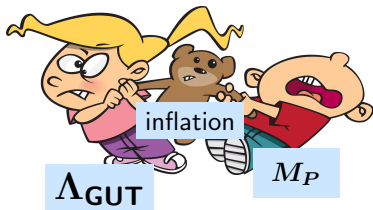


Towards chaotic inflation in SUSY GUTs

The rebirth of D-term inflation

[arxiv\[hep-ph\] 1406.6300](https://arxiv.org/abs/hep-ph/1406.6300), [1410.xxxx](#)



Valerie Domcke

SISSA, INFN Trieste

in collaboration with
W. Buchmüller, K. Schmitz
and T. Yanagida



in**o**visibles



Motivation

Data: $r \sim 0.1 (?) \rightarrow V_0^{1/4} \sim 10^{16} \text{ GeV} \sim \Lambda_{\text{GUT}}$



→ How can we link the scale of inflation to a scale of particle physics?

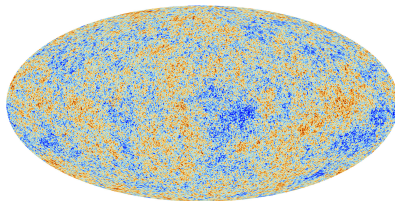
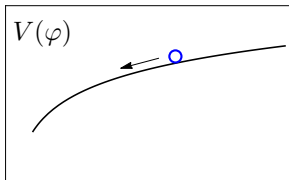
Theory: supersymmetric D-term hybrid inflation

- consistent in supergravity
- no cosmic string problem
- new regime with very different predictions

- small- and large-field models of inflation
- the chaotic regime of D-term hybrid inflation
- dynamical generation of a Fayet-Illiopoulos term
- conclusion and outlook

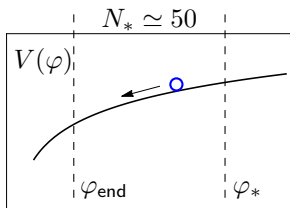
Some observations about slow-roll inflation

The basics:



Some observations about slow-roll inflation

The basics:



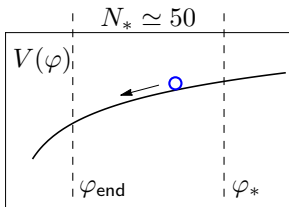
$$\epsilon = \frac{1}{2} \left(\frac{V'}{V} \right)^2, \quad \eta = \frac{V''}{V}$$

$$A_s = \frac{V_*}{24\pi^2 \epsilon_*}, \quad n_s = 1 - 6\epsilon_* + 2\eta_*, \quad r = 16\epsilon_*$$

$$\Delta\varphi \gtrsim 5 M_{\text{P}} \frac{N}{50} \sqrt{\frac{r}{0.1}} \quad [\text{Lyth '96}]$$

Some observations about slow-roll inflation

The basics:

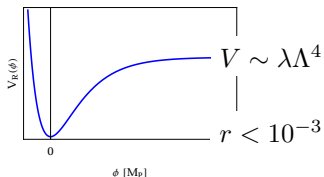


$$\epsilon = \frac{1}{2} \left(\frac{V'}{V} \right)^2, \quad \eta = \frac{V''}{V}$$

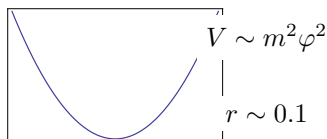
$$A_s = \frac{V_*}{24\pi^2 \epsilon_*}, \quad n_s = 1 - 6\epsilon_* + 2\eta_*, \quad r = 16\epsilon_*$$

$$\Delta\phi \gtrsim 5 M_{\text{P}} \frac{N}{50} \sqrt{\frac{r}{0.1}} \quad [\text{Lyth '96}]$$

Small-field models:



Large-field models:



Can large-field models be connected to grand unification?

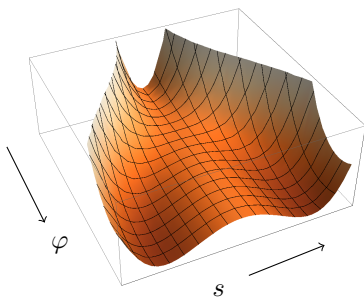
D-term hybrid inflation in a nutshell

A supersimple model: [Binetruy, Dvali '96; Halyo '96] [Kawasaki, Yamaguchi, Yanagida '00]

- $U(1)$ gauge symmetry with Fayet-Iliopoulos term ξ
- $W = \lambda\phi S_+ S_-$, $K = \frac{1}{2}(\phi + \bar{\phi})^2 + S_+ \bar{S}_+ + S_- \bar{S}_-$

- ξ : vacuum energy
- ϕ : inflaton φ , tree-level flat
- S_{\pm} : waterfall field, $|S_+|$ destabilize at

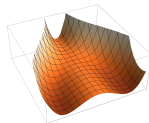
$$\varphi_c = \frac{g}{\lambda} \sqrt{2\xi}$$



D-term hybrid inflation in a nutshell

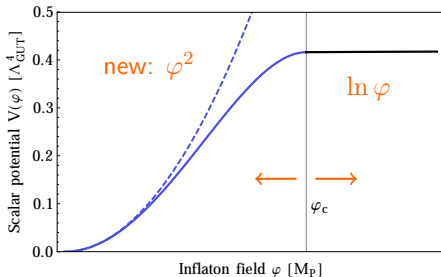
A supersimple model: [Binetruy, Dvali '96; Halyo '96] [Kawasaki, Yamaguchi, Yanagida '00]

- $U(1)$ gauge symmetry with Fayet-Iliopoulos term ξ
- $W = \lambda\phi S_+ S_-$, $K = \frac{1}{2}(\phi + \bar{\phi})^2 + S_+ \bar{S}_+ + S_- \bar{S}_-$



- ξ : vacuum energy
- ϕ : inflaton φ , tree-level flat
- S_{\pm} : waterfall field, $|S_+|$ destabilize at

$$\varphi_c = \frac{g}{\lambda} \sqrt{2\xi}$$



$\lambda \gg g$: D-term hybrid inflation. $\lambda \ll g$: new regime

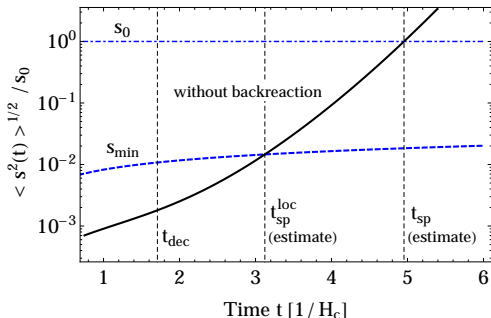
Quantum mechanics: tachyonic growth

Close to the critical point:

$$\varphi(t) \simeq \varphi_c + \dot{\varphi}_c t \quad \rightarrow \quad V(s; t) \simeq \frac{1}{2} g^2 \xi^2 - \frac{1}{2} D^3 t s^2 + \mathcal{O}(t^2, s^4)$$

→ tachyonic growth of s_k for small k : [Asaka, Buchmüller, Covi '01]

$$\ddot{s}_k + \left(k^2 e^{-2Ht} - \frac{9}{4} H^2 - D^3 t \right) s_k = 0$$

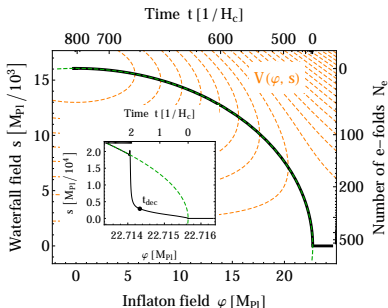
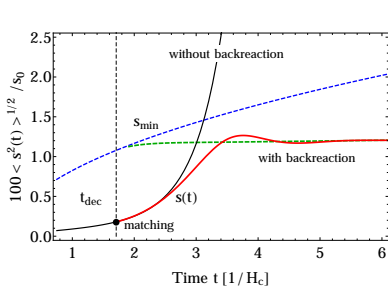


- spinodal time
 $\langle s^2(t) \rangle = s_{\text{vacuum}}$
- decoherence time
 $|s_k(t) \pi_{s_k}(t)| \gg \hbar$
 $\dot{s}/H > H/(2\pi)$

Classical evolution: backreaction

Backreaction sets in once we overshoot the minimum:

$$\ddot{\varphi} + 3H\dot{\varphi} + \frac{1}{2}\lambda^2 s^2 \varphi = 0, \quad \ddot{s} + 3H\dot{s} - \left(g^2 \xi - \frac{\lambda^2}{2} \varphi^2 \right) s + \frac{g^2}{2} s^3 = 0$$



inflaton field tracks *local* minimum. λ/g sets ratio of semi-axes.

Implications for cosmology and particle physics

After the critical point: $V(\varphi, s_{\min}(\varphi)) = \frac{1}{2}\lambda^2\xi\varphi^2\left(1 - \frac{1}{2}\frac{\varphi^2}{\varphi_c^2}\right)$
with $\varphi_c \propto g/\lambda > M_P$. $\rightarrow m^2\varphi^2$ inflation + small correction!

predictions:

$$A_s \simeq \frac{\lambda^4 \varphi_*^4 \varphi_c^2}{192\pi^2 g^2} \left(1 + \frac{1}{2} \frac{\varphi_*^2}{\varphi_c^2}\right), \quad n_s \simeq 1 - \frac{8}{\varphi_*^2} \left(1 - \frac{1}{4} \frac{\varphi_*^2}{\varphi_c^2}\right), \quad r \simeq \frac{32}{\varphi_*^2} \left(1 - \frac{\varphi_*^2}{\varphi_c^2}\right).$$

e.g. $g^2 = \frac{1}{2}$, $\lambda = 5 \cdot 10^{-4}$: $\varphi_* = 14.5 M_P$, $n_s = 0.963$, $r = 0.083$

constraints:

sufficient inflation after φ_c : $\lambda \lesssim 10^{-3}$, $\sqrt{\xi} \gtrsim 2 \times 10^{16}$ GeV

validity of this formalism: $\lambda \gtrsim 10^{-4}$, $\sqrt{\xi} \lesssim 2 \times 10^{17}$ GeV

no cosmic string problem!

chaotic inflation with slightly reduced r . GUT scale FI-term.

Final remarks on D-term 'hybrid' inflation

- We retrieve chaotic inflation from D-term hybrid inflation
- Inflation is linked to the decay of a GUT-scale false vacuum
- Interplay of quantum and classical dynamics
- If BICEP is dust, 'usual' D-term hybrid inflation is still viable for λ/g large enough (but tension with cosmic string bound)

[Buchmüller, VD, Kamada '13]

- Supergravity embedding of FI-term tricky

[Komargodski, Seibergy '09 and '10; Dienes, Thomas '10]

Dynamical supersymmetry breaking

based on strong $SP(1) \simeq SU(2)$ gauge group [Izawa, Yanagida '96, Intriligator, Thomas '96]

- field content below Λ : 4 'quarks' Q , 6 'mesons' M , 6 singlets Z
- perturbative superpotential $W = \lambda_i M_i Z_i$
- deformed moduli constraint

$$Pf(M) = M_1 M_2 - M_3 M_4 + M_5 M_6 = \Lambda^2, [Pf(M)]^2 = \det(M)$$

$\Rightarrow \langle M_i \rangle \neq 0 \rightarrow \langle F_{Z_i} \rangle \neq 0 \Rightarrow$ susy breaking

scale Λ via dimensional transmutation:

$$\Lambda = M_P \exp\left(-\frac{8\pi^2}{b g_s(M_P)}\right)$$

e.g. $g_s(M_P) = 1 \dots 4\pi \rightarrow \Lambda = 10^{10} \dots 10^{18}$ GeV

supersymmetry breaking by strong dynamics at scale Λ

Dynamical 'FI-term'

Gauge a $U(1)$ subgroup of the flavour symmetry

work in progress w K. Schmitz and T. Yanagida

$$\begin{array}{l} \text{fields} \\ U(1) \end{array} \left\| \begin{array}{ccccccc} & \overbrace{Q_1 \quad Q_2}^{M_+} & & \overbrace{Q_3 \quad Q_4}^{M_-} & & Z_a^0 & Z_+ & Z_- \\ +1/2 & +1/2 & -1/2 & -1/2 & 0 & +1 & -1 \end{array} \right.$$

D-term potential:

$$V_D = \frac{1}{2} g^2 (|M_+|^2 - |M_-|^2 + \dots)^2$$

with $\langle M_{\pm} \rangle^2 \simeq \frac{\lambda_{\mp}}{\lambda_{\pm}} \Lambda^2$ (and all other fields stabilized at zero):

$$\xi = \Lambda^2 \left(\frac{\lambda_-}{\lambda_+} - \frac{\lambda_+}{\lambda_-} \right) (1 - \mathcal{O}(g^2))$$

Effective GUT-scale FI-term generated

Conclusion and Outlook

Realization of D-term hybrid inflation with dynamical FI-term

- consistent in supergravity
- no cosmic string problem
- GUT scale arises naturally
- chaotic regime of D-term inflation with sizeable r

Outlook

- Further applications of dynamical FI-term
- $U(1)_{\text{FI}} \leftrightarrow G_{\text{GUT}}$?
- more results on the search of primordial B-modes coming soon !

Thank you!

Questions?