in CE6SSM

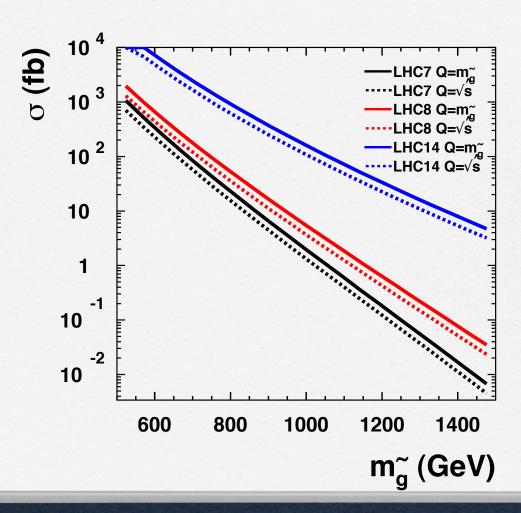
Athron, King, Miller, Moretti, Nevzorov

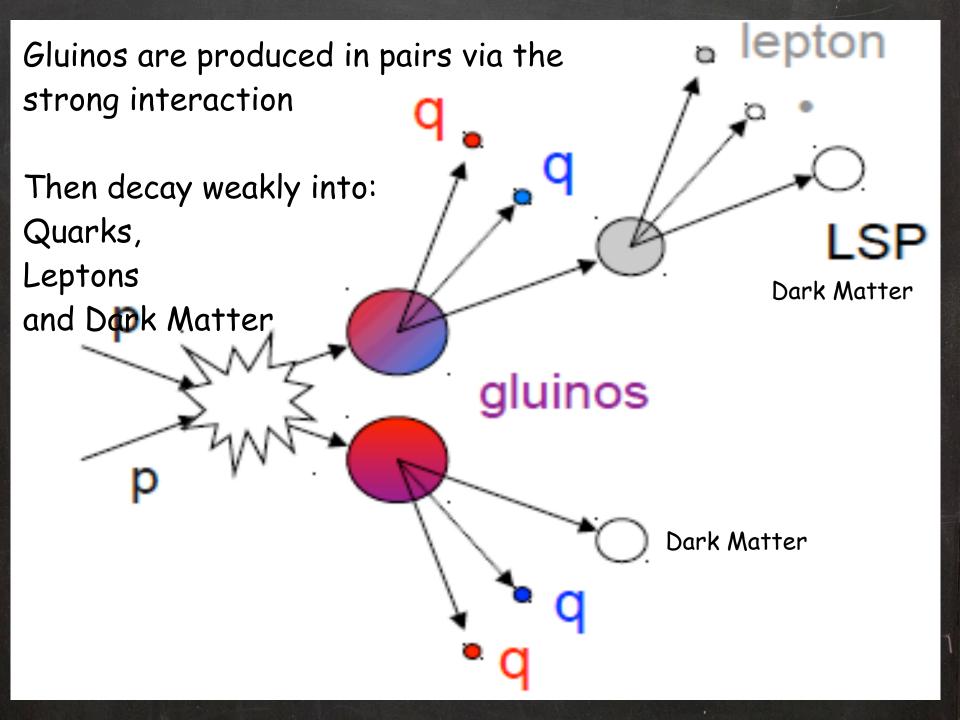
$$\tan \beta = 10, \, \lambda_{12} = 0.1, \, s = 10 \, \text{TeV}.$$





## Gluino pair production





#### Neutralinos in E<sub>6</sub>SSM

Hall, King

- 3 Higgs families = 1 MSSM family  $H_u H_d + 2$  inert families  $H_{u1} H_{d1} H_{u2} H_{d2}$
- $\square$  3 famílies of Singlets = 1 NMSSM singlet  $\Rightarrow$  + 2 inert singlets  $S_1 S_2$

The full neutralino mass matrix

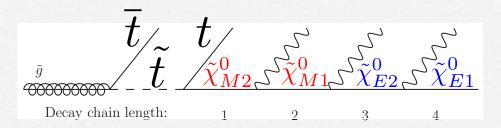
$$\tilde{\chi}_{\text{int}}^0 = (\begin{array}{ccc|c} \tilde{B} & \tilde{W}^3 & \tilde{H}_d^0 & \tilde{H}_u^0 \end{array} \middle| \begin{array}{ccc|c} \tilde{S} & \tilde{B}' & \tilde{H}_{d2}^0 & \tilde{H}_{u2}^0 & \tilde{S}_2 & \tilde{H}_{d1}^0 & \tilde{H}_{u1}^0 & \tilde{S}_1 \end{array})^{\text{T}}$$

$$M_{\rm E_6SSM}^n = \begin{pmatrix} M_{
m USSM} & D_2 & D_1 \\ B_2^{
m T} & A_{22} & A_{21} \\ B_1^{
m T} & A_{21}^{
m T} & A_{11}^{
m T} \end{pmatrix}$$

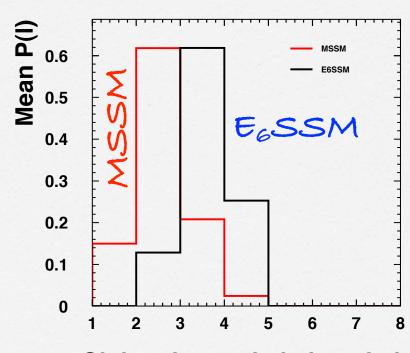
12x12 matrix!!

## Gluino decay chains

Belyaev, Hall, King, Svantesson (preliminary)



E<sub>6</sub>SSM has longer gluino chains with a lighter LSP at the end

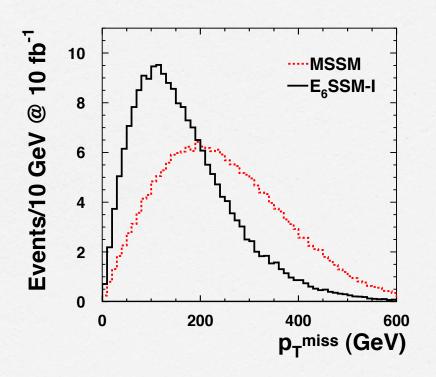


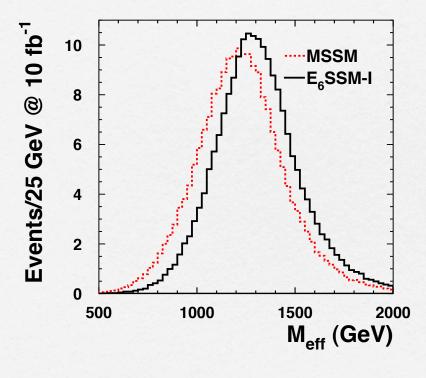
Gluino decay chain length, I

	MSSM	$E_6SSM-I$	E <sub>6</sub> SSM-II	E <sub>6</sub> SSM-III	E <sub>6</sub> SSM-IV	$E_6SSM-V$	E <sub>6</sub> SSM-VI	
$\tan \beta$	10	1.5	1.42	1.77	3	1.42	1.42	
λ		0.497	0.598	-0.462	-0.4	0.598	0.598	
s		5180	5268	5418	5500	5268	5268	
$\mu$	1578	(1820)	(2228)	(1770)	(-1556)	(2228)	(2228)	
$A_t = A_b = A_t$		-3110	-3100	476.2	4638	-2684	-2684	
$M_A$	302.5	3666	4365	2074	4341	4010	4000	$\overline{Q}$
$M_1$	150	150	150	150	150	150	150	[GeV]
$M_2$	285	300	300	300 151	300	300	300	
$M_{1'}$	800.2	151 800.0	151 800.0	800.0	151 800.0	151 800.0	151 800.0	
$m_{\tilde{g}}$	148.7	148.9	149.1	151.2	150.6	149.1	149.1	
$m_{ ilde{\chi}_{M1}^0}$	302.2	296.1	296.8	303.7	301.7	296.8	296.8	
$m_{ ilde{\chi}^0_{M2}}$								
$m_{ ilde{\chi}^0_{M3}}$	1582	1763	2233	1766	1557	2233	2233	
$m_{ ilde{\chi}_{M4}^0}$	1584	1823	2246	1771	1558	2246	2246	
$m_{ ilde{\chi}_{M1}^{\pm}}$	302.2	299.0	299.2	300.9	300.4	299.2	299.2	
$m_{ ilde{\chi}_{M2}^{\pm}}$	1584	1822	2229	1771	1557	2229	2229	
$m_{ ilde{\chi}_{U1}^0}$		1878	1835	1909	1937	1835	1835	$\overline{Q}$
$m_{ ilde{\chi}_{U2}^0}$		1973	2003	2062	2087	2003	2003	[GeV]
$m_{ ilde{\chi}_{E1}^0}$		62.7	43.5	45.2	0	0	0.00011	
$m_{ ilde{\chi}^0_{E2}}$		62.8	48.6	53.2	0	0	1.53	
$m_{ ilde{\chi}_{E3}^0}$		119.8	131.3	141.6	164.1	119.9	120.1	
$m_{ ilde{\chi}_{E4}^0}$		121.0	163.6	187.4	164.1	119.9	122.8	
$m_{ ilde{\chi}_{E5}^0}$	-	183.0	197.0	227.8	388.9	185.8	185.8	
$m_{ ilde{\chi}^0_{E6}}$		184.4	224.3	265.6	388.9	185.8	187.0	
$m_{ ilde{\chi}_{E1}^{\pm}}$		109.8	119.9	122.7	164.1	119.9	119.9	
$m_{ ilde{\chi}_{E2}^{\pm}}$		117.7	185.8	225.1	388.9	185.8	185.8	
$m_h^{\kappa_{E2}}$	124.4	125.4	133.8	116.3	124.7	126.1	125.8	
P(l=1)	0.188	$< 10^{-9}$	$< 10^{-5}$	$< 10^{-5}$	0.1727	$< 10^{-8}$	$< 10^{-12}$	
P(l=2)	0.812	$< 10^{-4}$	0.01524	0.1723	0.8273	0.01	$< 10^{-5}$	
P(l=3)	0	0.1746	0.2336	0.7986	$< 10^{-6}$	0.2	0.1721	
P(l=4)	0	0.8196	0.7512	0.02915	$< 10^{-15}$	0.8	0.8280	
P(l=5)	0	0.0058	$< 10^{-7}$	0	0	$< 10^{-15}$	0	
$\Omega h^2$	0.00628	0.00114	0.0006842	0.0006937	0.101	0.00154		
$\sigma_{SI}$	$0.401 \times 10^{-9}$	$15.34 \times 10^{-8}$	$9.35 \times 10^{-8}$	$16.35 \times 10^{-8}$	$3.75 \times 10^{-11}$	$3.98 \times 10^{-13}$		[pb]

## E<sub>6</sub>SSM gluino gives less p<sub>T</sub>miss

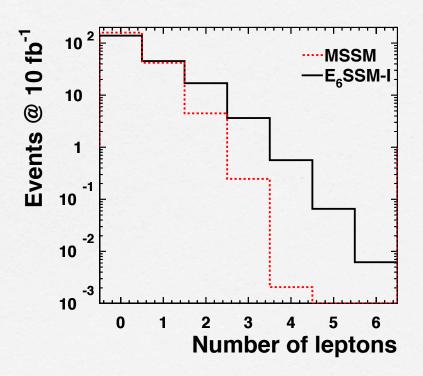
Belyaev, Hall, King, Svantesson (preliminary)

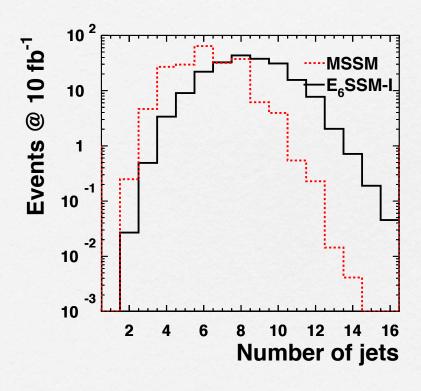




## E<sub>6</sub>SSM gluino gives more jets and leptons

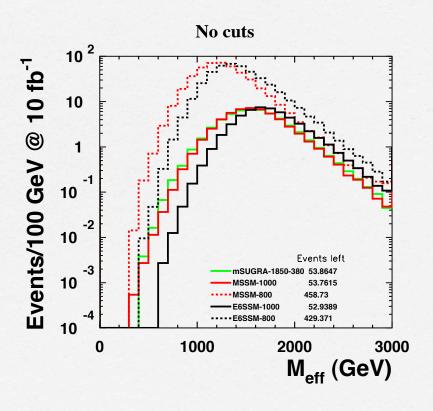
Belyaev, Hall, King, Svantesson (preliminary)

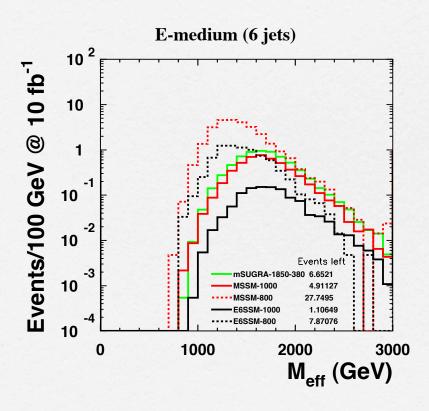




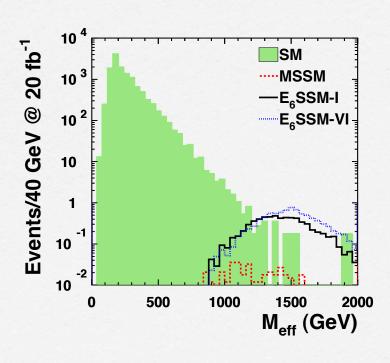
LHC@7 TeV, gluino@800 GeV

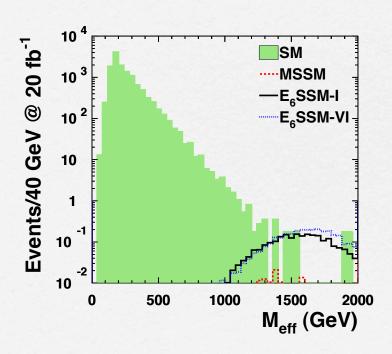
# E<sub>6</sub>SSM gluino harder to see in 6 jet channel





# E<sub>6</sub>SSM gluino easier to see in 3 lepton channel at 8 TeV

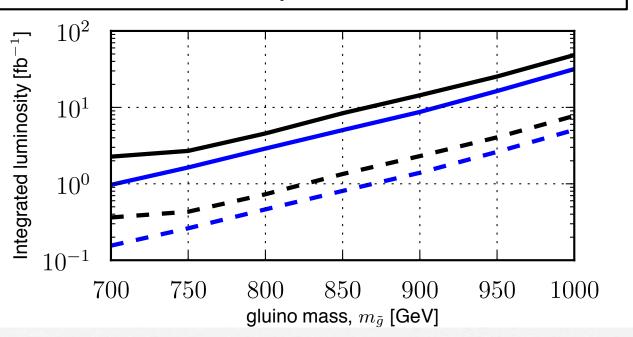




## Luminosity required for $5\sigma$ discovery and $2\sigma$ exclusion, $\sqrt{s}=8{\rm TeV}$ in 3 lepton channel

E<sub>6</sub>SSM-I:  $5\sigma$  discovery **–** • E<sub>6</sub>SSM-I:  $2\sigma$  exclusion

 $\mathsf{E}_6\mathsf{SSM}\text{-VI: }5\sigma$  discovery  $\bullet$   $\bullet$   $\mathsf{E}_6\mathsf{SSM}\text{-VI: }2\sigma$  exclusion



## **Summary of Part 2**

- □ E6SSM is richer theory with many LHC signals
- □ Matter spectrum with 3 families of 27 dimensional particle reps and a Z'
- ☐ Higgs at 125 GeV possible in E6SSM
- □ Typical spectrum is heavy squarks but lighter gluino
- Gluino has longer decay chains with more jets and leptons and less missing transverse momentum
- □ E6SSM Gluino is harder to see in 6 jet channel
- ☐ But easier to see in 3 lepton channel
- We eagerly await the next LHC results!

#### **Extra Slides**

#### F-Theory GUTs: a 12d string theory

Heckman and Vafa "6d spheres" with "2d fibres" "4d Flatlander" SU(2) SU(3) gluons Quarks Unification We live close to E<sub>8</sub> point SU(5)<sub>GUT</sub> on the extra dimensional "6d sphere" Steve King, Sout 3/24/11

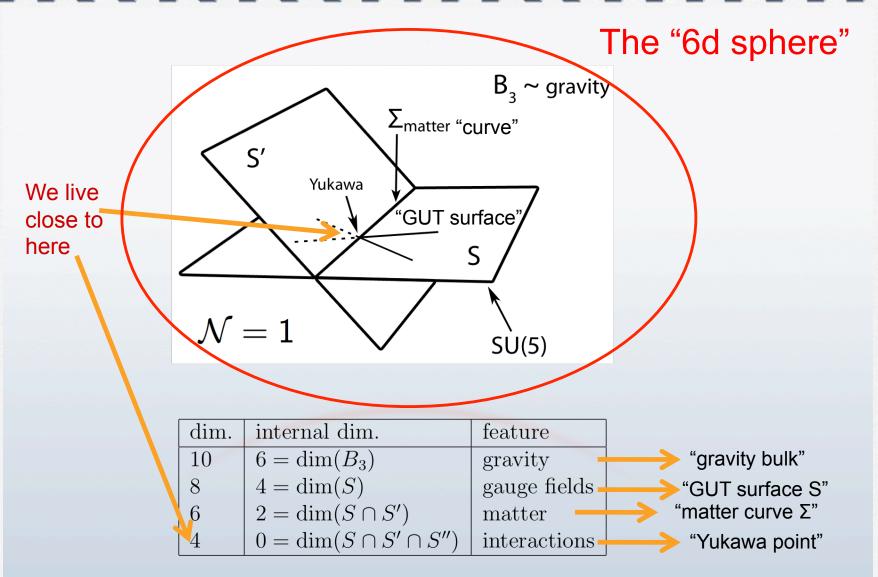


Figure 1: The structure of an F-theory GUT

## GUT breaking is achieved not with Higgs but with Hypercharge Flux

$$SU(5)\supset SU(3)_C\times SU(2)_L\times U(1)_Y$$
 Think of magnetic 
$$5\to (1,2)_{1/2}+(3,1)_{-1/3}.$$
 2-d Matter curve  $\Sigma$ 

Index theorem gives number of chiral doublets and triplets (think of Gauss's law):

$$(1,2)_{1/2}: n_L - n_R = 3 \int_{\Sigma} F_{U(1)_Y} + q \int_{\Sigma} F_{U(1)_{\perp}}$$
$$(3,1)_{-1/3}: n_L - n_R = -2 \int_{\Sigma} F_{U(1)_Y} + q \int_{\Sigma} F_{U(1)_{\perp}}$$

Doublet-triplet Higgs splitting requires:

Higgs: 
$$\int_{\Sigma} F_{U(1)_Y} \neq 0$$
Matter: 
$$\int_{\Sigma} F_{U(1)_Y} = 0.$$

Typically predicts exotics

Callaghan, King, Leontaris, Ross

### E6SSM from F-theory

$E_6$	<i>SO</i> (10)	SU(5)	Weight vector	$N_Y$	$M_{U(1)}$	SM particle content	Low energy spectrum	
$27_{t'_1}$	16	$\overline{5}_3$	$t_1 + t_5$	1	4	$4d^c + 5L$	$3d^c + 3L$	
$27_{t'_1}$	16	10 <sub>M</sub>	$t_1$	-1	4	$4Q + 5u^c + 3e^c$	$3Q + 3u^c + 3e^c$	
$27_{t_1'}$	16	$\theta_{15}$	$t_1 - t_5$	0	$n_{15}$	$3v^c$	-	
$27_{t_1'}$	10	51	$-t_1 - t_3$	-1	3	$3D+2H_u$	$3D+2H_u$	
$27_{t_1'}$	10	$\overline{5}_2$	$t_1 + t_4$	1	3	$3\overline{D} + 4H_d$	$3\overline{D} + 3H_d$	
$27_{t_1'}$	1	$\theta_{14}$	$t_1 - t_4$	0	$n_{14}$	$ heta_{14}$	<u>-</u>	
$27_{t_3'}$	16	$\overline{5}_5$	$t_3 + t_5$	-1	-1	$\overline{d^c} + 2\overline{L}$	-	
$27_{t_{3}'}$	16	102	$t_3$	1	-1	$\overline{Q} + 2\bar{u^c}$	-	
$27_{t_{3}'}$	16	$\theta_{35}$	$t_3 - t_5$	0	n <sub>35</sub>	<u>-</u>	-	
$27_{t_3'}$	10	$5_{H_u}$	$-2t_{1}$	1	0	$H_u$	$H_u$	
$27_{t_{3}'}$	10	$\overline{5}_4$	$t_3 + t_4$	-1	0	$\overline{H_d}$	-	
		-		-				

F-theory model predicts incomplete multiplets with matter content of 3 copies of 27s of E6

#### **CMSSM Dark Matter**

Neutralino mass matrix

$$ilde{B}$$
  $ilde{W}_3$   $ilde{H}_d$   $ilde{H}_u$ 

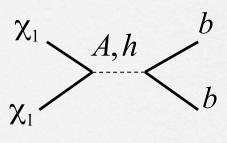
$$egin{pmatrix} M_1 & & & & & \\ & M_2 & & & & \\ & & 0 & -\mu & \\ & & -\mu & 0 \end{pmatrix}$$

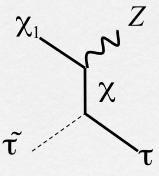
$$\chi_1 = N_1 \tilde{B} + N_2 \tilde{W} + N_3 \tilde{H}_d + N_4 \tilde{H}_u$$

$$\Omega_{DM}h^2 = C \frac{T_0^3}{M_P^2} \frac{1}{\langle \sigma v \rangle}$$

$$\chi_1$$
 $\tilde{f}$ 
 $\chi_1$ 
 $f$ 

 $\chi_1$   $\chi^{\pm}$   $\chi_{W}$ 





Bulk

$$m_{\tilde{f}} \approx m_{\chi_1}$$

Focus

Funnel

$$m_{A,h} \approx 2m_{\chi_1}$$

Co-annihilation

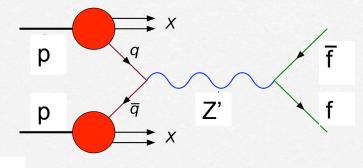
$$m_{\tilde{\tau}} \approx m_{\chi_1}$$

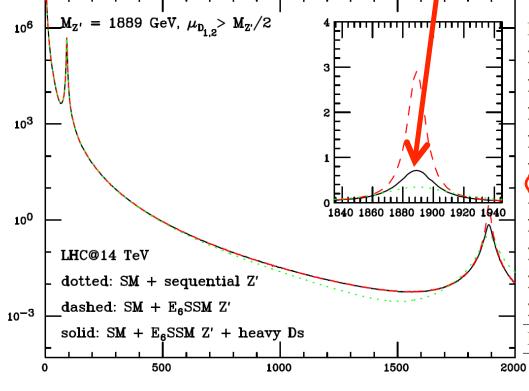
Athron, SFK, Miller, Moretti, Nevzorov

 $Z'_N$ 

 $d\sigma/dM_{l+l-}$  (fb/GeV)

Latest ATLAS limit is  $M_{Z_N^\prime} > 1520~GeV$  but exotic decays makes Z' peak smaller and harder to discover





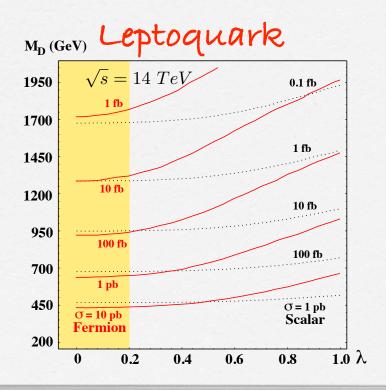
 $M_{l+l-}$  (GeV)

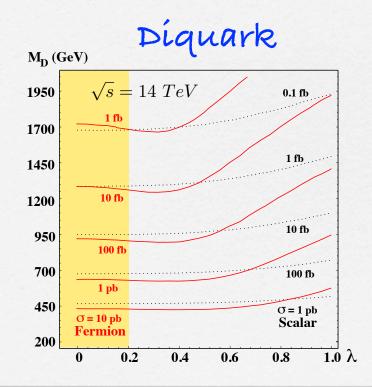
$\Gamma(Z_N' \to l^+ l^-) \ (l = e, \mu \text{ or } \tau)$	0.77
$\Sigma_l \Gamma(Z_N' \to \nu_l \overline{\nu}_l)$ (all neutrinos)	1.64
$\Sigma_l \Gamma(Z_N' \to l^+ l^-, \nu_l \overline{\nu}_l)$ (all leptons)	3.96
$\Sigma_q \Gamma(Z_N' \to q\bar{q})$ (all quarks)	10.08
$\Sigma_i \Gamma(Z_N' \to D_i \bar{D}_i)$ (exotic fermions)	0.00
$\sum_{\alpha} \Gamma(Z'_N \to \tilde{H}_{\alpha} \tilde{H}_{\alpha})$ (inert Higgsinos	5.19
$\Sigma_{\alpha}\Gamma(Z_N' \to \tilde{S}_{\alpha}\tilde{S}_{\alpha}) \text{ (singlinos)}$	7.63
$\Sigma_i \Gamma(Z_N' \to D_i D_i)$ (exotic scalars)	0.19
$\Sigma_f \Gamma(Z_N' \to \tilde{f}\tilde{f})$ (sfermions)	0.010
$\Sigma_{\alpha}\Gamma(Z_N' \to H_{\alpha}H_{\alpha})$ (inert Higgses)	0.39
$\Sigma_j \Gamma(Z_N' \to \tilde{\chi}_j \tilde{\chi}_j)$ (gauginos)	$7.92 \times 10^{-5}$
$\Gamma_{\text{tot}}$ (all)	27.45

### **Exotic D-particles**

Kang, Langacker, Nelson

D-particles are coloured and may be pair produced at LHC D-particles may be Leptoquarks D→LQ or Diquarks D→QQ



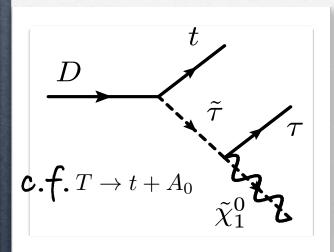


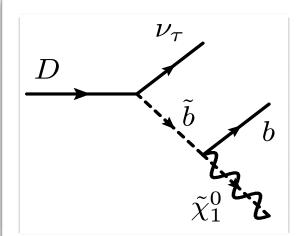
### **D-fermion decays**

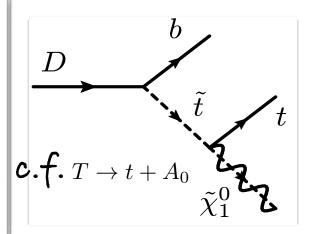
Leptoquark

Leptoquark

Díquark







$$pp \to t\bar{t}\tau^+\tau^- + E_T^{miss} + X$$

$$pp \to b\bar{b} + E_T^{miss} + X$$

$$pp \to t\bar{t}b\bar{b} + E_T^{miss} + X$$

spectacular signals!

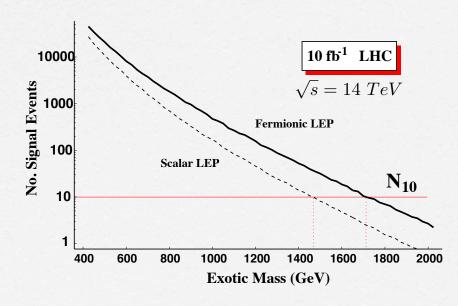
Kang, Langacker, Nelson

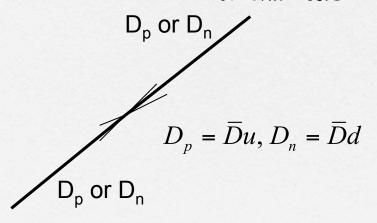
#### D-fermion as R-hadron

- Imposing B and Lall couplings DFF forbidden
- D-particle are quasi-stable R-hadrons, decay via

dim5 :  $D^cQH_dS$ ,  $D^cQQu^c$ ,  $D^cQL\nu^c$ .

punch through to muon chambers





Belyaev, Hall, King, Svantesson (preliminary)

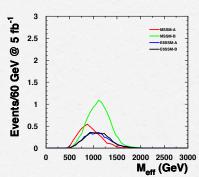
## Meff

$$M_{eff} = p_T^{miss} + \sum_{jets} |p_T^{jet}|$$

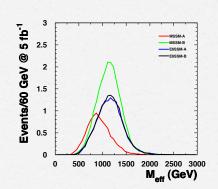
#### Meff does not really help to distinguish models

#### ATLAS style cuts

	MSSM-A								
No.	limit	Eff.	${\rm Frac.}$	Eff.	Frac.	Eff.	Frac.	Eff.	Frac
0	no cut	0.00	1.00	0.00	1.00	0.00	1.00	0.00	1.00
1	$p_T > 130$	0.12	0.88	0.19	0.81	0.40	0.60	0.40	0.60
2	$\begin{array}{c} p\!\!\!/_{T} > \!\!\!130 \\ p_{T}^{jet_{1}} > \!\!\!130 \end{array}$	0.42	0.51	0.04	0.77	0.03	0.58	0.03	0.59
3	$p_T^{jet_2} > 40$	0.13	0.44	0.01	0.76	0.00	0.58	0.00	0.58
4	$p_T^{jet_3} > 40$	0.36	0.28	0.11	0.68	0.04	0.56	0.04	0.56
5	$p_T^{jet_4} > 40$	0.55	0.13	0.20	0.54	0.11	0.50	0.11	0.50
$6 \Delta \phi$	$b(p_T, jet)_{min} > 0.4$	0.28	0.09	0.37	0.34	0.59	0.20	0.58	0.21
7	$p_T/M_{eff} > 0.25$	0.15	0.08	0.49	0.17	0.69	0.06	0.68	0.07



	0 500 1000 1500 2000 2500 3000 M <sub>eff</sub> ( <b>GeV</b> )
	0
Eve	1 / / \
inte	
9/8	2   // \
Events/60 GeV @	3 / //
>	4 / // \
(8)	5 — MSSM-B — E6SSM-A — E6SSM-B
5 fb	——MSSM-A ——MSSM-B
7	



#### CMS style cuts

CUT			MSSM-A		MSSM-B		E <sub>6</sub> SSM-A		E <sub>6</sub> SSM-B	
No.	limit	Eff.	Frac.	Eff.	Frac.	Eff.	Frac.	Eff.	Frac.	
0	no cut	0.00	1.00	0.00	1.00	0.00	1.00	0.00	1.00	
1	$H_T > 200 \text{ GeV}$	0.58	0.42	0.34	0.66	0.47	0.53	0.47	0.53	
2	$p_T^{jet_1} > 50 \text{ GeV}$	0.00	0.42	0.00	0.66	0.00	0.53	0.00	0.53	
3	$p_T^{jet_2} > 50 \text{ GeV}$	0.13	0.37	0.02	0.64	0.01	0.52	0.01	0.53	
4	$p_T^{jet_3} > 50 \text{ GeV}$	0.43	0.21	0.13	0.56	0.06	0.49	0.06	0.50	
5	$\Delta \phi(p_T, jet_1) > 0.5$	0.02	0.21	0.02	0.55	0.03	0.48	0.03	0.48	
6	$\Delta \phi(p_T, jet_2) > 0.5$	0.05	0.19	0.08	0.50	0.12	0.42	0.12	0.42	
7	$\Delta \phi(p_T, jet_3) > 0.3$	0.04	0.19	0.07	0.47	0.10	0.38	0.10	0.38	
8	$\Delta R(jet, lep)_{min} < 0.3$	0.18	0.15	0.24	0.36	0.37	0.24	0.36	0.25	
9	$H_T > 800 \text{ GeV}$	0.88	0.02	0.49	0.18	0.38	0.15	0.38	0.15	

#### Two potential problems: rapid proton decay + FCNCs

- FCNC problem may be tamed by introducing a  $Z_2^H$  under which third family Higgs and singlet are even all else odd  $\rightarrow$  only allows Yukawa couplings involving third family Higgs and singlet  $H_u$ ,  $H_d$ , S
- Z<sub>2</sub><sup>H</sup> also forbids all DFF and hence forbids D decay (and p decay)
- $\rightarrow$   $Z_2^H$  cannot be an exact symmetry!

How do we reconcile D decay with p decay?

In E<sub>6</sub>SSM can have extra discrete symmetries:

 $Z_2^L$  under which L are odd  $\rightarrow$  forbids DQL, allows DQQ  $\rightarrow$  exotic D are diquarks

 $Z_2^B$  with L & D odd  $\rightarrow$  forbids DQQ, allows DQL  $\rightarrow$  exotic D are leptoquarks

**Or:--** small DFF couplings  $\sim 10^{\text{-}12}$  will suppress p decay sufficiently while couplings  $\sim 10^{\text{-}12}$  will allow D decay with lifetime <0.1 s (nucleosynth) N.B.  $\Gamma_{\text{D}} \propto g^{\text{2}}$ ,  $\Gamma_{\text{p}} \propto g^{\text{4}}$  (Howl, SFK)