

# CP Violation (and Baryogenesis) in Two Higgs Doublet Models

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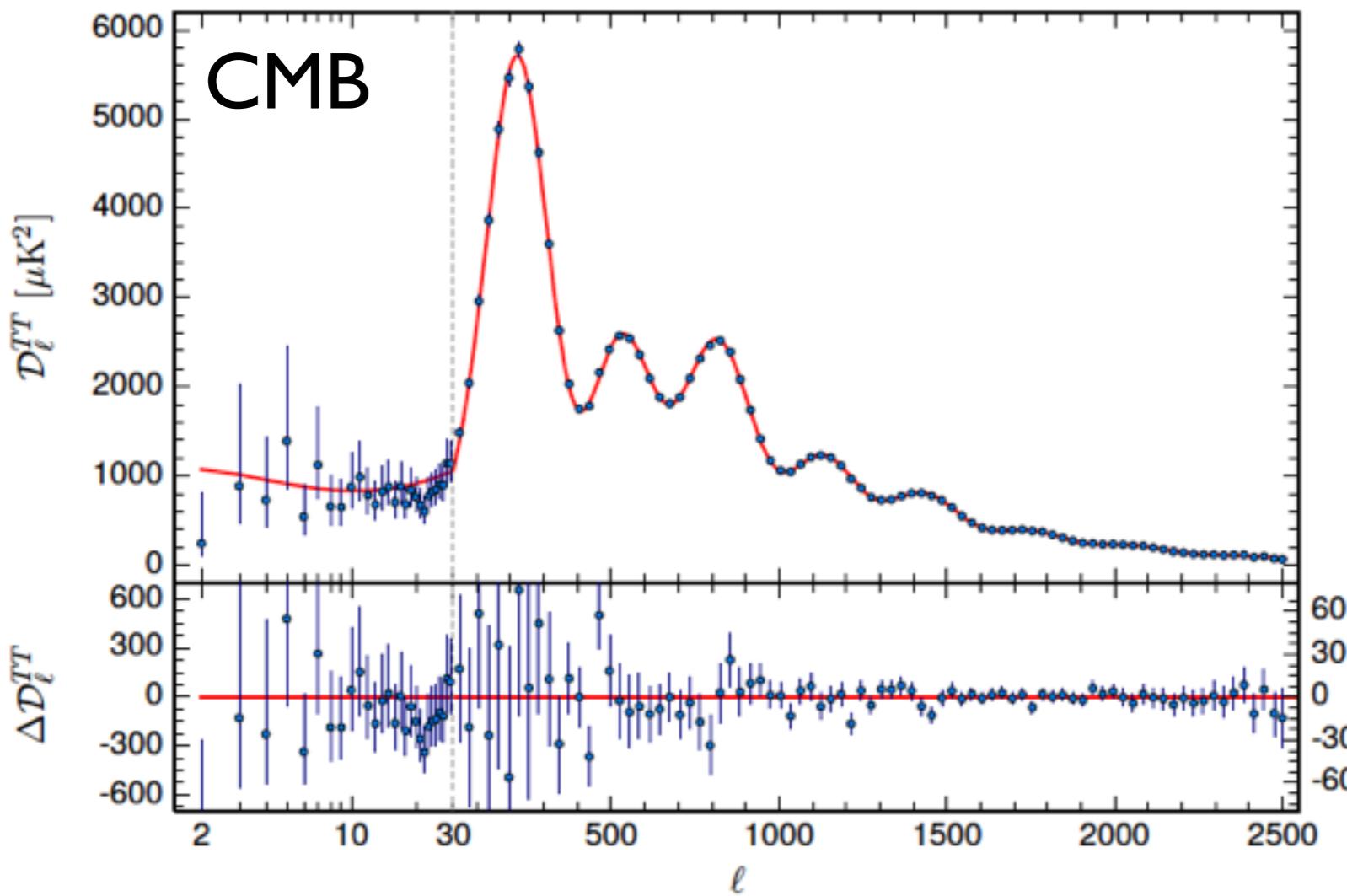
University of Sussex  
May 28, 2015

Ipek, PRD D89 (2014) 073012, arXiv:1310.6790 + work in progress

# Sneak peek

- There is more matter than antimatter - *baryogenesis*
- SM cannot explain this:
  - There is baryon number violation
  - Not enough CP violation
  - No out-of-equilibrium processes
- Solution: Two Higgs Doublet Model (2HDM)
- Constraints on CP violation: EDM experiments
- Small CP violation gives an “easy”, perturbative approach
- Study the phase structure of CP violating 2HDM?

# There is more matter than antimatter



$$\Omega_\Lambda \sim 0.69$$

$$\Omega_{\text{DM}} \sim 0.27$$

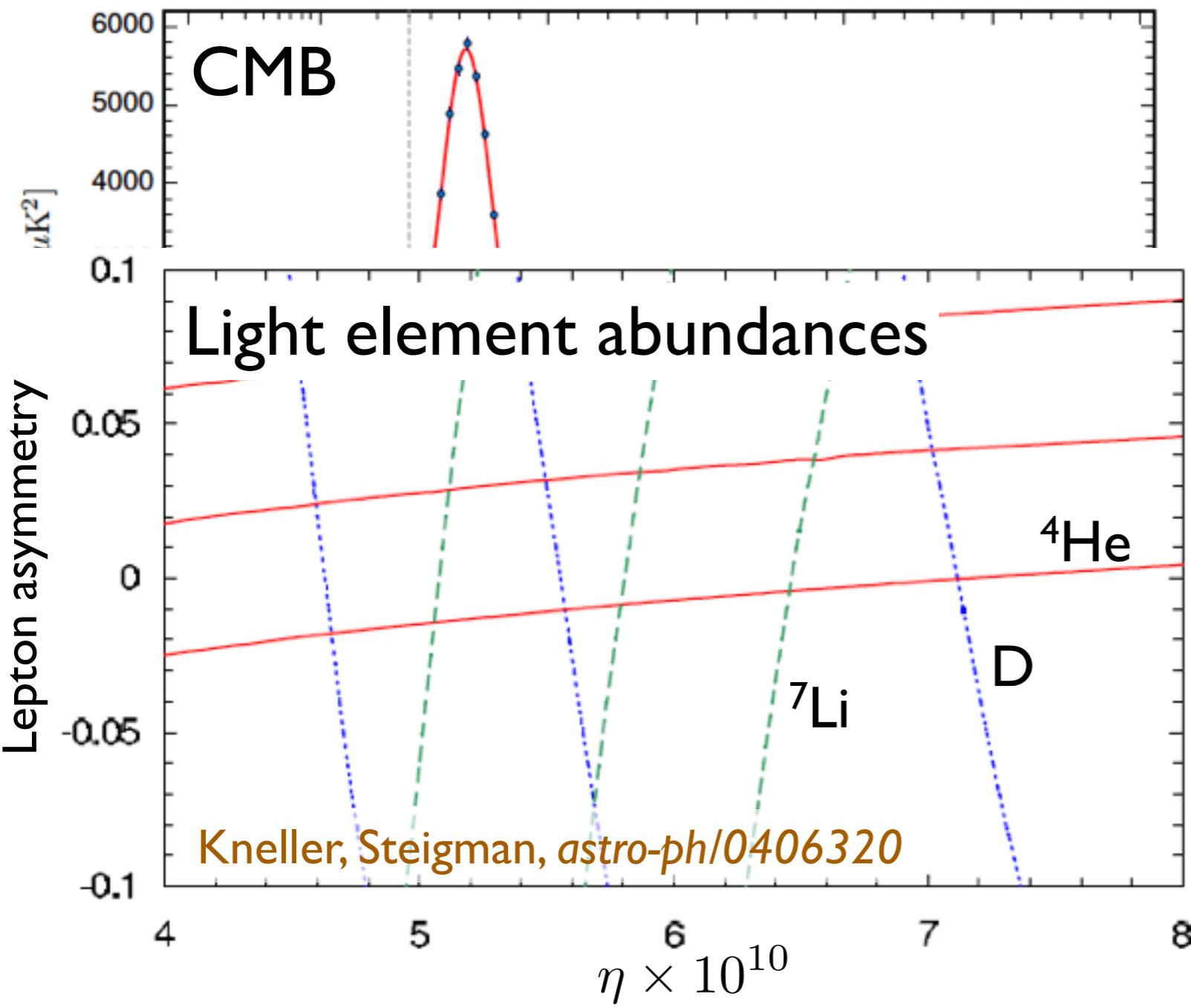
$$\Omega_B \sim 0.04$$

number of baryons:

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma}$$

$$\simeq 6 \times 10^{-10}$$

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$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma}$$

$$\approx 6 \times 10^{-10}$$

# Sakharov Conditions

Sakharov, *JETP Lett.* 5, 24 (1967)

Need to produce 1 extra quark for every 10 billion antiquarks!

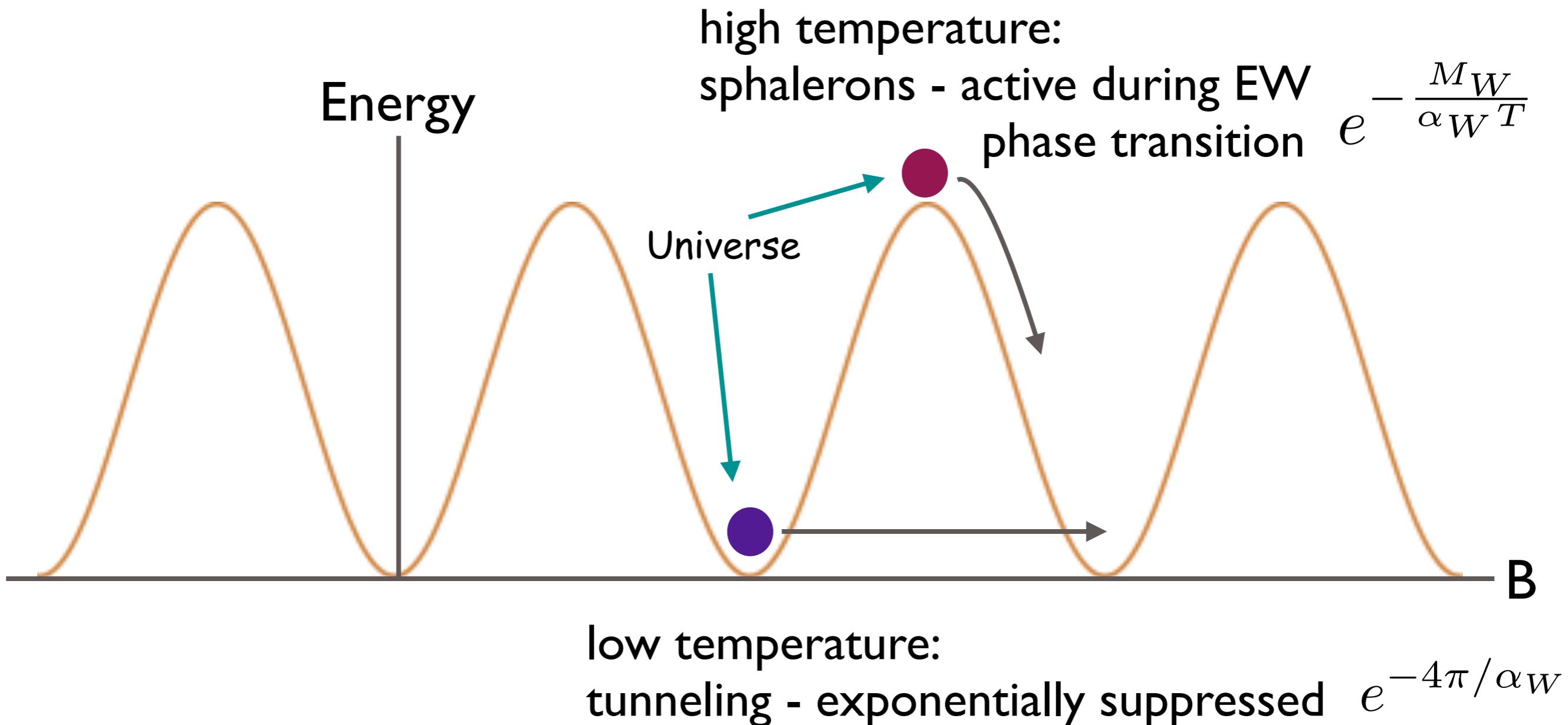
But how?

Three conditions must be satisfied:

- 1) Baryon number (B) must be violated  
can't have a baryon asymmetry w/o violating baryon number!
- 2) C and CP must be violated  
a way to differentiate matter from antimatter
- 3) B and CP violating processes must happen out of equilibrium  
equilibrium destroys the produced baryon number

# I) SM violates baryon number ✓

Baryon number is anomalously violated

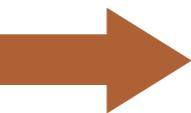


# 2) SM has CP violation

But not enough!

Only source of CP violation: quark mixing matrix:

$$-\frac{g}{\sqrt{2}} (\bar{u} \ \bar{c} \ \bar{t}) \gamma^\mu W_\mu^+ \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{\text{CKM}}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

3 mixing angles + 1 phase  CP violation!

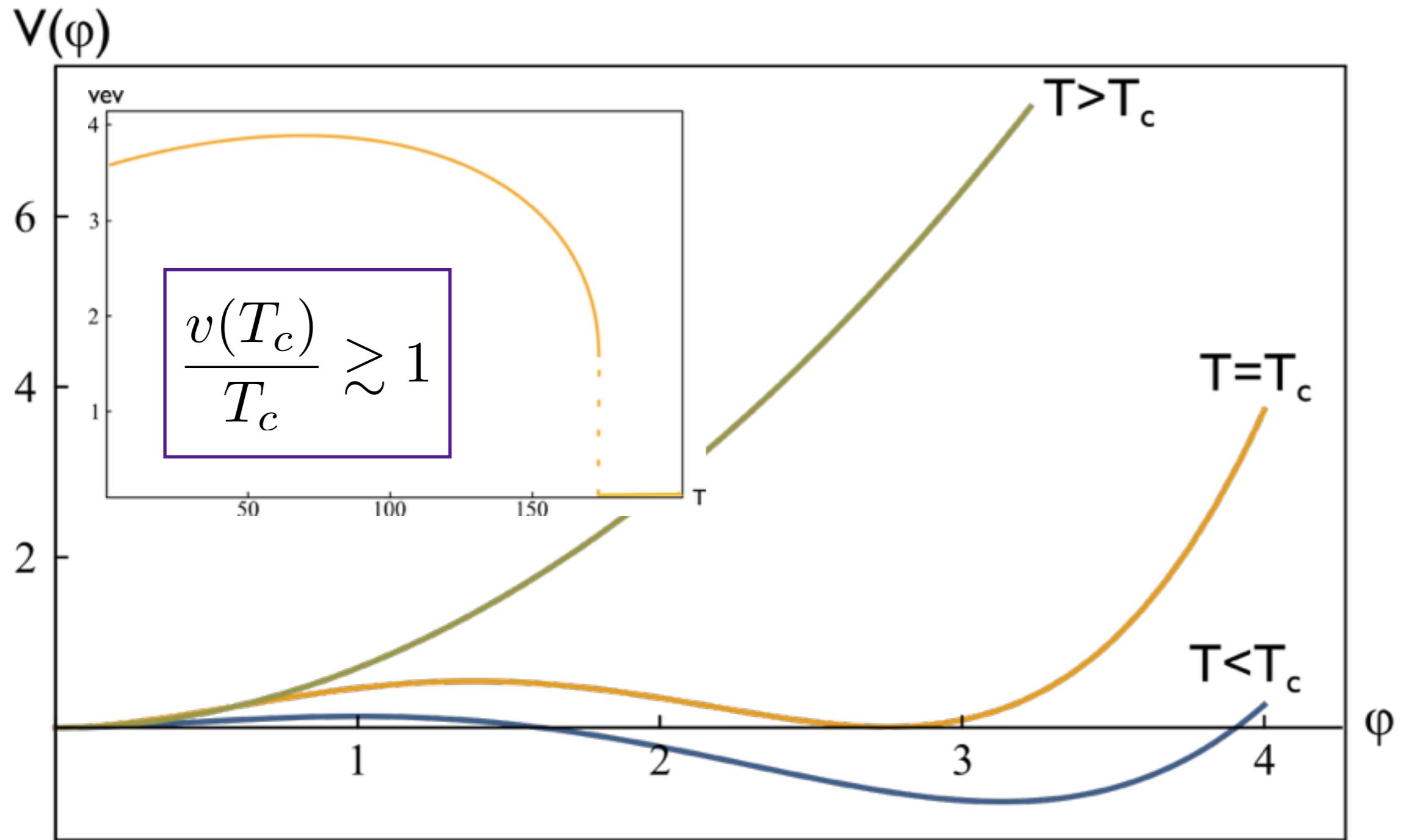
BUT

too small + gets suppressed by small Yukawa's

  $\eta \sim 10^{-20}$  

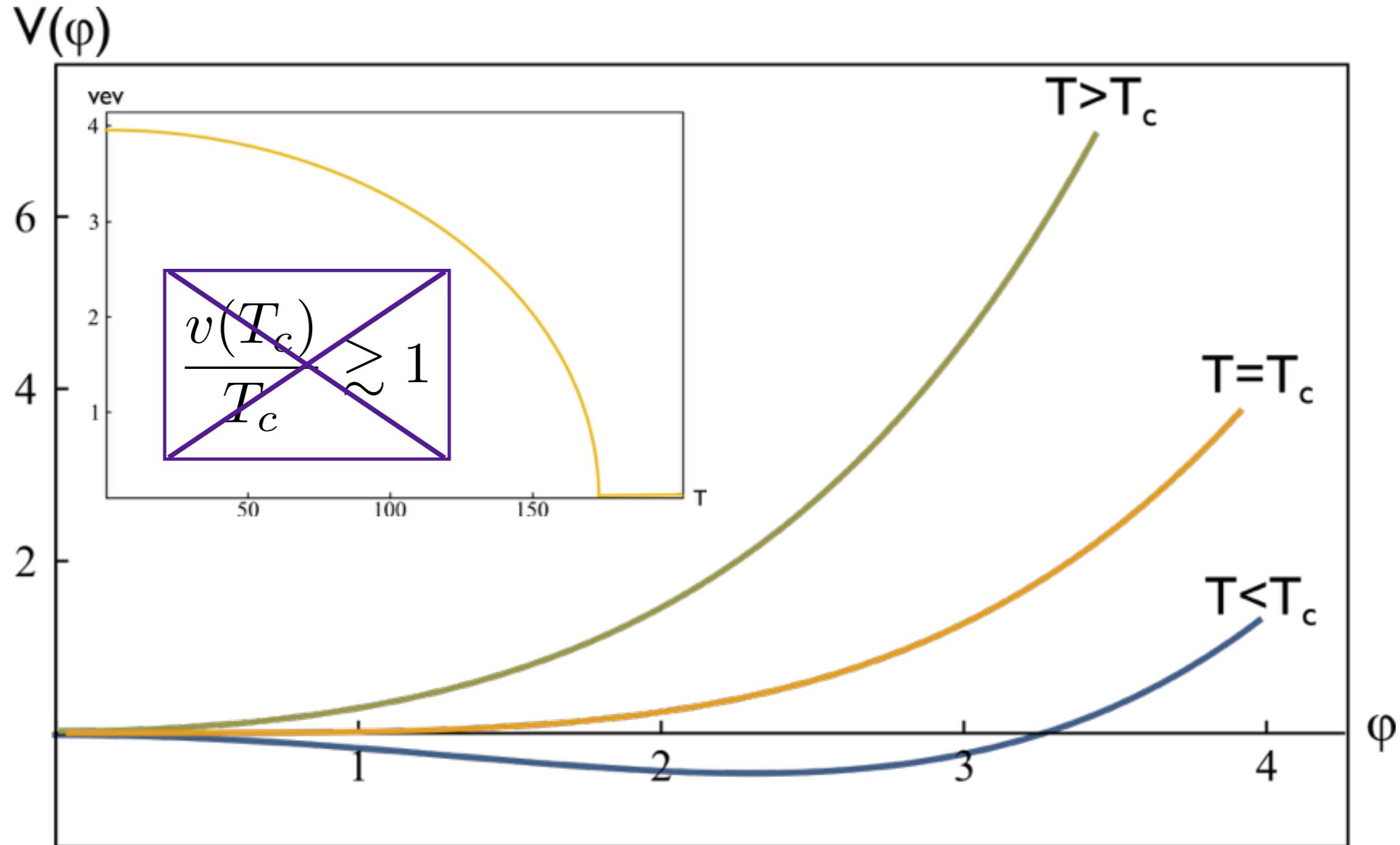
# 3) SM has EW phase transition X

But not out-of-equilibrium!



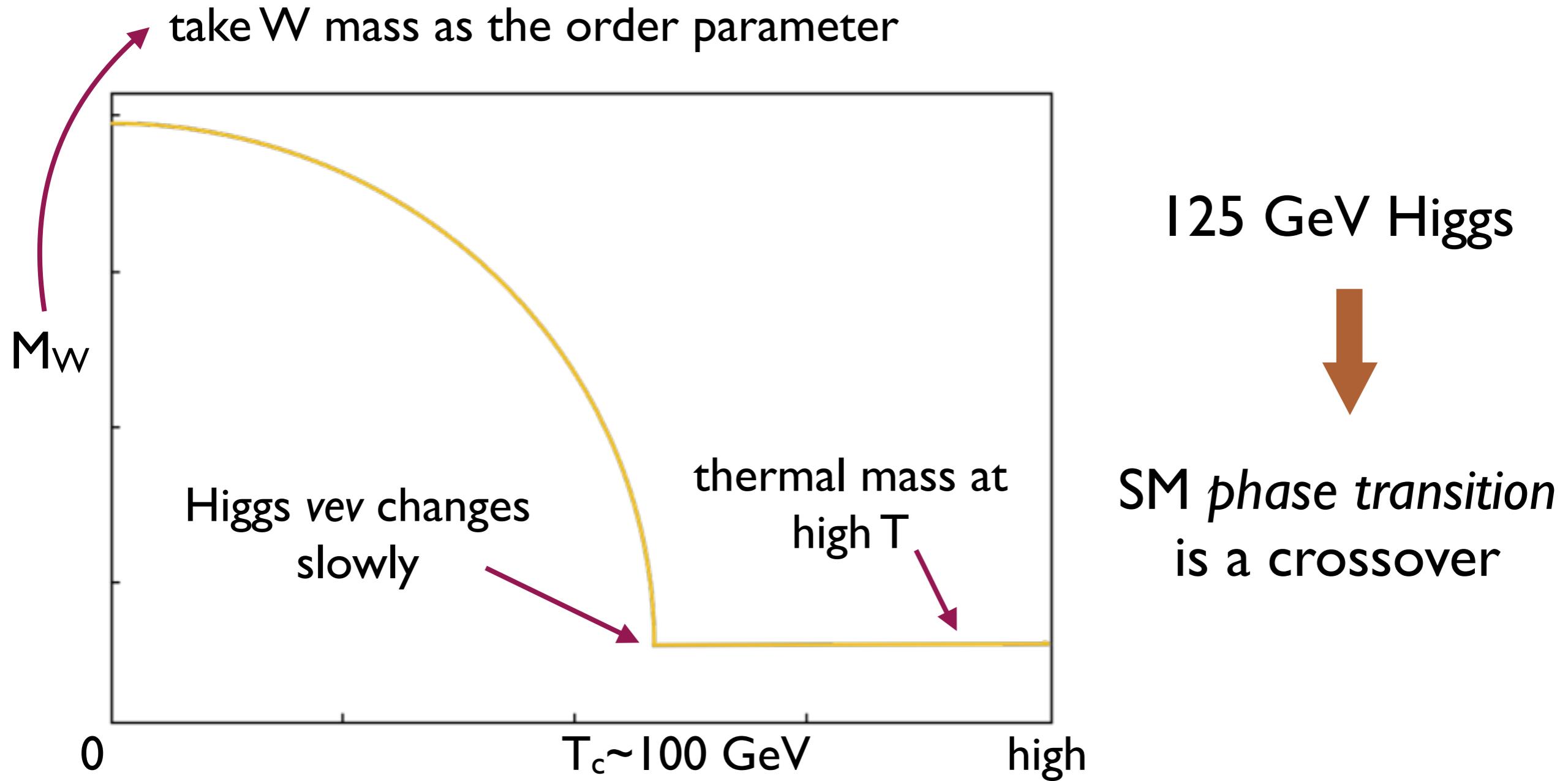
# 3) SM has EW phase transition X

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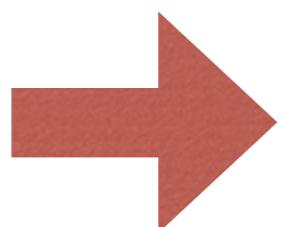


# We need New Physics

Couple to the SM

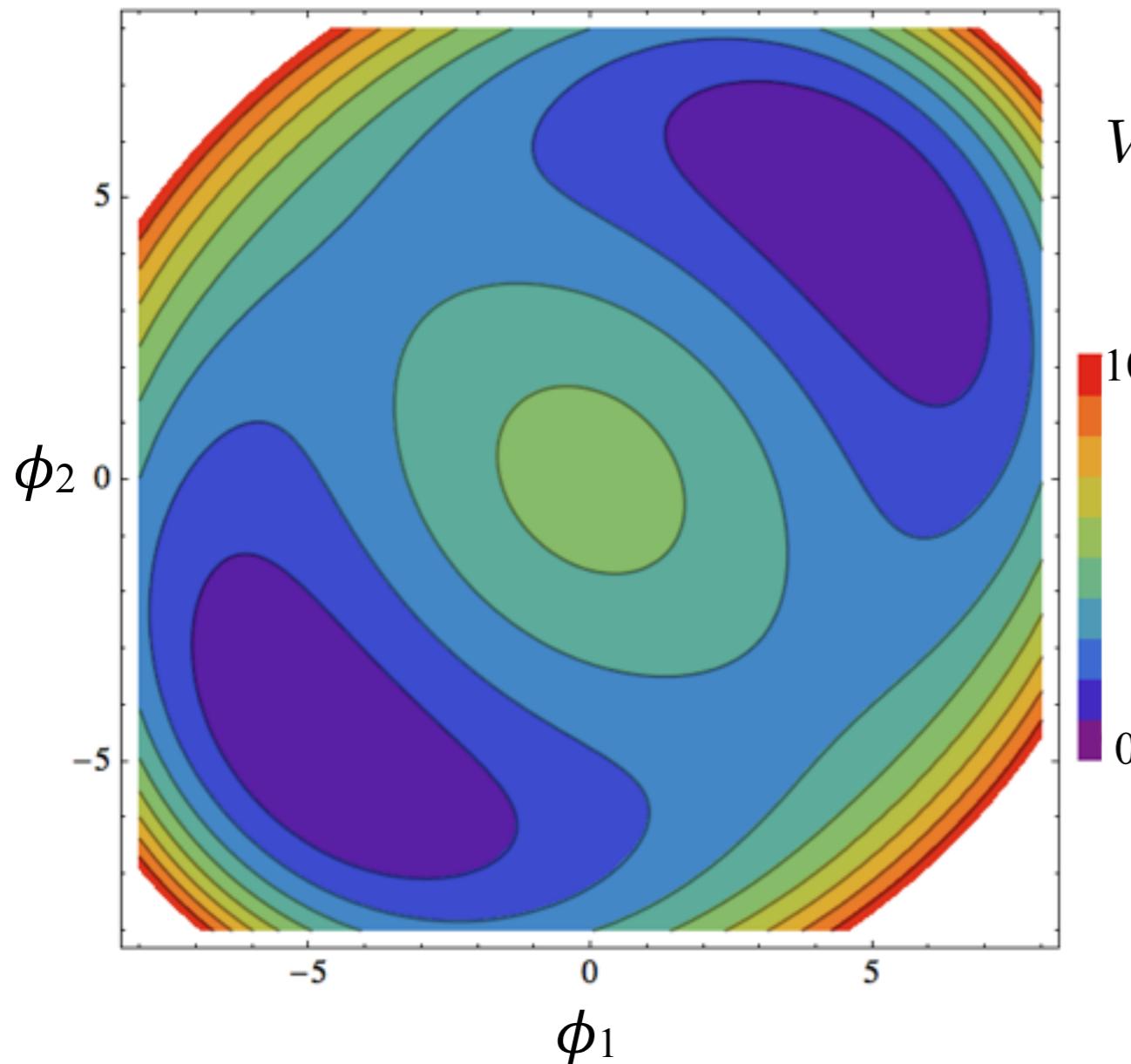
Extra CP violation

New light scalars can give a first  
order phase transition



Two Higgs Doublet Model

# Two is better than one



$$\begin{aligned}
 V = & h_1 \left( \Phi_1^\dagger \Phi_1 - \frac{v_1^2}{2} \right)^2 + h_2 \left( \Phi_2^\dagger \Phi_2 - \frac{v_2^2}{2} \right)^2 \\
 & + h_3 \left[ \left( \Phi_1^\dagger \Phi_1 - \frac{v_1^2}{2} \right) + \left( \Phi_2^\dagger \Phi_2 - \frac{v_2^2}{2} \right) \right]^2 \\
 & + h_4 \left[ (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) - (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) \right] \\
 & + h_5 \left( \text{Re}(\Phi_1^\dagger \Phi_2) - \frac{v_1 v_2}{2} \cos \xi \right)^2 \\
 & + h_6 \left( \text{Im}(\Phi_1^\dagger \Phi_2) - \frac{v_1 v_2}{2} \sin \xi \right)^2
 \end{aligned}$$

minimum at:

$$\langle \Phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix} \quad \langle \Phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 e^{i\xi} \end{pmatrix}$$

$\sqrt{v_1^2 + v_2^2} = 246 \text{ GeV}$

CP violation!

# There are 5 Higgs bosons

2 Higgs doublets

$$\Phi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1^+ \\ v_1 + \rho_1 + i\eta_1 \end{pmatrix} \quad \Phi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_2^+ \\ v_2 e^{i\xi} + \rho_2 + i\eta_2 \end{pmatrix}$$

→ 8 degrees of freedom →<sup>W,Z</sup> 5 Higgs bosons

2 charged Higgses:  $H^\pm = -\sin \beta \phi_1^\pm + \cos \beta \phi_2^\pm$

$$\tan \beta = \frac{v_2}{v_1}$$

Assume no CP violation  $\xi=0$ , 3 neutral Higgses:

$$A^0 = -\sin \beta \eta_1 + \cos \beta \eta_2 \quad \text{--- CP odd}$$

$$H^0 = \cos \alpha \rho_1 + \sin \alpha \rho_2$$

125 GeV Higgs →

$$h^0 = -\sin \alpha \rho_1 + \cos \alpha \rho_2 \quad \boxed{\quad} \quad \text{CP even}$$

# New Higgs couplings

Assume no CP violation,  $\xi=0$

Type II 2HDM:

$\Phi_1$  couples to down-type quarks and leptons

$\Phi_2$  couples to up-type quarks,

Higgs boson(s)  
couplings in units of  
125 GeV Higgs  
couplings:

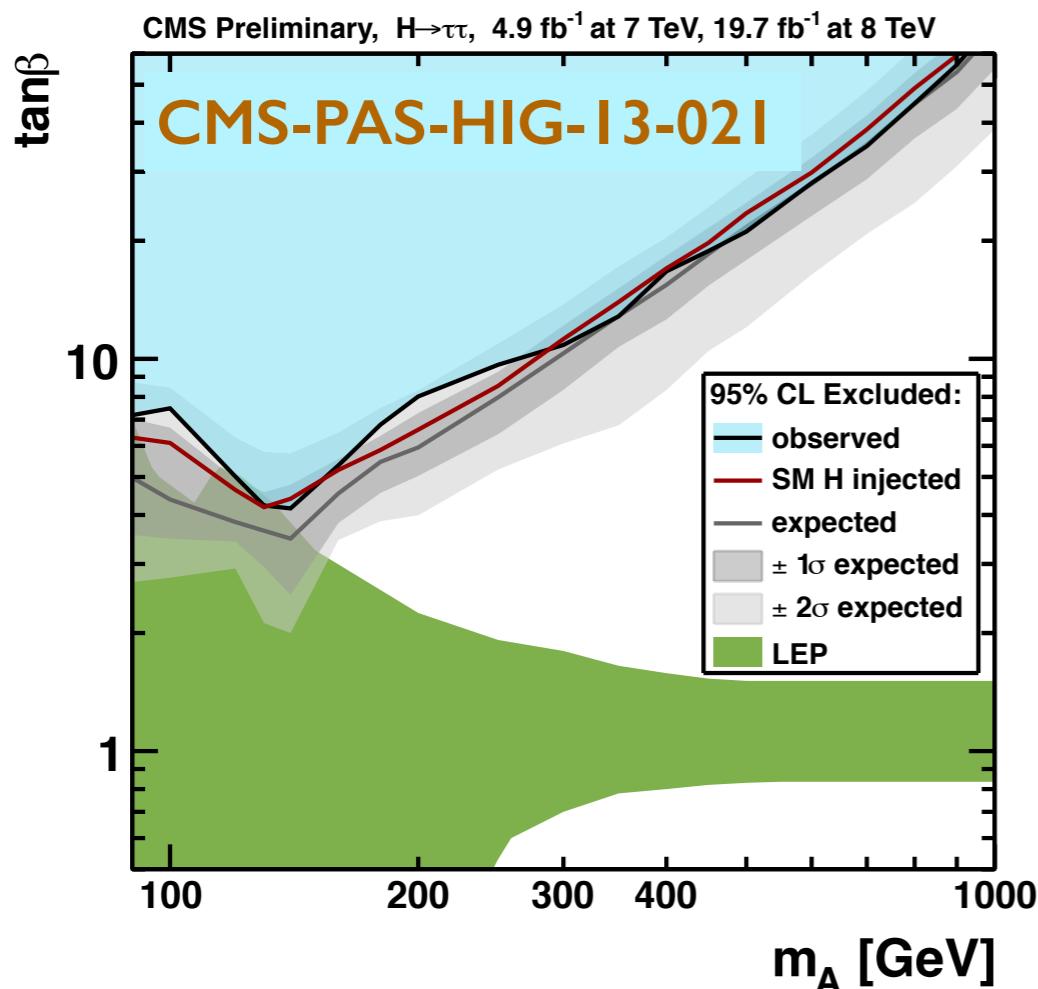
	$h^0$	$H^0$	$A^0$
$\chi_u$	$\frac{\cos \alpha}{\sin \beta}$	$\frac{\sin \alpha}{\sin \beta}$	$-\cot \beta$
$\chi_{d,e}$	$-\frac{\sin \alpha}{\cos \beta}$	$\frac{\cos \alpha}{\cos \beta}$	$-\tan \beta$
$\chi_V$	$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$	X

alignment limit:  $\beta - \alpha = \frac{\pi}{2} \rightarrow h^0$  couplings are SM-like

# How to find Heavy Higgses

Assume no CP violation,  $\xi=0$

- Heavy Higgs decays:  $H \rightarrow \tau\tau$

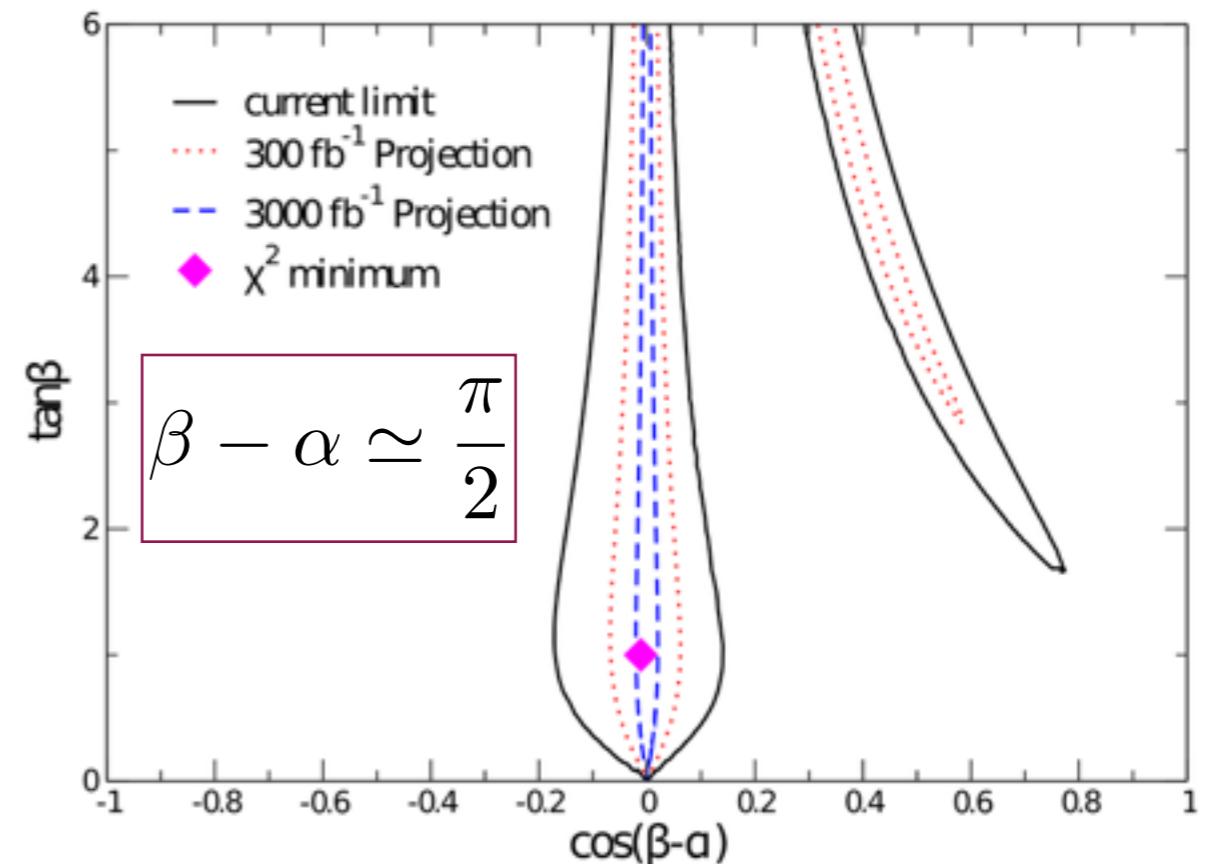


- Electroweak precision parameters

$$M_{H^\pm} \simeq M_{A^0/H^0/h^0}$$

Grimus, et al, arxiv: 0711.4022

- Deviations from SM Higgs couplings Chen, Dawson, arxiv: 1305.1624



- $\bar{B} \rightarrow X_s \gamma$  decays:

$$M_{H^\pm} \geq 380 \text{ GeV}$$

Hermann, et al, arxiv: 1208.2788

# New Higgs potential at high T

Assume no CP violation,  $\xi=0$

Consider the zero temperature 2HDM potential:

$$V_0 = -\mu_1^2 \Phi_1^\dagger \Phi_1 - \mu_2^2 \Phi_2^\dagger \Phi_2 + \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \lambda_2 (\Phi_2^\dagger \Phi_2)^2$$

$$+ \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{\lambda_5}{2} \left[ (\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2 \right]$$

simplified

At finite temperature, Coleman-Weinberg type loop expansion:

over simplified

e.g. light Higgses  
 $\lambda \ll g_W^2$

$$V_T = \mu^2 T^2 v^2 - \delta T v^3 + V_0(v_1, v_2)$$

$g_W^2 + \left(\frac{v_2}{v} y_t\right)^2$

$g_W^3 + \left(\frac{v_2}{v} y_t\right)^3$

\* Top-Yukawa is the most important one

competition between  $T^2 v^2$  and  $T v^3$  terms



order of the phase transition

# New Higgs potential at high T

Assume no CP violation,  $\xi=0$

In general, with a softly broken  $Z_2$ :

$$V = -\mu_1^2 \Phi_1^\dagger \Phi_1 - \mu_2^2 \Phi_2^\dagger \Phi_2 - \frac{\mu^2}{2} (\Phi_1^\dagger \Phi_2 + \Phi_2^\dagger \Phi_1) + \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\ + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2]$$

We have  $m_{H^\pm}, m_{h^0} \gg M_W$  not simple!

Finite T potential  
depends on:

- 2 mixing angles
- $\mu$
- 5 Higgs masses

- W/Z masses
- Top-quark mass

we know

set  $m_{h^0} = 125$  GeV

→ Left with:  $m_{A^0}, m_{H^0}, m_{H^\pm}, \mu, \beta, \alpha$

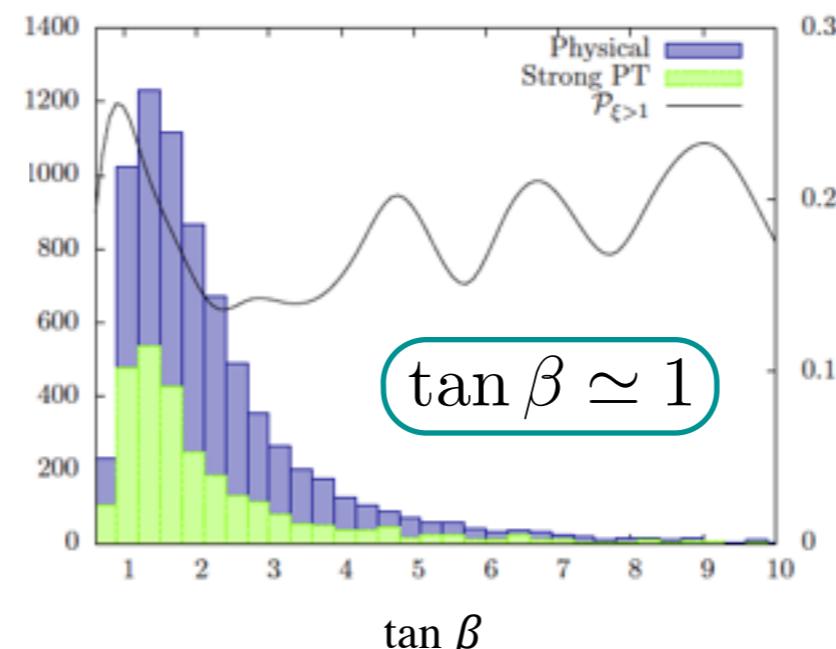
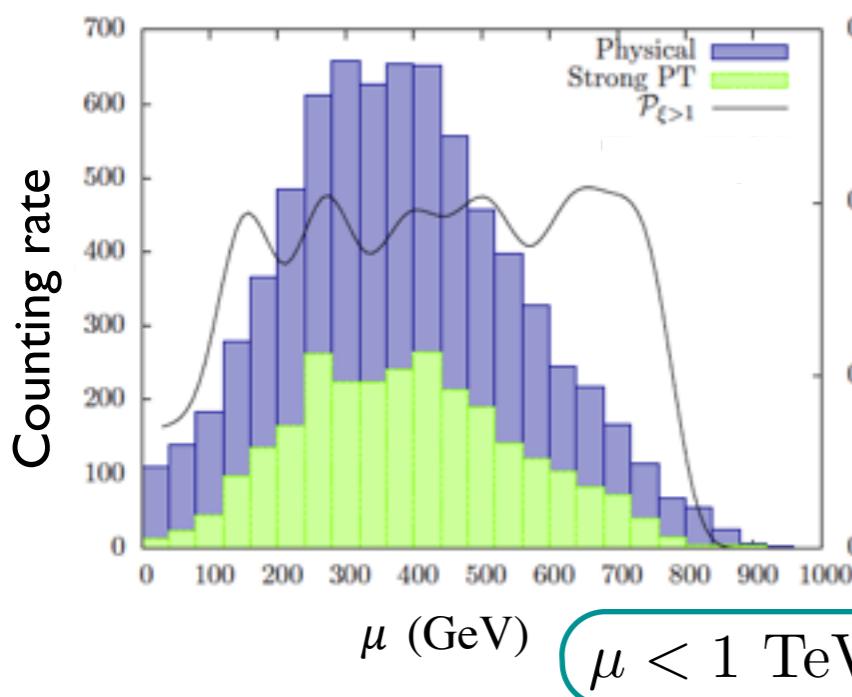
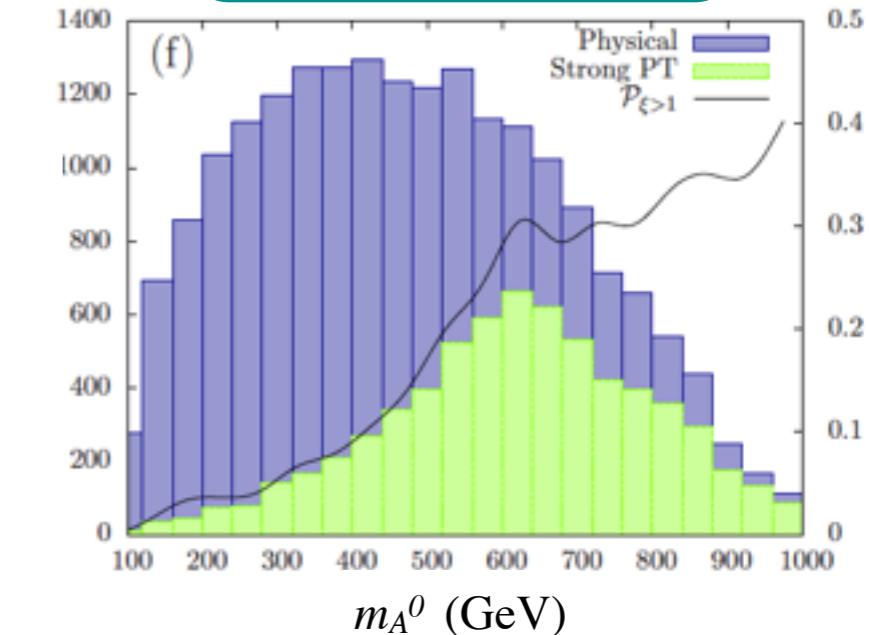
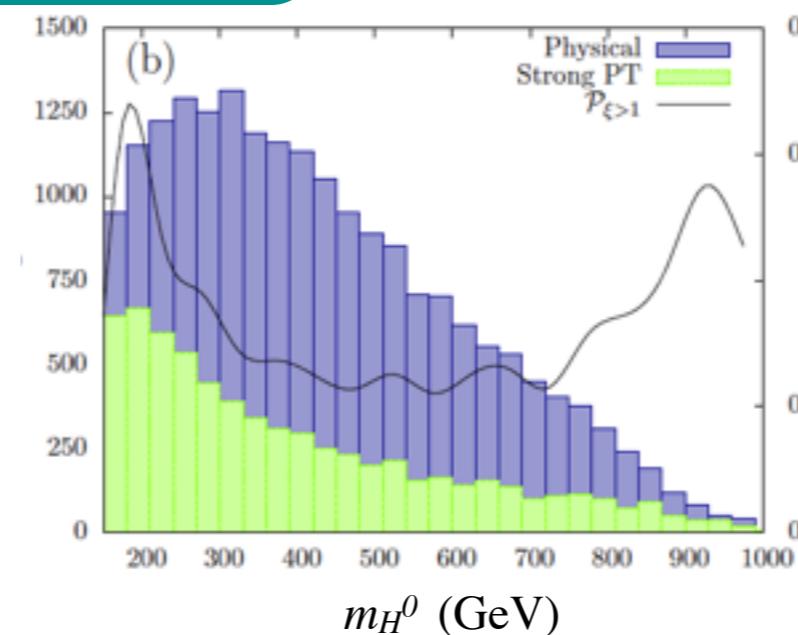
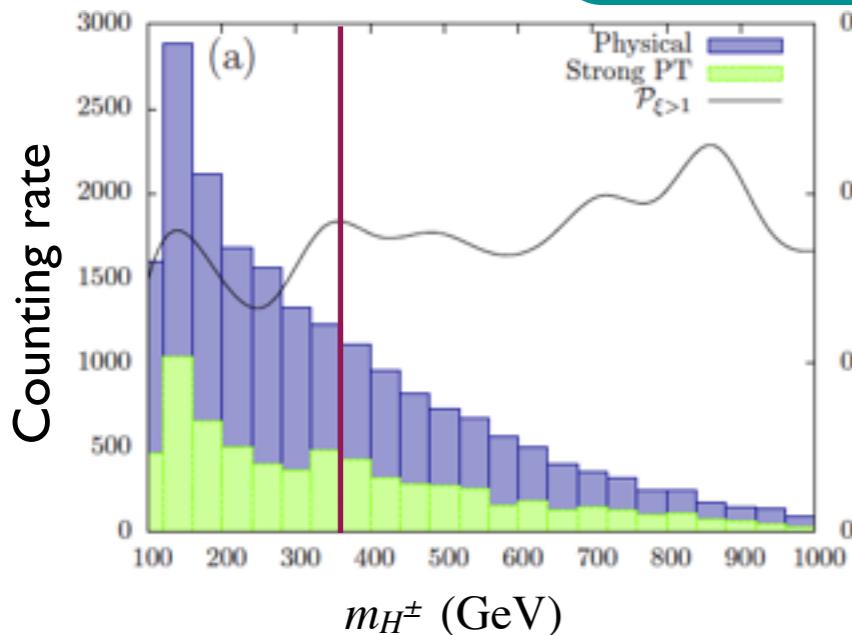
# First order phase transition ✓

Assume no CP violation,  $\xi=0$

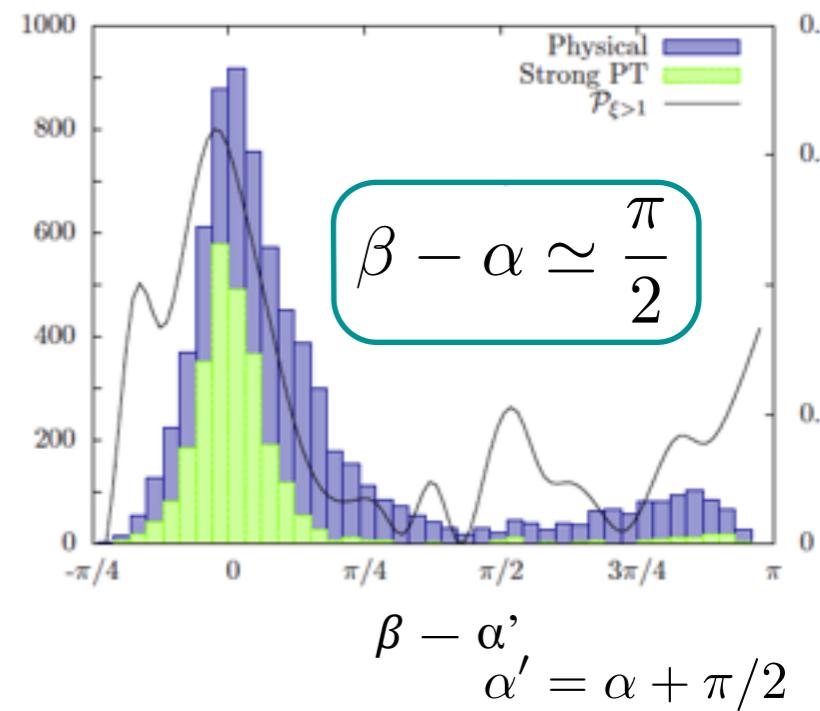
Dorsch, Huber, No, arxiv: 1305.6610

$$m_{A^0} > m_{H^0} \gtrsim m_{H^\pm}$$

$$m_{A^0} \gtrsim 400 \text{ GeV}$$



$$\tan \beta \simeq 1$$



# We also need CP violation

2HDMs can have a  
first order  
EW phase transition



Provides  
out-of-equilibrium  
conditions

First order phase transition is awesome,  
but not enough!

We also need extra CP violation



2HDMs can have CP violation in the Higgs sector

But complicated...

# CP violation mixes 3 Higgses

There is CP violation:  $\xi \neq 0$

$$\begin{aligned}
V = & h_1 \left( \Phi_1^\dagger \Phi_1 - \frac{v_1^2}{2} \right)^2 + h_2 \left( \Phi_2^\dagger \Phi_2 - \frac{v_2^2}{2} \right)^2 + h_3 \left[ \left( \Phi_1^\dagger \Phi_1 - \frac{v_1^2}{2} \right) + \left( \Phi_2^\dagger \Phi_2 - \frac{v_2^2}{2} \right) \right]^2 \\
& + h_4 \left[ (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) - (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) \right] \\
& + \boxed{h_5 \left( \text{Re}(\Phi_1^\dagger \Phi_2) - \frac{v_1 v_2}{2} \cos \xi \right)^2 + h_6 \left( \text{Im}(\Phi_1^\dagger \Phi_2) - \frac{v_1 v_2}{2} \sin \xi \right)^2}
\end{aligned}$$

- Charged Higgses are the same with mass  $\frac{h_4}{2} v^2$
- Neutral Higgses are a mixture of CP-odd ( $A^0$ ) and CP-even ( $\rho_1$  and  $\rho_2$ ) states:

$$\begin{pmatrix} \tilde{h} \\ \tilde{H} \\ \tilde{A} \end{pmatrix} = \begin{pmatrix} -s_\alpha c_{\alpha_b} & c_\alpha c_{\alpha_b} & s_{\alpha_b} \\ s_\alpha s_{\alpha_b} s_{\alpha_c} + c_\alpha c_{\alpha_c} & s_\alpha c_{\alpha_c} - c_\alpha s_{\alpha_b} s_{\alpha_c} & c_{\alpha_b} s_{\alpha_c} \\ s_\alpha s_{\alpha_b} c_{\alpha_c} - c_\alpha s_{\alpha_c} & -s_\alpha s_{\alpha_c} - c_\alpha s_{\alpha_b} c_{\alpha_c} & c_{\alpha_b} c_{\alpha_c} \end{pmatrix} \begin{pmatrix} \rho_1 \\ \rho_2 \\ A^0 \end{pmatrix}$$

# Mass states are a mess

There is CP violation:  $\xi \neq 0$

Mass<sup>2</sup> matrix in the basis  $(\rho_1, \rho_2, A^0)^T$

$$M^2 = v^2 \begin{pmatrix} 2(h_1 + h_3)c_\beta^2 + \frac{1}{2}(h_5c_\xi^2 + h_6s_\xi^2)s_\beta^2 & \left[2h_3 + \frac{1}{2}(h_5c_\xi^2 + h_6s_\xi^2)\right]s_\beta c_\beta & \frac{1}{2}(h_6 - h_5)s_\beta c_\xi s_\xi \\ \left[2h_3 + \frac{1}{2}(h_5c_\xi^2 + h_6s_\xi^2)\right]s_\beta c_\beta & 2(h_2 + h_3)s_\beta^2 + \frac{1}{2}(h_5c_\xi^2 + h_6s_\xi^2)c_\beta^2 & \frac{1}{2}(h_6 - h_5)c_\beta c_\xi s_\xi \\ \frac{1}{2}(h_6 - h_5)s_\beta c_\xi s_\xi & \frac{1}{2}(h_6 - h_5)c_\beta c_\xi s_\xi & \frac{1}{4}(h_5 + h_6 + (h_6 - h_5)c_{2\xi}) \end{pmatrix}$$

If  $h_5 = h_6 \rightarrow$  No dependence on  $\xi$  !

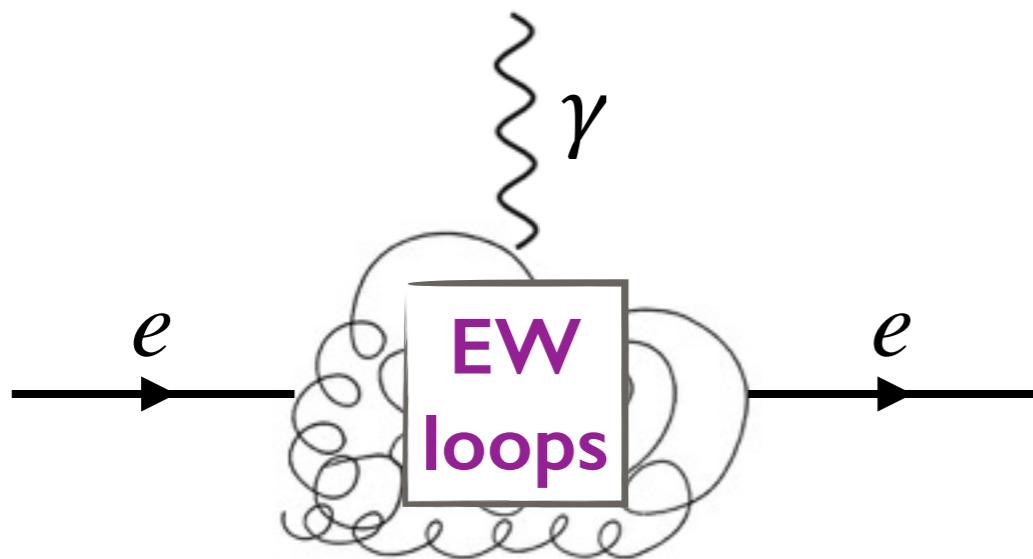
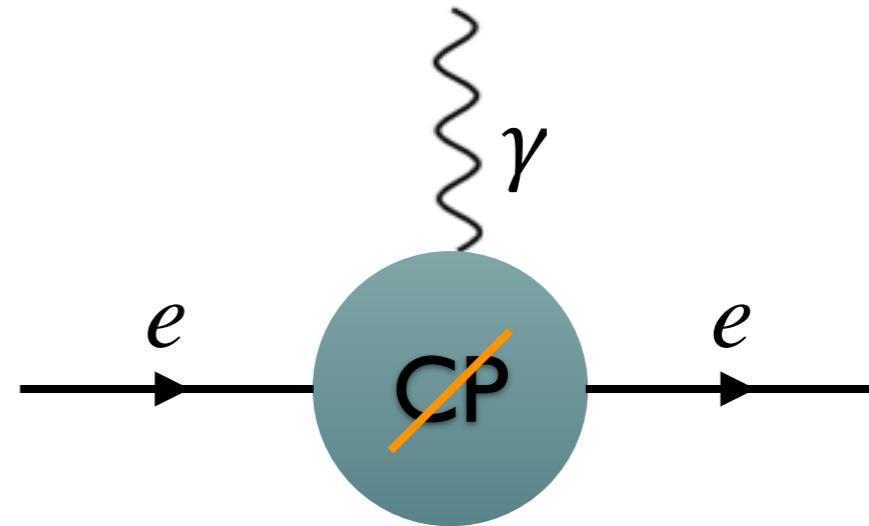
Diagonalize to find the mass eigenstates with:

$$R = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha_c & \sin \alpha_c \\ 0 & -\sin \alpha_c & \cos \alpha_c \end{pmatrix} \begin{pmatrix} \cos \alpha_b & 0 & \sin \alpha_b \\ 0 & 1 & 0 \\ -\sin \alpha_b & 0 & \cos \alpha_b \end{pmatrix} \begin{pmatrix} -\sin \alpha & \cos \alpha & 0 \\ \cos \alpha & \sin \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

In general, it is a mess!  2 new angles

# CP violation shows up in EDMs

Electric Dipole Moments  
are a good measure of CP  
violation



SM prediction for electron  
EDM is very small:

$$d_e < 10^{-38} \text{ } e\cdot\text{cm}$$

10 orders of magnitude!  
But still very small

Experimental bound:  $d_e < 0.87 \times 10^{-28} \text{ } e\cdot\text{cm}$

ACME, Science, 343 (2014)

# We expect small CP violation

Let's do some estimates

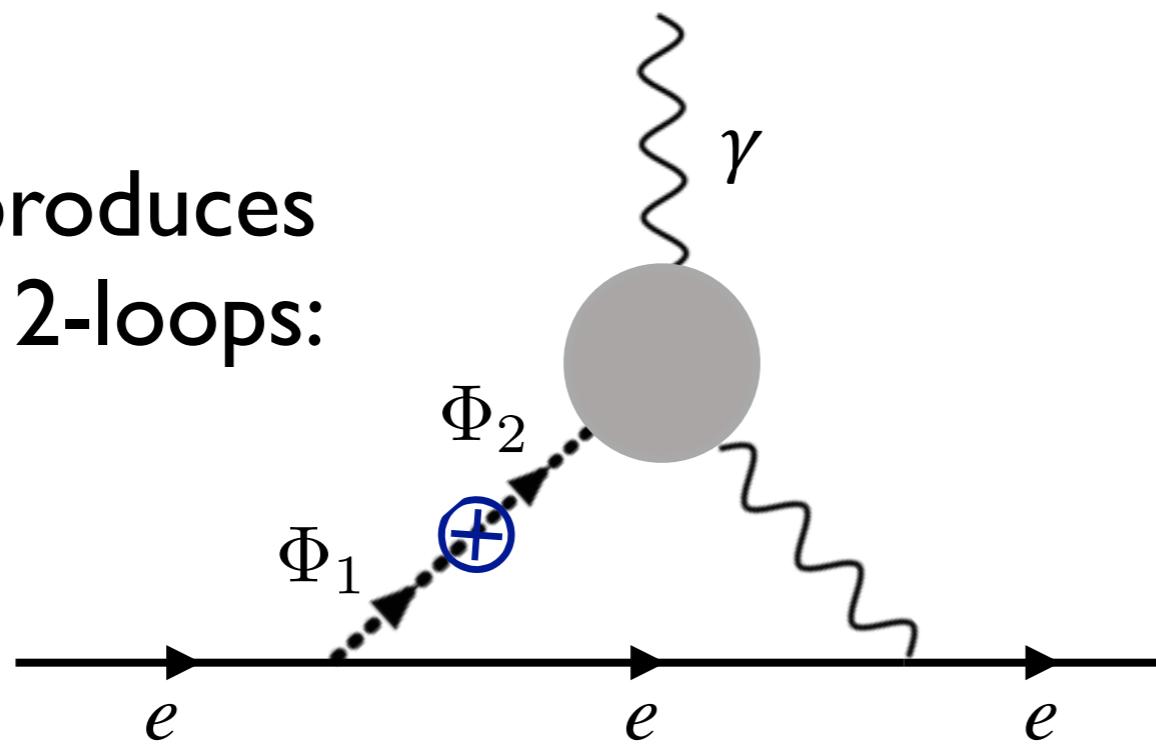
Experiment says:  $d_e < 0.87 \times 10^{-28} \text{ e}\cdot\text{cm}$

New Physics gives:  $d_e \sim \left(\frac{\alpha_e}{4\pi}\right)^n \frac{m_e}{\Lambda^2} \sin \xi$

$n$ : number of loops

$\Lambda$ : scale of NP

2HDM produces  
EDMs in 2-loops:



For  $\Lambda \sim O(100 \text{ GeV})$

$$\boxed{\sin \xi < 0.1}$$

# Perturbative mass matrix

$\sin \xi < 0.1$

Ipek, PRD D89 (2014) 073012, arXiv:1310.6790

$$M^2 \simeq v^2 \begin{pmatrix} 2(h_1 + h_3)c_\beta^2 + \frac{h_5}{2}s_\beta^2 & \frac{1}{4}(4h_3 + h_5)s_{2\beta} & 0 \\ \frac{1}{4}(4h_3 + h_5)s_{2\beta} & 2(h_2 + h_3)c_\beta^2 + \frac{h_5}{2}s_\beta^2 & 0 \\ 0 & 0 & \frac{h_6}{2} \end{pmatrix}$$

CP conserving mass<sup>2</sup> matrix

$$+ \frac{h_6 - h_5}{2}v^2 \begin{pmatrix} \xi^2 s_\beta^2 & \frac{1}{2}\xi^2 s_{2\beta} & \xi s_\beta \\ \frac{1}{2}\xi^2 s_{2\beta} & \xi^2 c_\beta^2 & \xi c_\beta \\ \xi s_\beta & \xi c_\beta & -\xi^2 \end{pmatrix}$$

Perturbation due to CP violation, mixing

Expansion works when

$$C \equiv \frac{h_6 - h_5}{2} \simeq \frac{m_{A^0}^2}{v^2} - \frac{\mu^2}{v^2 \sin(2\beta)} : \text{not too large}$$

# Perturbative mass eigenstates

Use non-degenerate perturbation theory



$$m_{A^0} \neq m_{H^0}, m_{h^0}$$

these are neutral Higgs masses when  $\xi = 0$

eigenstates

$$\tilde{h} \simeq h^0 - \xi a \cos \gamma A^0$$

masses

$$m_{\tilde{h}}^2 \simeq m_{h^0}^2 + \xi^2 C v^2 \cos^2 \gamma (1 - a)$$

$$\tilde{H} \simeq H^0 + \xi b \sin \gamma A^0$$

$$m_{\tilde{H}}^2 \simeq m_{H^0}^2 + \xi^2 C v^2 \sin^2 \gamma (1 - b)$$

$$\begin{aligned} \tilde{A} \simeq & A^0 + \xi a \cos \gamma h^0 \\ & - \xi b \sin \gamma H^0 \end{aligned}$$

$$\begin{aligned} m_{\tilde{A}}^2 \simeq & m_{A^0}^2 - \xi^2 C v^2 \sin^2 \gamma (1 - b) \\ & - \xi^2 C v^2 \cos^2 \gamma (1 - a) \end{aligned}$$

$$a \equiv \frac{C v^2}{m_{A^0}^2 - m_{h^0}^2}, \quad b \equiv \frac{C v^2}{m_{A^0}^2 - m_{H^0}^2}, \quad C v^2 \simeq m_{A^0}^2 - \frac{\mu^2}{\sin(2\beta)}, \quad \gamma \equiv \alpha + \beta$$

# Perturbative rotation matrix

$$R = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha_c & \sin \alpha_c \\ 0 & -\sin \alpha_c & \cos \alpha_c \end{pmatrix} \begin{pmatrix} \cos \alpha_b & 0 & \sin \alpha_b \\ 0 & 1 & 0 \\ -\sin \alpha_b & 0 & \cos \alpha_b \end{pmatrix} \begin{pmatrix} -\sin \alpha & \cos \alpha & 0 \\ \cos \alpha & \sin \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

diagonalizes mass<sup>2</sup> matrix  $R M^2 R^{-1}$

New mixing  
angles:

- $\sin \alpha_b \simeq -\xi \frac{m_{A^0}^2 - \bar{\mu}^2}{m_{A^0}^2 - m_{h^0}^2} \sin \gamma$        $\bar{\mu}^2 \equiv \frac{\mu^2}{\sin(2\beta)}$
- $\sin \alpha_c \simeq \xi \frac{m_{A^0}^2 - \bar{\mu}^2}{m_{A^0}^2 - m_{H^0}^2} \cos \gamma$        $\gamma \equiv \alpha + \beta$

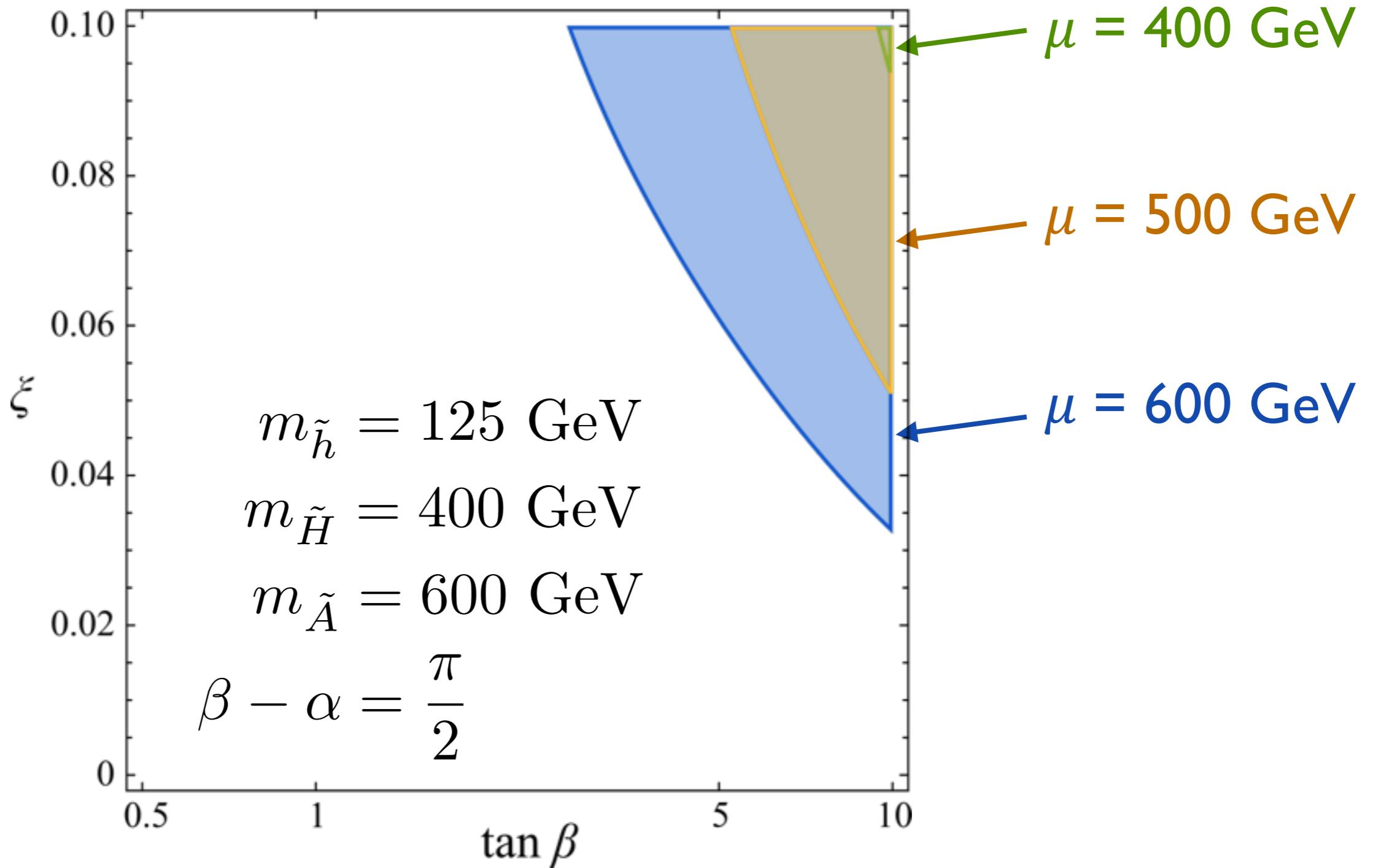
New mass eigenstates:

$$\star \tilde{h} \simeq h^0 + \sin \alpha_b A^0 \quad \star \tilde{H} \simeq H^0 + \sin \alpha_c A^0$$

$$\star \tilde{A} \simeq A^0 - \sin \alpha_b h^0 - \sin \alpha_c H^0$$

# How good is it?

Regions of  $\tan\beta - \xi$  space where corrections are larger than 10%



# Perturbative Yukawa couplings

Coupling constants in units of 125 GeV Higgs couplings

	$h^0$	$H^0$	$A^0$
$\chi_u$	$\frac{\cos \alpha}{\sin \beta}$	$\frac{\sin \alpha}{\sin \beta}$	
$\tilde{\chi}_u$			$-\cot \beta$
$\chi_{d,e}$	$-\frac{\sin \alpha}{\cos \beta}$	$\frac{\cos \alpha}{\cos \beta}$	
$\tilde{\chi}_{d,e}$			$-\tan \beta$
$\chi_V$	$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$	

# Perturbative Yukawa couplings

Coupling constants in units of 125 GeV Higgs couplings

	$\tilde{h}$	$\tilde{H}$	$\tilde{A}$
$\chi_u$	$\frac{\cos \alpha}{\sin \beta}$	$\frac{\sin \alpha}{\sin \beta}$	$\frac{\xi}{\sin \beta} (a \cos \gamma \cos \alpha + b \sin \gamma \sin \alpha)$
$\tilde{\chi}_u$	$\xi a \cos \gamma \cot \beta$	$-\xi b \sin \gamma \cot \beta$	$-\cot \beta$
$\chi_{d,e}$	$-\frac{\sin \alpha}{\cos \beta}$	$\frac{\cos \alpha}{\cos \beta}$	$\frac{\xi}{\cos \beta} (b \sin \gamma \cos \alpha - a \cos \gamma \sin \alpha)$
$\tilde{\chi}_{d,e}$	$\xi a \cos \gamma \tan \beta$	$-\xi b \sin \gamma \tan \beta$	$-\tan \beta$
$\chi_V$	$\sin(\beta - \alpha)$	$\cos(\beta - \alpha)$	$\xi (a \cos \gamma \sin(\beta - \alpha) + b \sin \gamma \cos(\beta - \alpha))$

$$+O(\xi^2)$$

# Only one extra parameter

CP violating Higgs potential:

$$V = -\mu_1^2 \Phi_1^\dagger \Phi_1 - \mu_2^2 \Phi_2^\dagger \Phi_2 - (\mu^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}) + \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \lambda_2 (\Phi_2^\dagger \Phi_2)^2 \\ + \lambda_3 (\Phi_1^\dagger \Phi_1)(\Phi_2^\dagger \Phi_2) + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{\lambda_5}{2} \left[ (\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2 \right]$$

Finite T potential  
depends on:

- 2 mixing angles
- $\mu$
- 5 Higgs masses

- W/Z masses
- Top-quark mass

we know

set  $m_{h^0} = 125$  GeV

→ Left with:  $m_{A^0}, m_{H^0}, m_{H^\pm}, \mu, \beta, \alpha$   
+  $\xi$

# Constraints from EDMs

Inoue, Ramsey-Musolf, Zhang, arxiv: 1403.4257

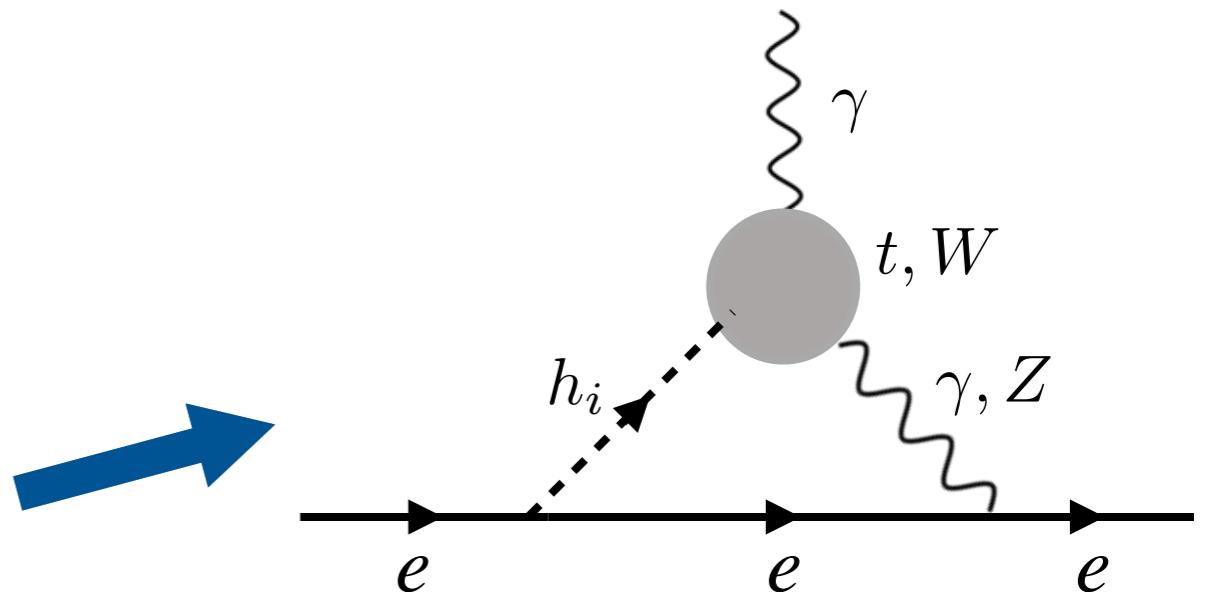
Chen, Dawson, Zheng, arxiv: 1503.01114

...

$$d_e < 0.87 \times 10^{-28} \text{ e}\cdot\text{cm}$$

All of the Higgses have CP-odd  
couplings to electrons

There are quite a few  
Barr-Zee diagrams, and more!

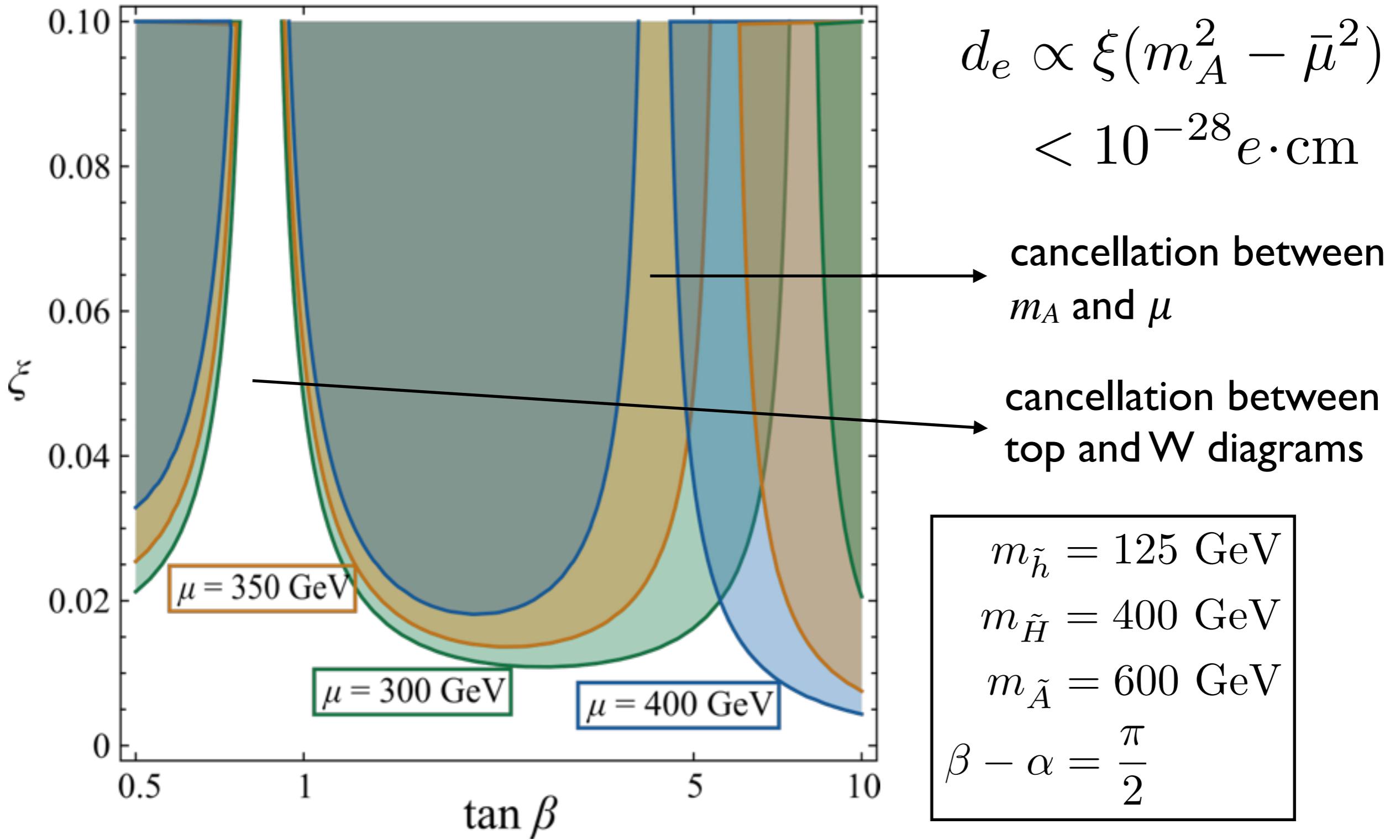


There are other EDMs too, e.g. neutron

Complicated problem with all the new angles and the mixings

But at the end:  $d_e \propto \xi(m_A^2 - \bar{\mu}^2) F(\text{masses, etc})$

# Constraints from EDMs



# Constraints from Higgs searches

SM Higgs couplings change for  $\beta - \alpha \neq \frac{\pi}{2}$

Chen, Dawson, Zheng, arxiv: 1503.01114

Also with  $\xi \neq 0$   $\longrightarrow$  
$$\frac{\Gamma(\tilde{h} \rightarrow bb)}{\Gamma_{\text{SM}}} \sim 1 + O(\xi^2)$$
 too small!

New vector boson couplings for  $\xi \neq 0$

# Summary

- We should be able to explain baryogenesis
- Progress has been made in finding a first-order phase transition in 2HDMs
- CP violation is crucial but complicated in 2HDM
- EDM constraints allow for a small CP-violating phase
- Would also affect Higgs measurements
- Perturbatively including CP violation in baryogenesis studies?