

Can one elucidate the dark matter problem if no new physics is found at LHC?

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“Long lived” expectations since the 1980s



Ref.TH.2910-CERN

THE INTEREST IN e^+e^- EXPERIMENTS AT LEP ENERGIES

John Ellis
CERN -- Geneva

ABSTRACT

A review is presented of various open theoretical questions and how LEP experiments may help answer them. Besides the conventional gauge theories of the strong, weak and electromagnetic interactions, LEP should also provide us with useful information relevant to grand unified theories, supersymmetry and dynamical symmetry breaking. New emphasis is laid on toponium as a laboratory for studying the weak interactions as well as the strong, and on the possibilities for testing technicolour at LEP.

Invited talk presented at the
International Conference on Experimentation at LEP
University of Uppsala, Sweden
June 16 to 20, 1980

Ref.TH.2910-CERN
29 June 1980

The interest in e^+e^- experiments at LEP energies

LEP should also provide us with useful information relevant to grand unified theories, supersymmetry and dynamical symmetry breaking

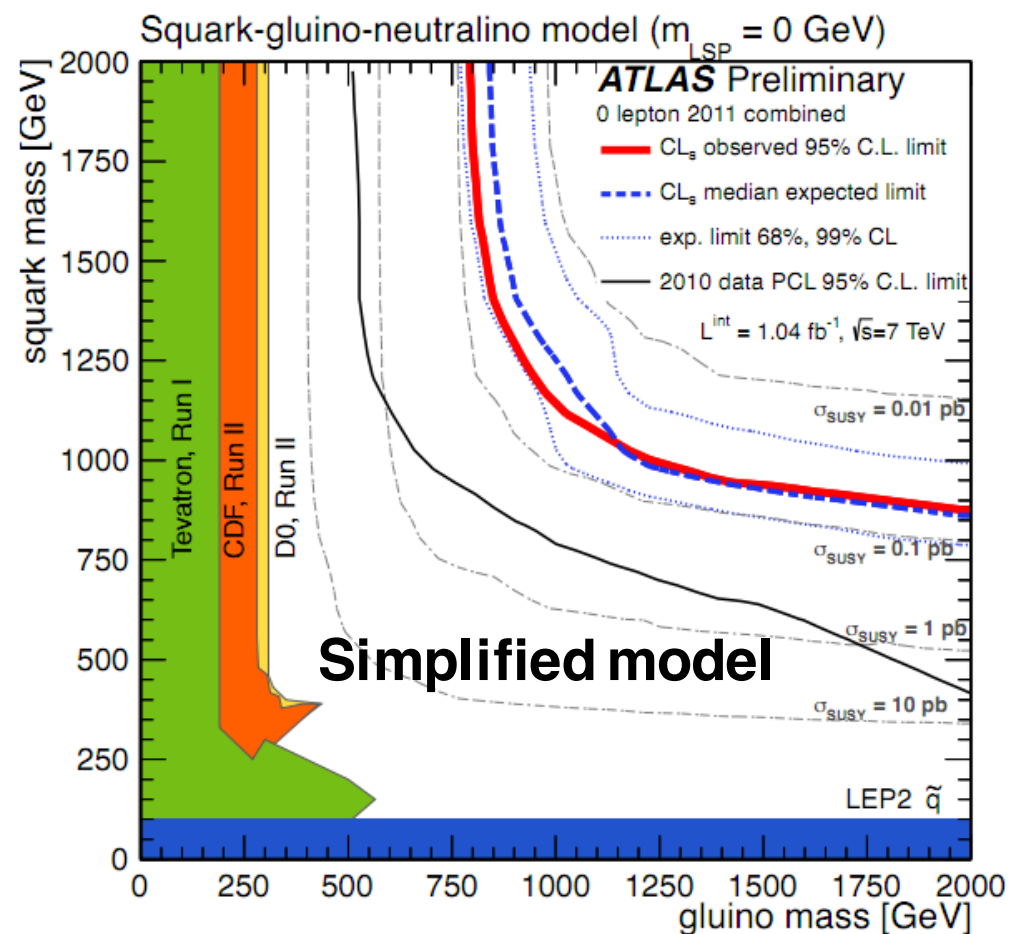
J. R. Ellis, 16th June, 1980

Would the LHC be the ultimate “kill-hope” machine ?

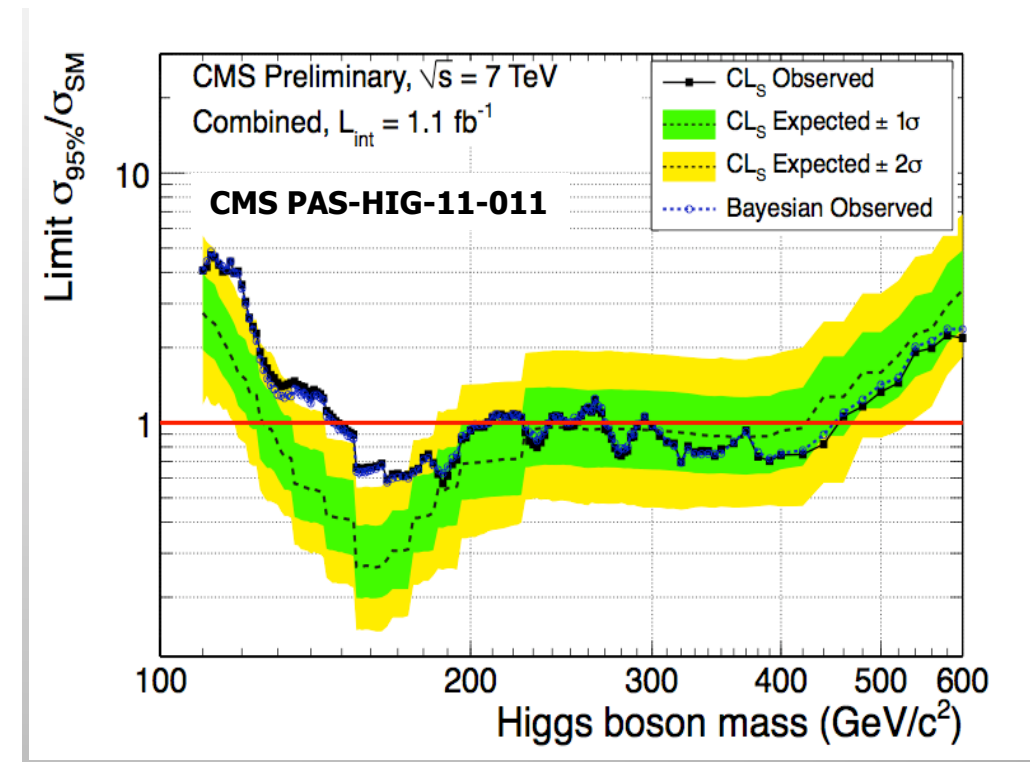
Quid of LHC?

LEP has not provide us with useful information relevant to grand unified theories, supersymmetry and dynamical symmetry breaking

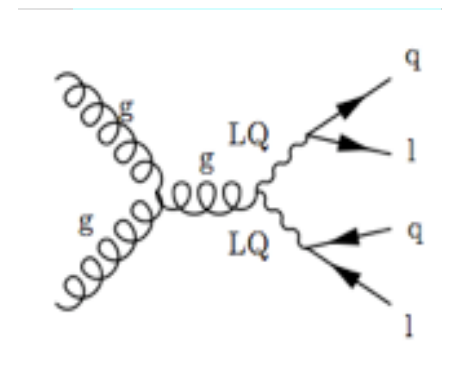
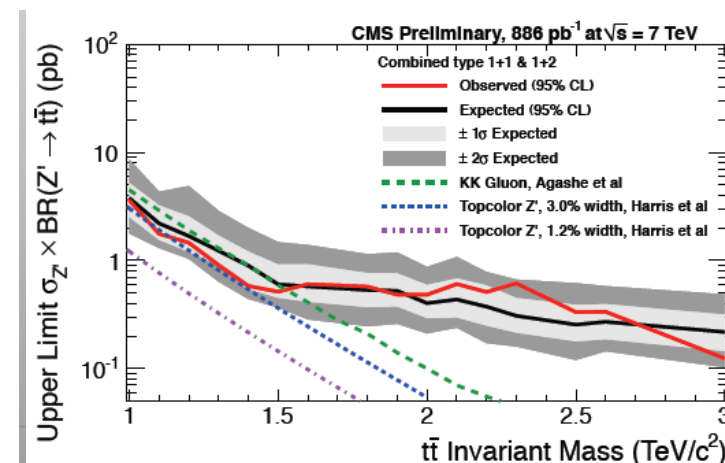
SUSY?



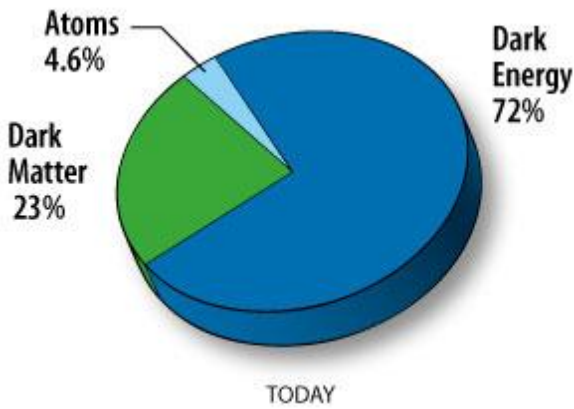
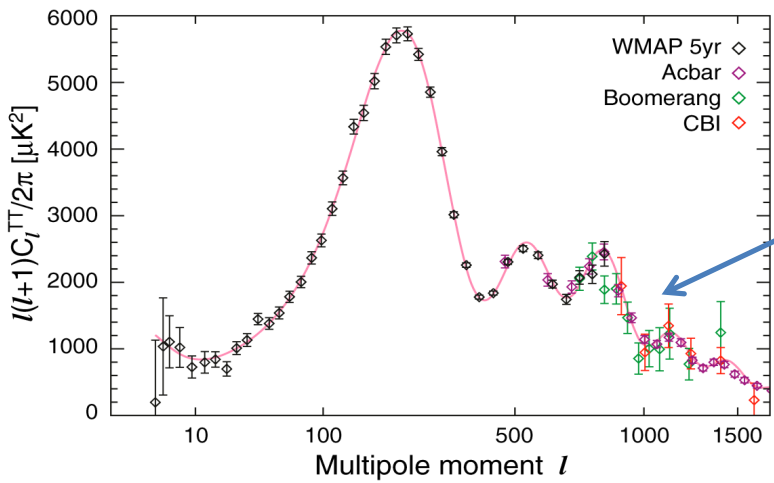
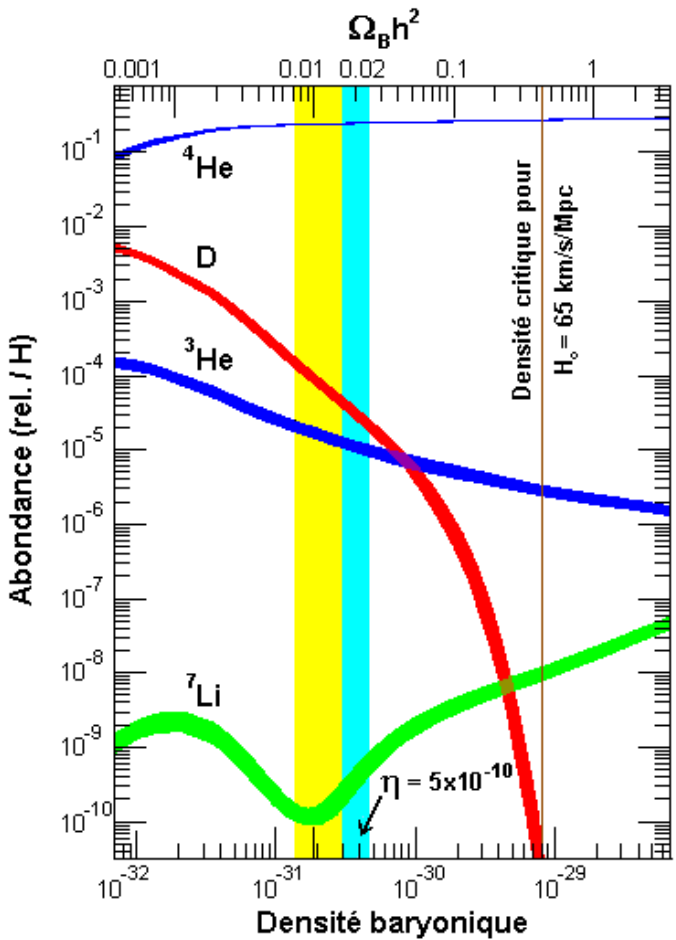
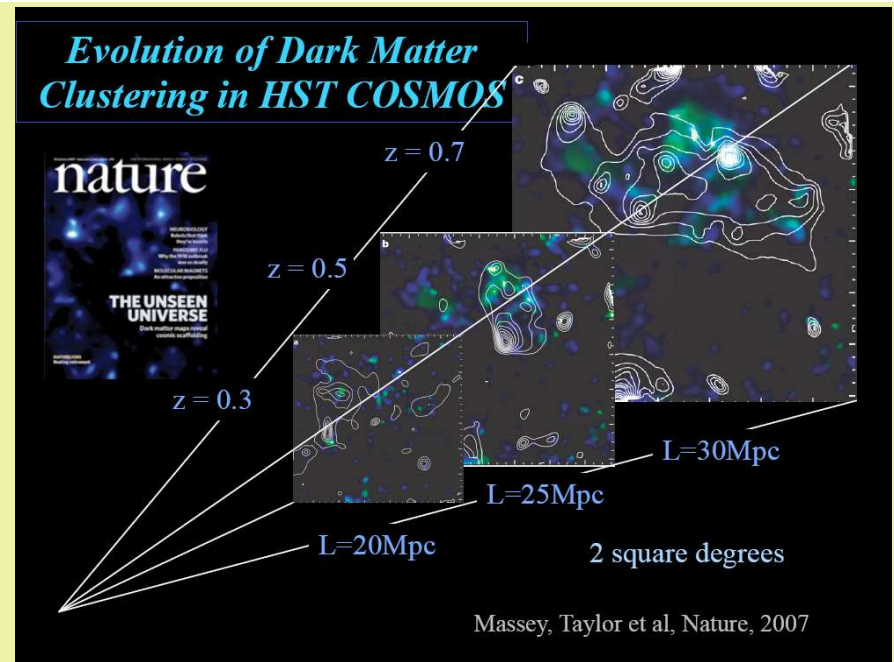
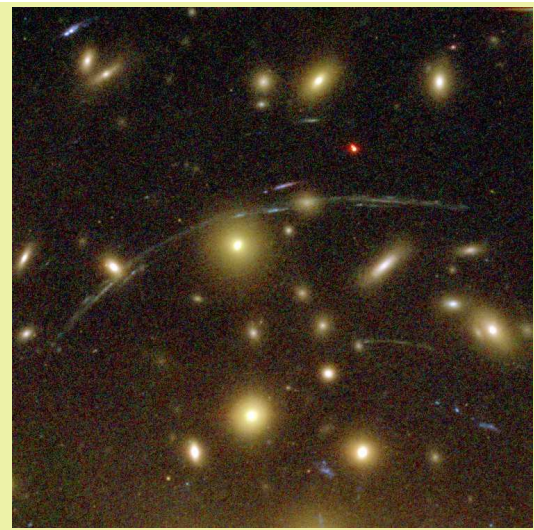
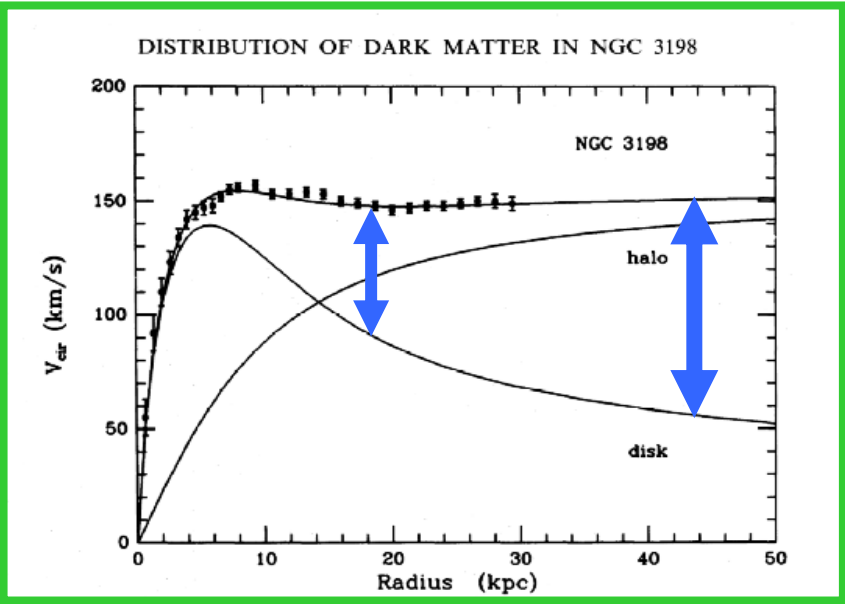
Higgs?



Others?



But there is definitely new physics (whatever it is...)



Two possible explanations...

Modifying gravity:

Bekenstein's theory (TeVeS)

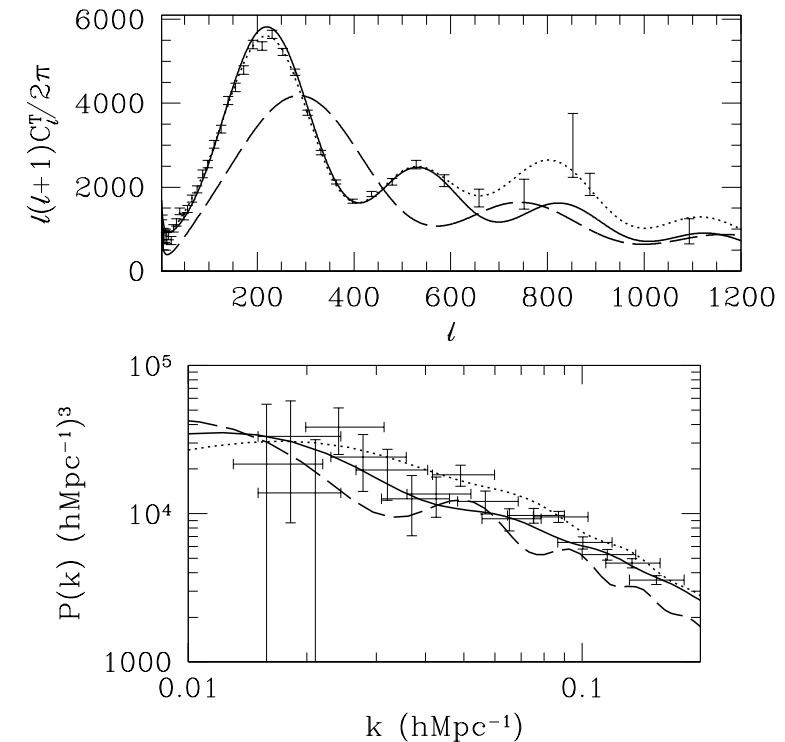


FIG. 4: The angular power spectrum of the CMB (top panel) and the power spectrum of the baryon density (bottom panel) for a MOND universe (with $a_0 \simeq 4.2 \times 10^{-8} \text{ cm/s}^2$) with $\Omega_\Lambda = 0.78$ and $\Omega_\nu = 0.17$ and $\Omega_B = 0.05$ (solid line), for a MOND universe $\Omega_\Lambda = 0.95$ and $\Omega_B = 0.05$ (dashed line) and for the Λ -CDM model (dotted line). A collection of data points from CMB experiments and Sloan are overplotted.

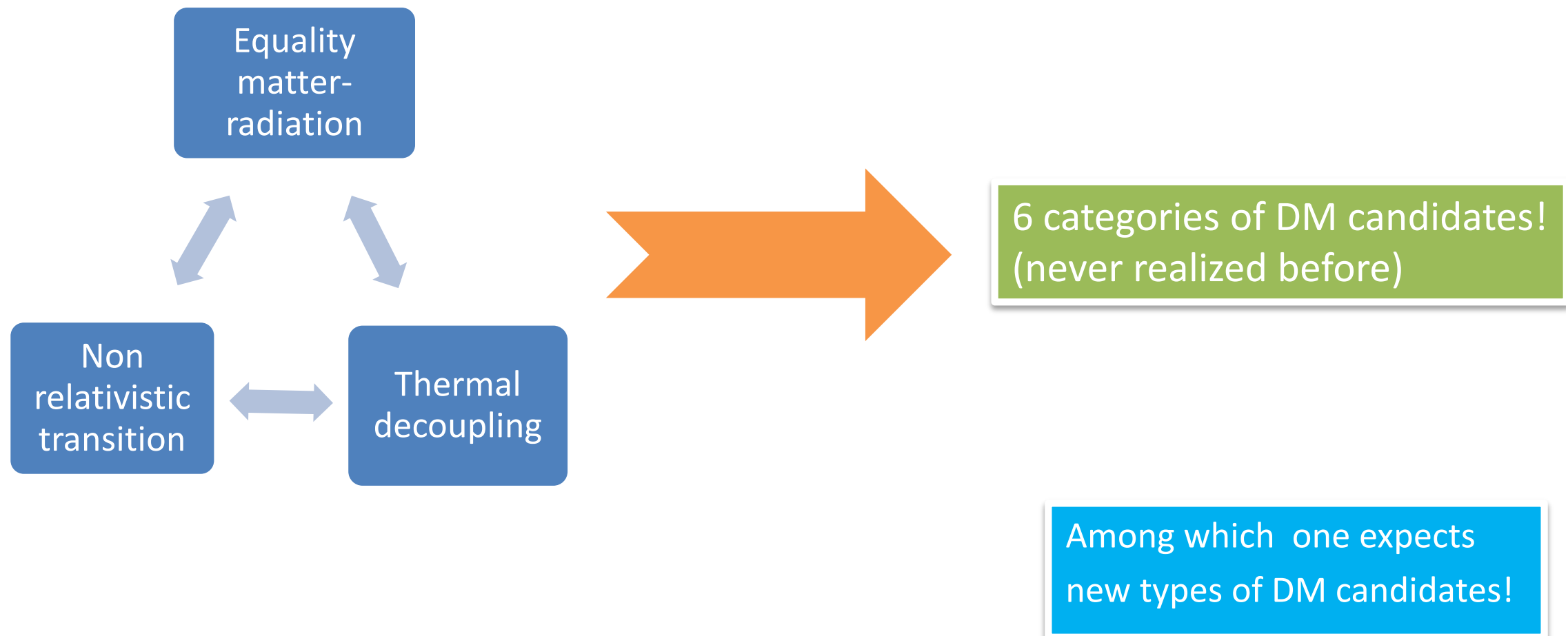
New neutral particles:

Hence “weakly” interacting when it comes to emitting photons but perhaps not that weakly interacting...

What kind of neutral particle?

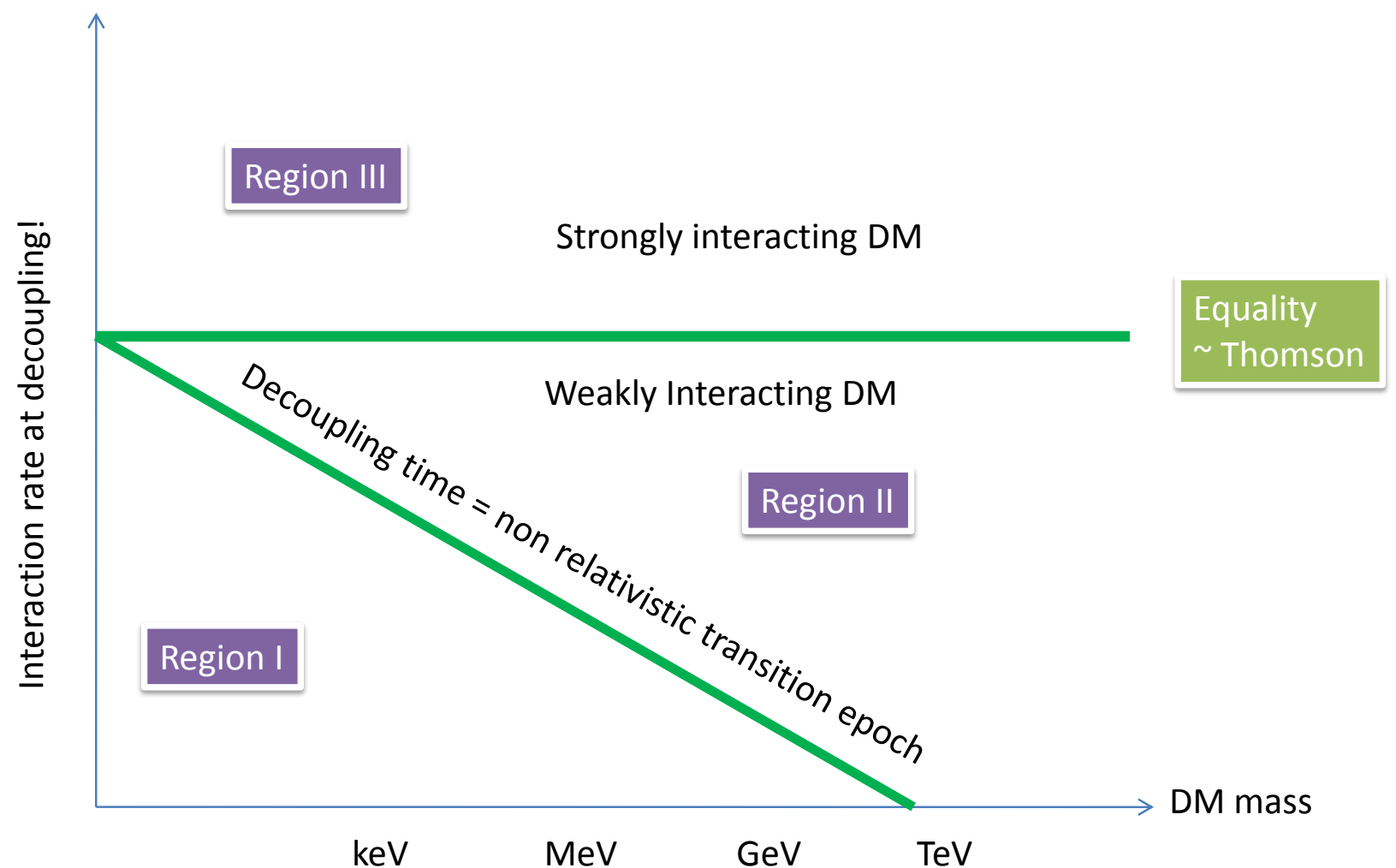
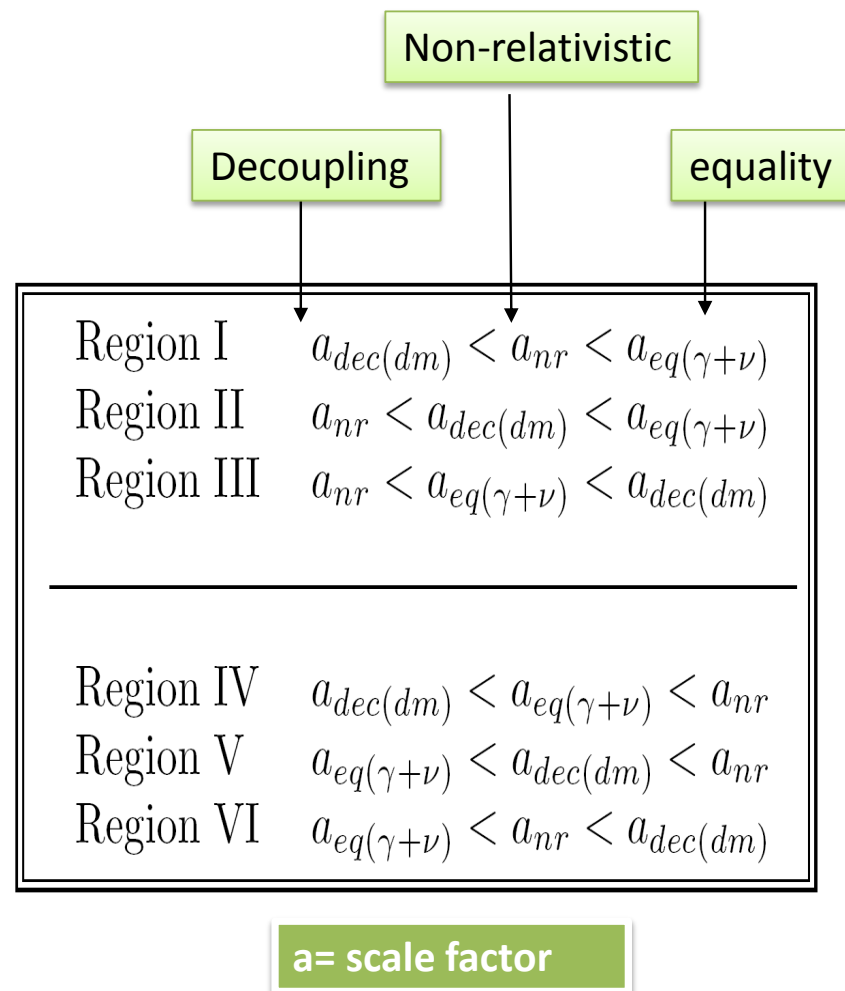
- * DM particles must have a mass
- * DM particles must have interactions
- * DM particles experience the thermal history of the Universe

- ✿ mass = non relativistic transition time
- ✿ interactions = decoupling time
- ✿ thermal history = equality time



Need for a neat classification of all types of DM candidates

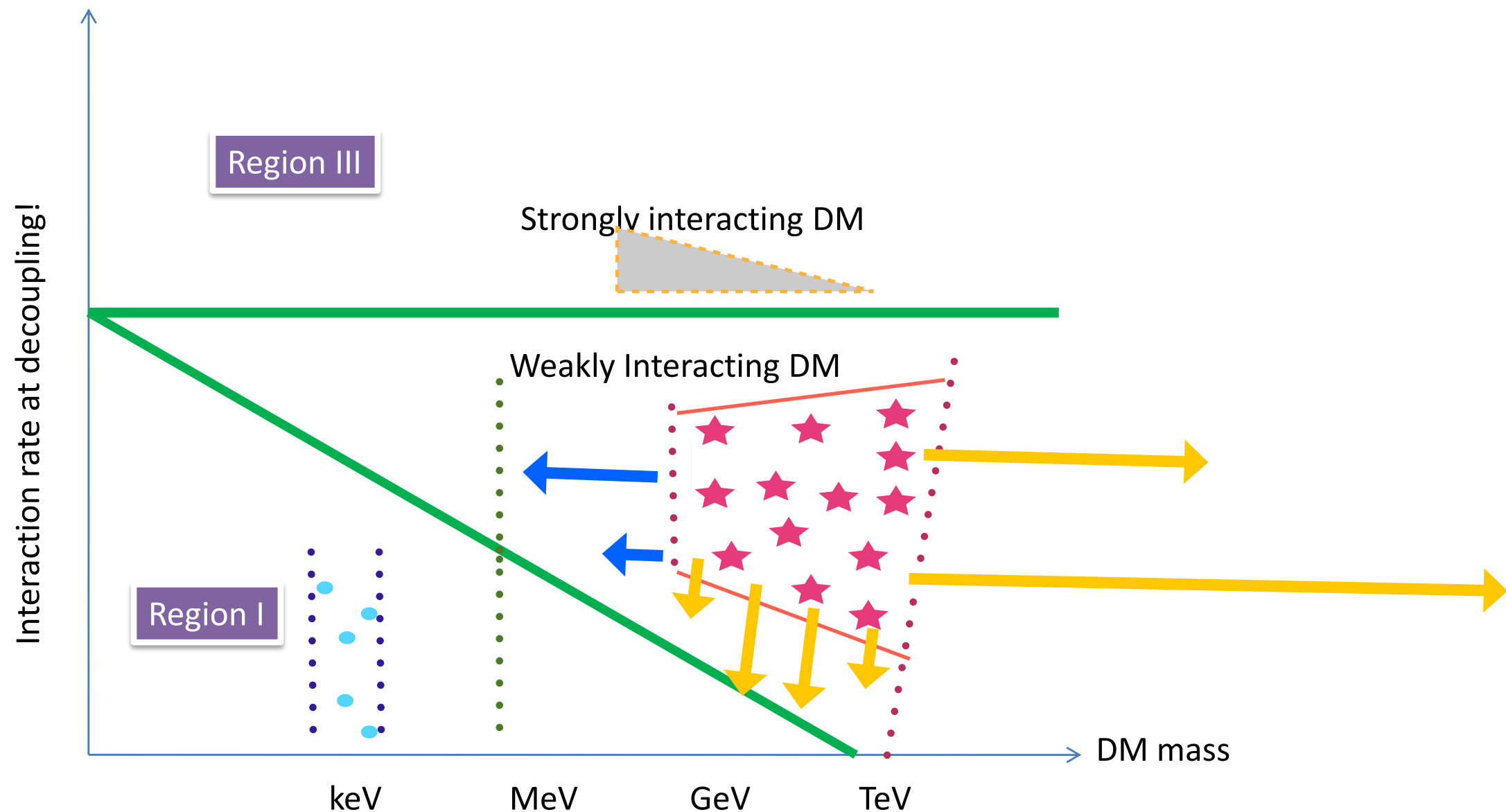
CB, P. Fayet and R. Schaeffer, astro-ph/0012504



Locate your favourite DM candidates

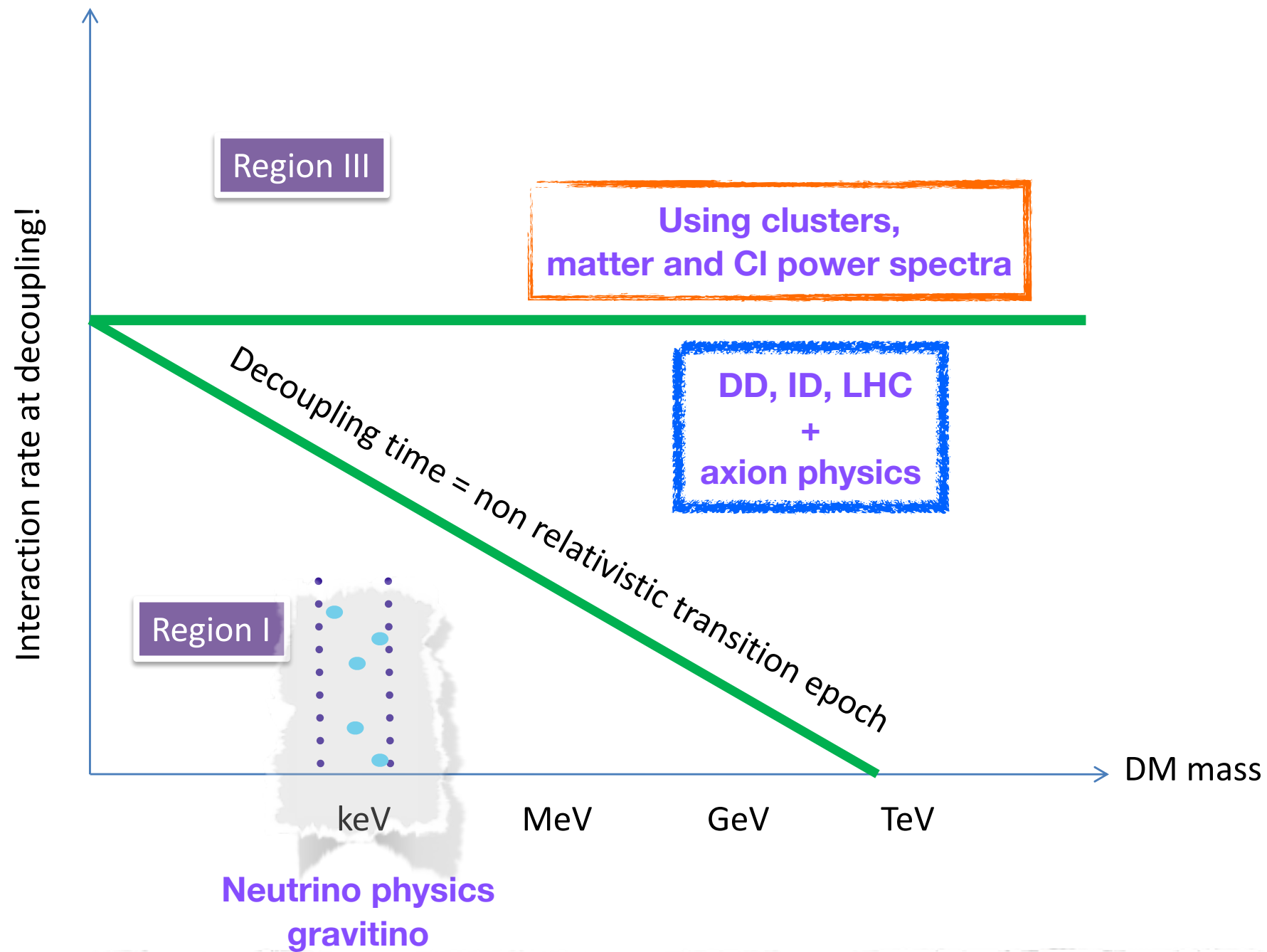
(forget about CDM and WDM for the moment)

Theoretical parameter space



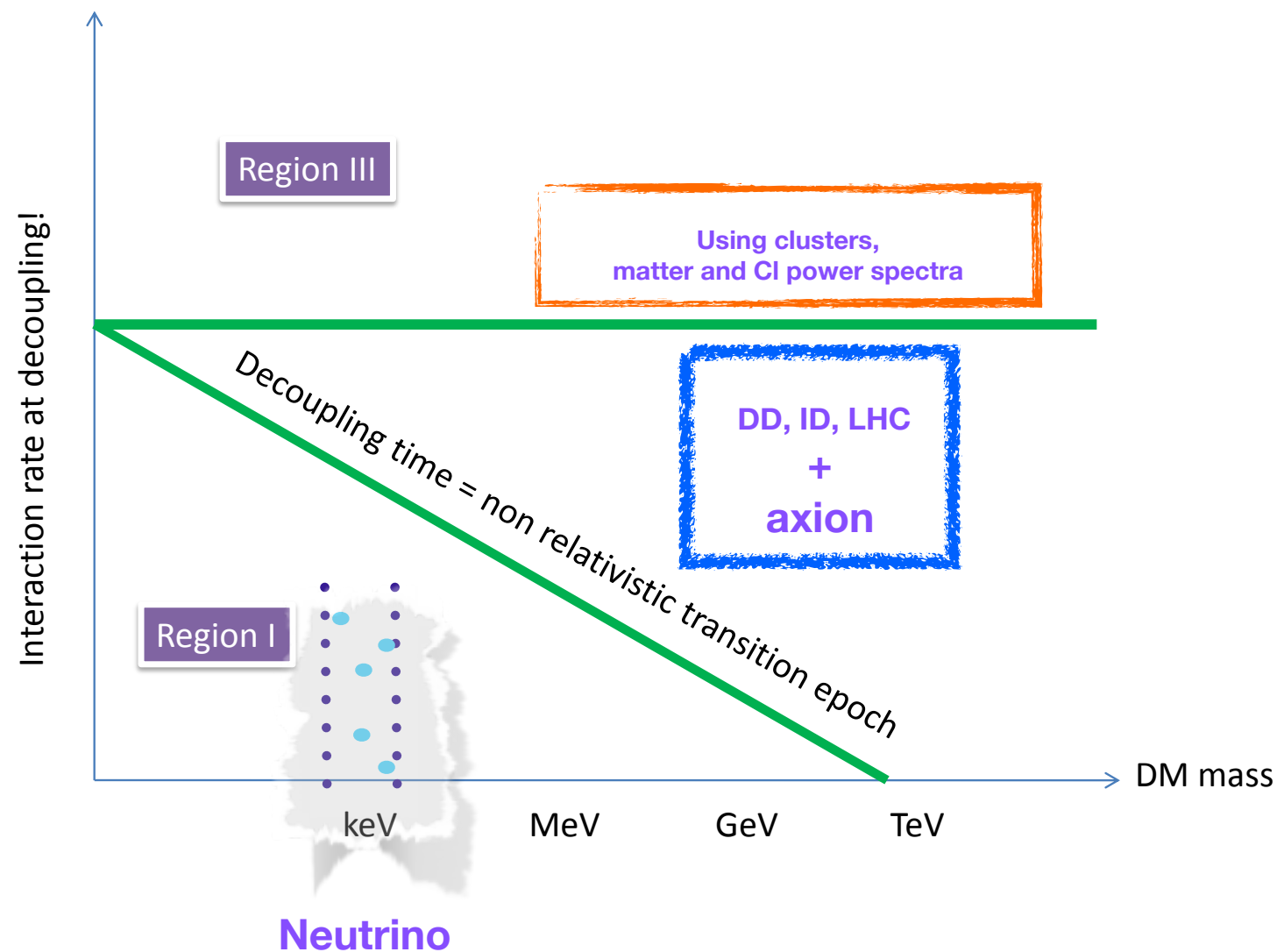
Theoretically, we have explored a small region of the parameter space.

DM parameter space investigated experimentally/observationally



Experimentally, explored about the same although a bit tricky

We have to be very LUCKY for the DM to be the type of particles that we can discover with our present tools!



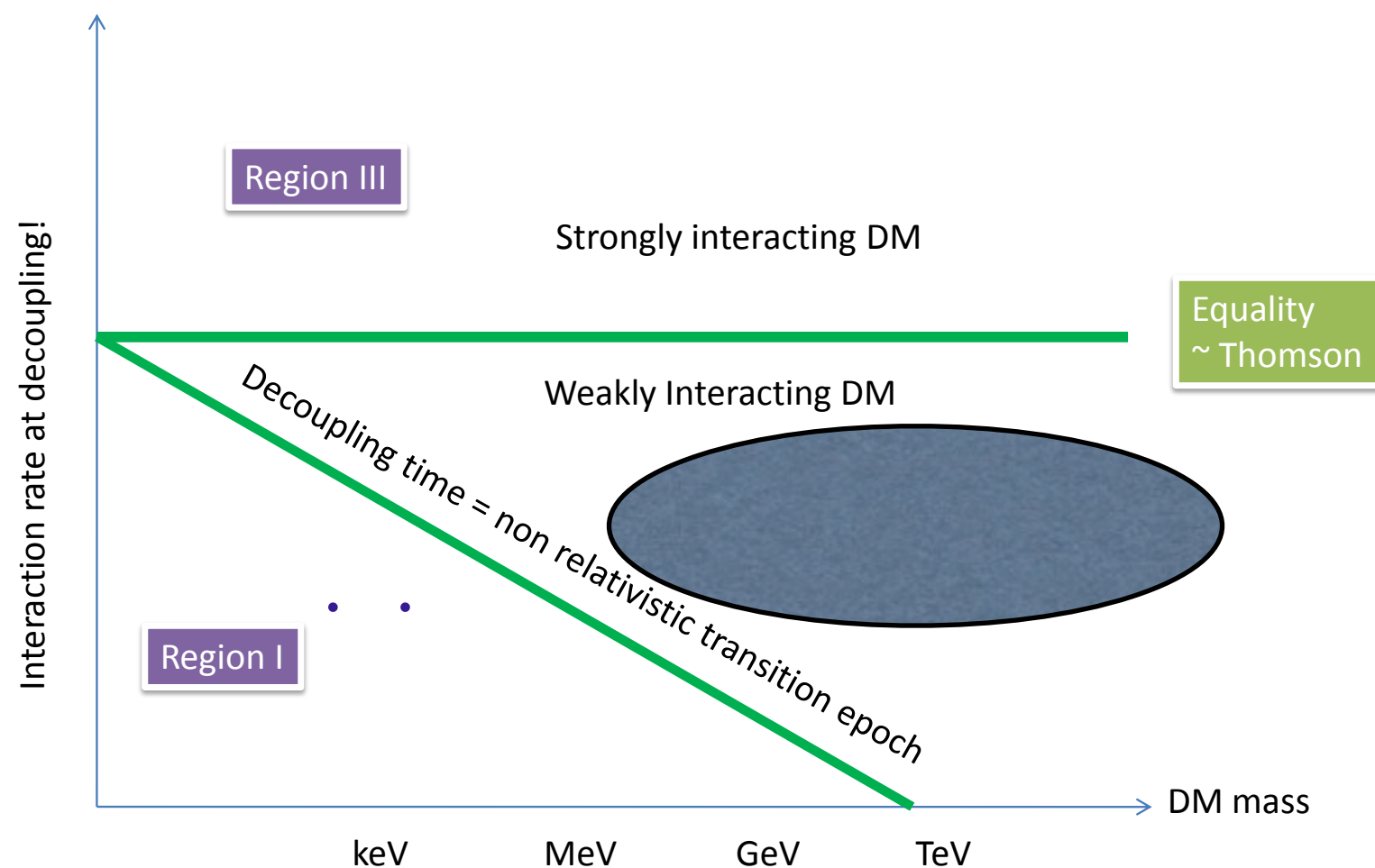
LHC

~~LEP~~ should ~~also~~ provide us with useful information relevant to grand unified theories, supersymmetry and dynamical symmetry

Experimentally, explored about the same although a bit tricky

The WIMP region

**‘thermal’ bulk only;
I am discarding non-thermal candidates**



The relic density paradigm

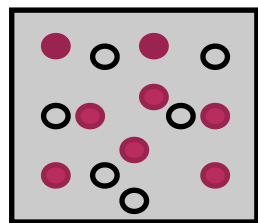
- Hypothesis until the 2000s, working hypothesis:

*thermal (annihilating)

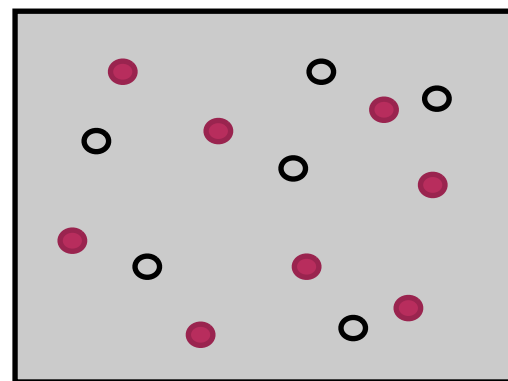
*symmetric ($n_{\text{DM}} = n_{\text{antiDM}}$)

*weakly interacting

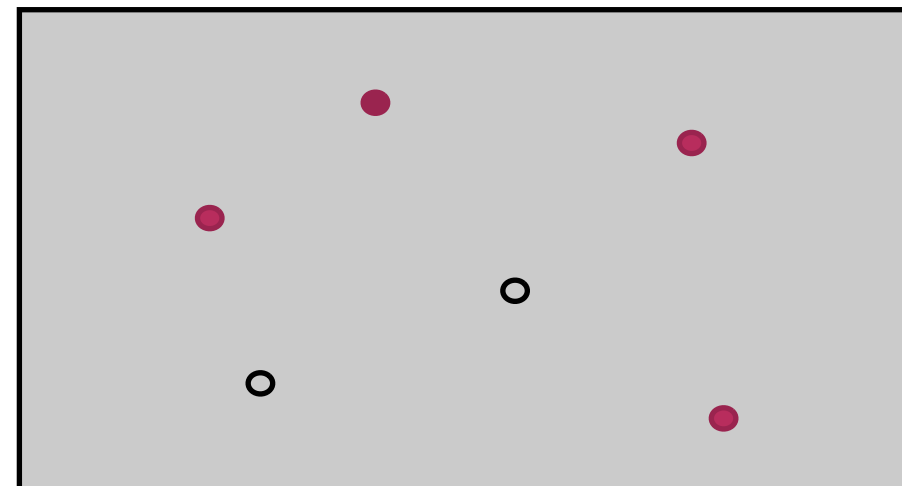
$$\frac{dn}{dt} = -3 H n - \langle \sigma v \rangle (n^2 - n_{eq}^2)$$



Thermal equilibrium



Expansion



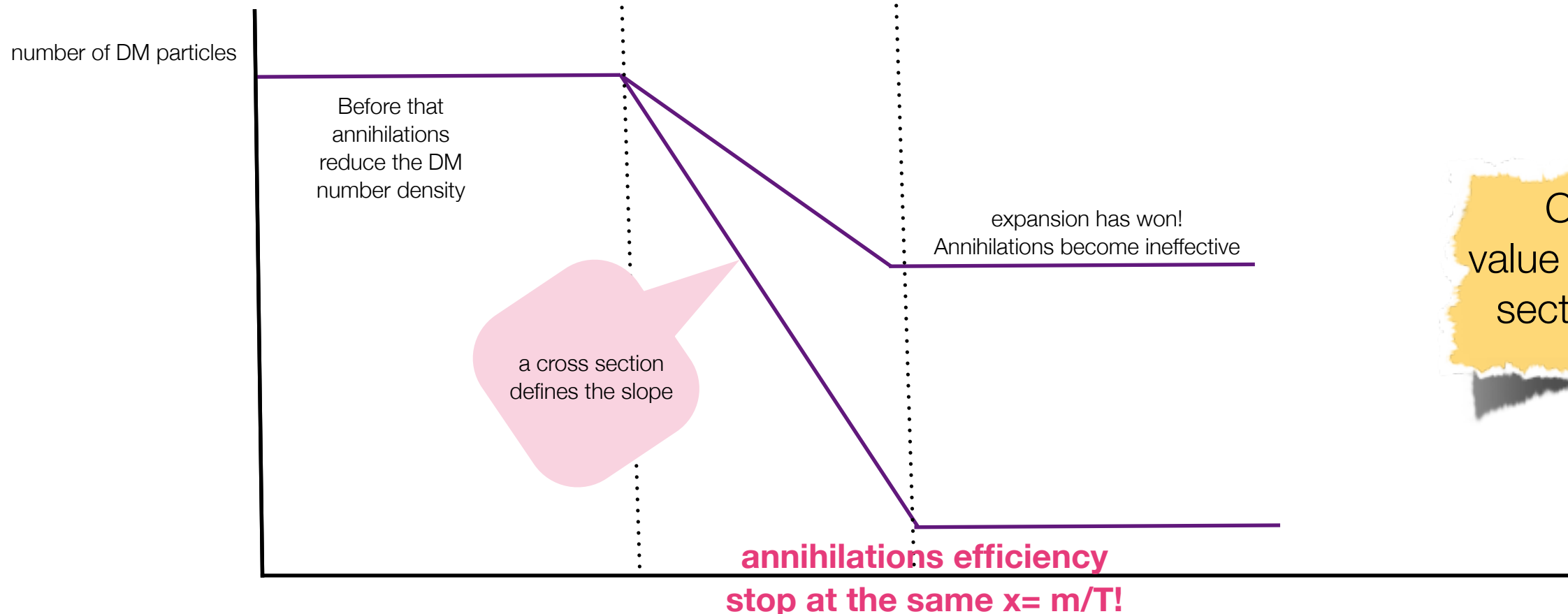
Annihilations

Boltzmann equation
no Particle Physics!

Why is/was the thermal hypothesis so popular?

- Solutions of the Boltzmann equation requires **NO** Particle Physics input !
- BUT ... it does give an information about DM particle properties:

$$\langle \sigma v \rangle = \frac{3 \cdot 10^{-27} \text{ cm}^3 / \text{s}}{\Omega_{dm} h^2}$$



Only one
value of the cross
section works!

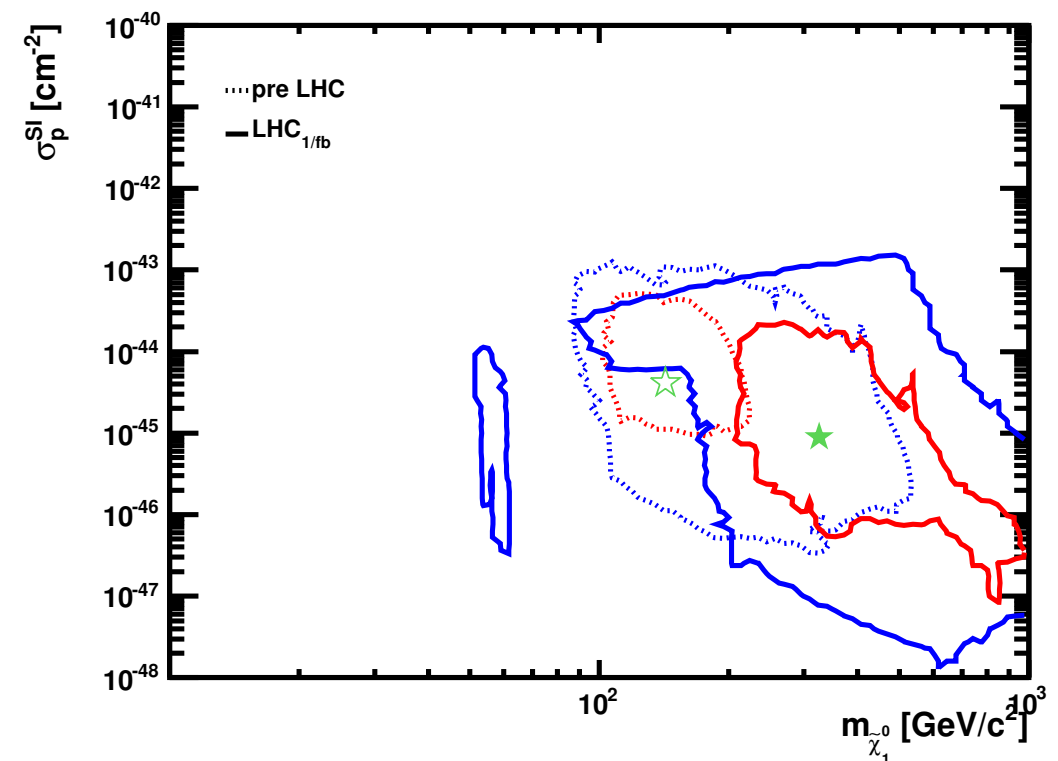
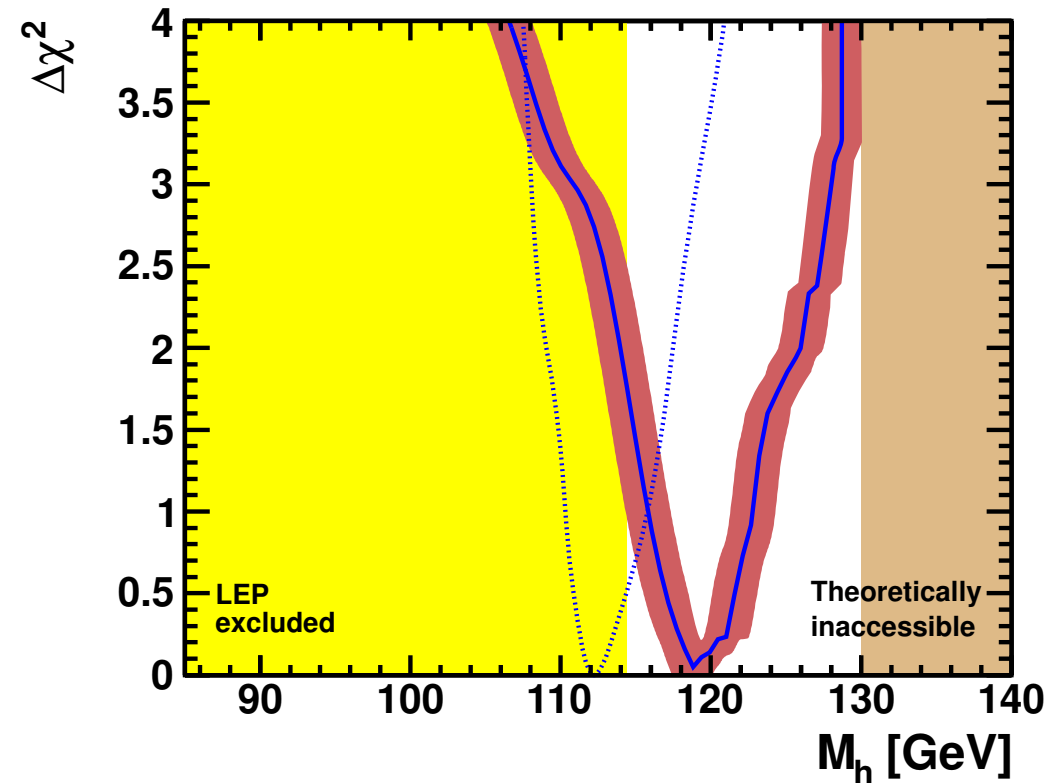
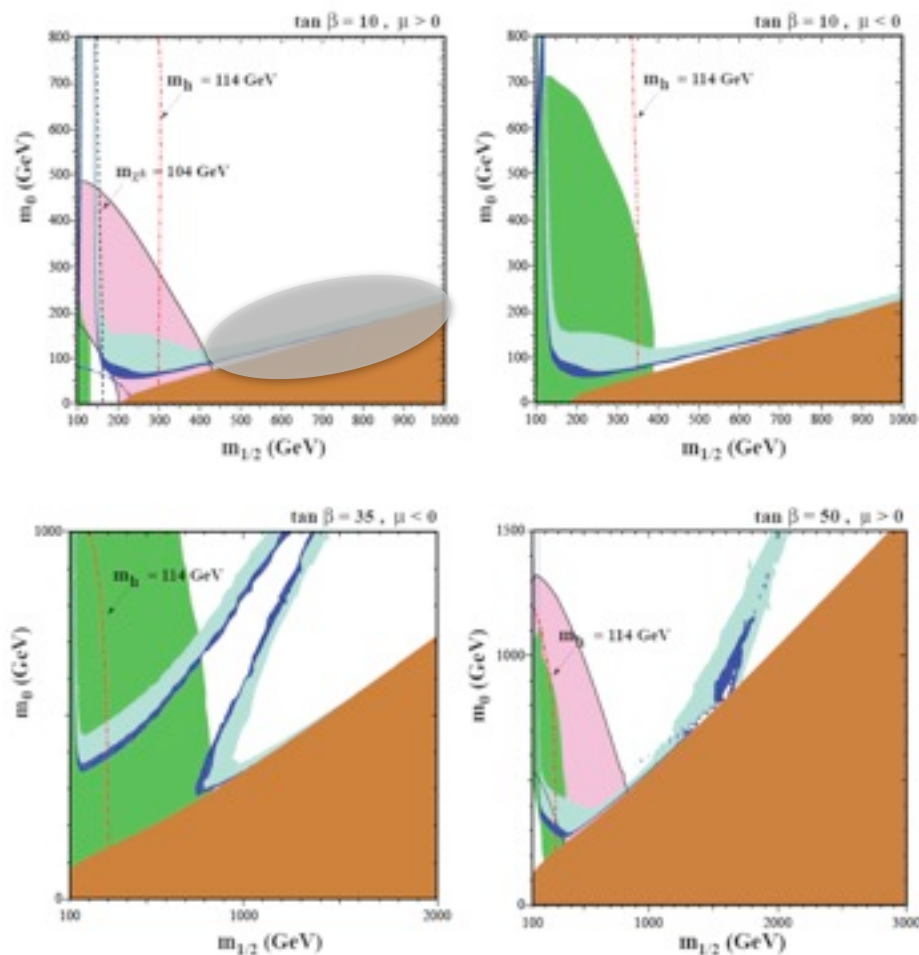
Will the CMSSM/mSUGRA survive?

Predictions are being tested at LHC: quid of the light Higgs?

arXiv:1110.3568

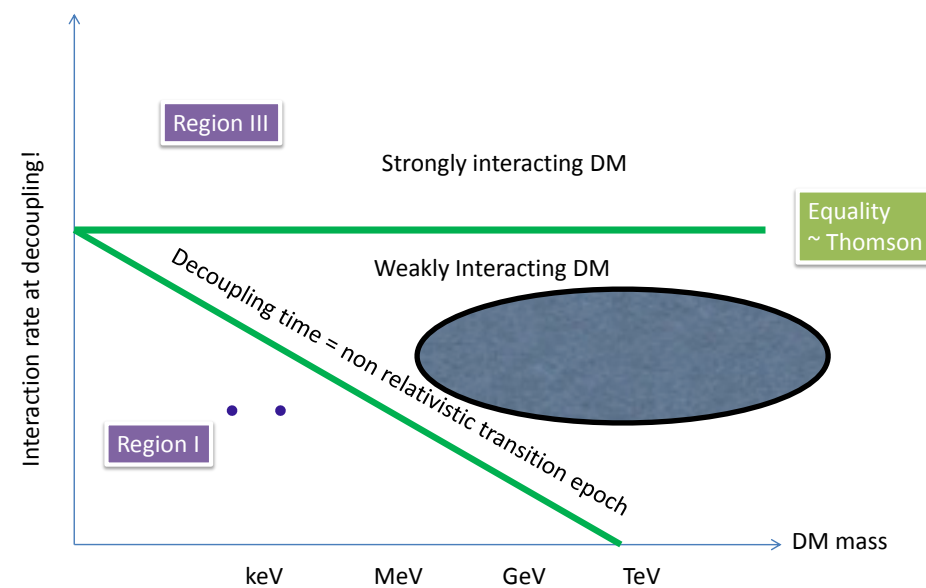
$$\frac{dn_1}{dt} = -3Hn_1 - \langle\sigma v\rangle_{ann}(n_1^2 - n_{1,0}^2) - \langle\sigma v\rangle_{coa}(n_1 n_2 - n_{1,0} n_{2,0})$$

$$\frac{dn_2}{dt} = -3Hn_2 - \langle\sigma v\rangle_{ann}(n_2^2 - n_{2,0}^2) - \langle\sigma v\rangle_{coa}(n_1 n_2 - n_{1,0} n_{2,0})$$



Summary for CMSSM

Likely to be excluded, we have to move on

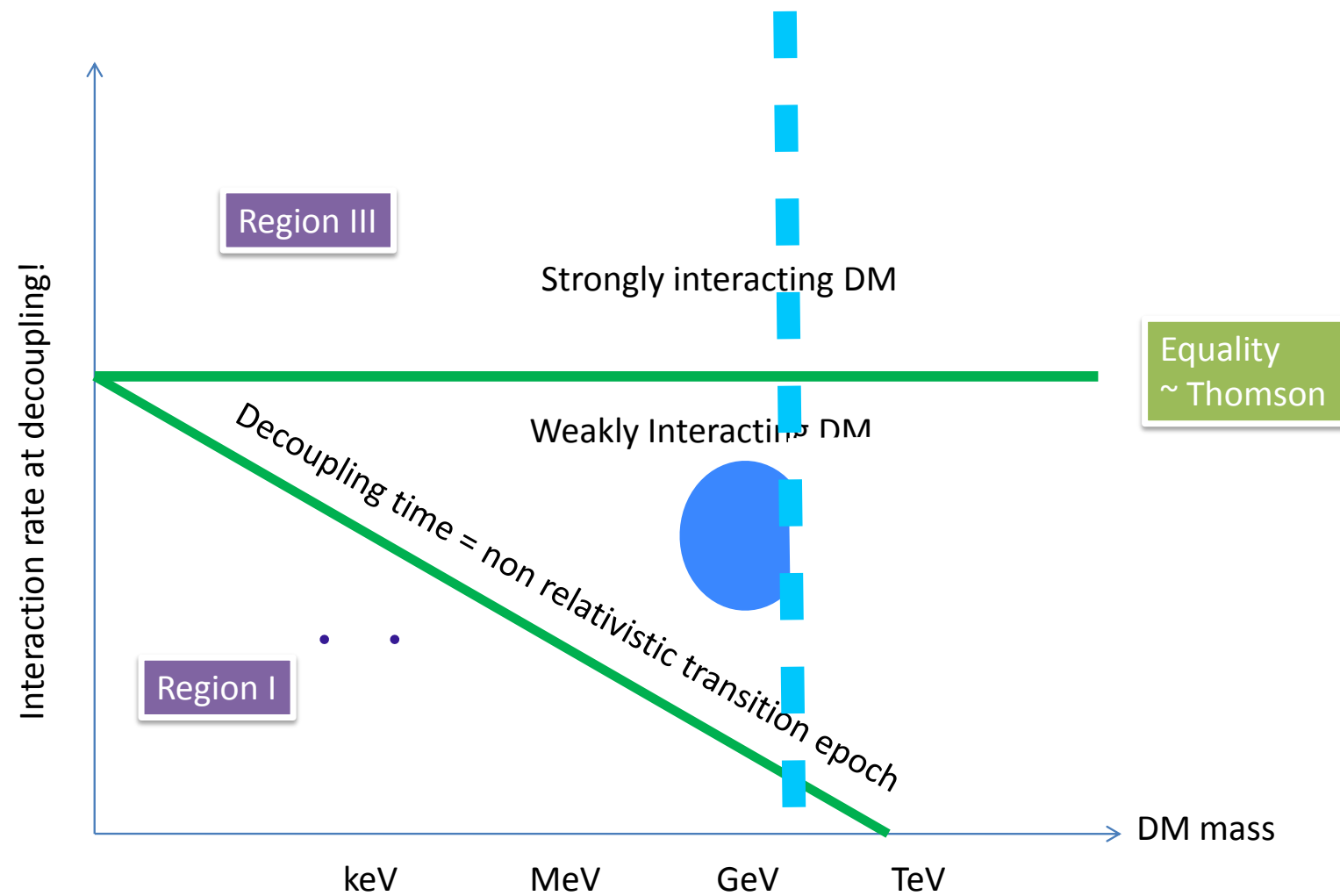


What (conventional) else in the thermal region?

(Can we still have supersymmetry?)

CMSSM $\longrightarrow$ pMSSM

Thermal, low mass, annihilating candidates



Are there light ([1-10] GeV) neutralinos in the pMSSM?

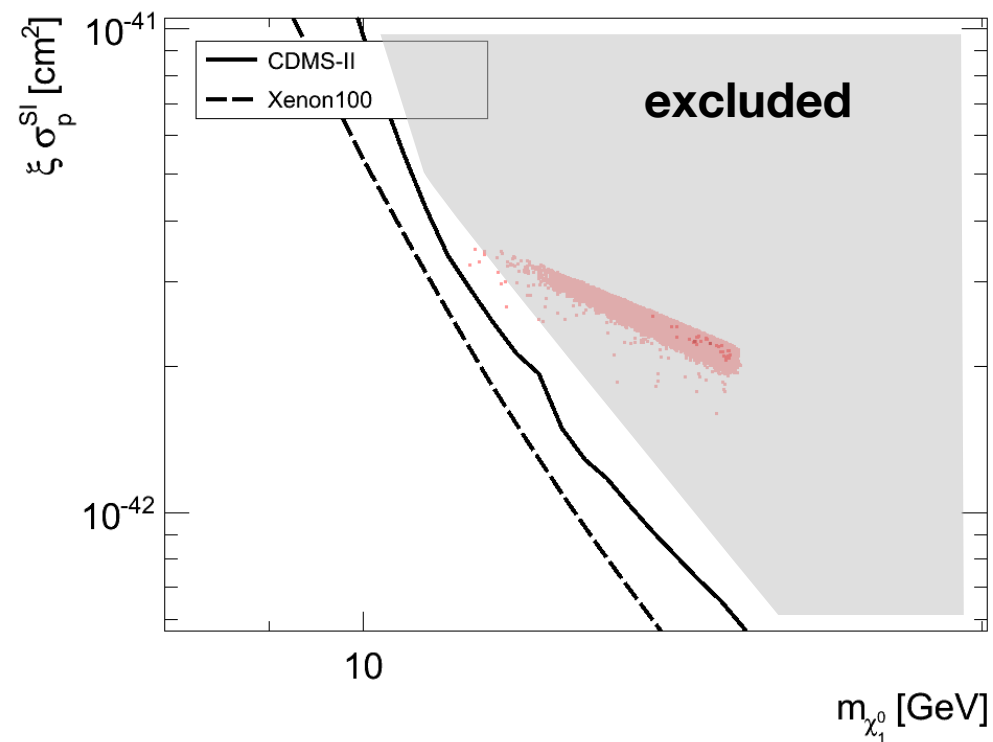


FIG. 3: MSSM-EWSB scenario with $\mu > 0$ and $m_\chi < 15$ GeV. Spin-independent cross section on proton times the fraction of neutralinos in the Milky Way dark halo (ξ) versus the neutralino mass m_χ . The dark red (light pink) points have a likelihood greater than 99.4% (68%).

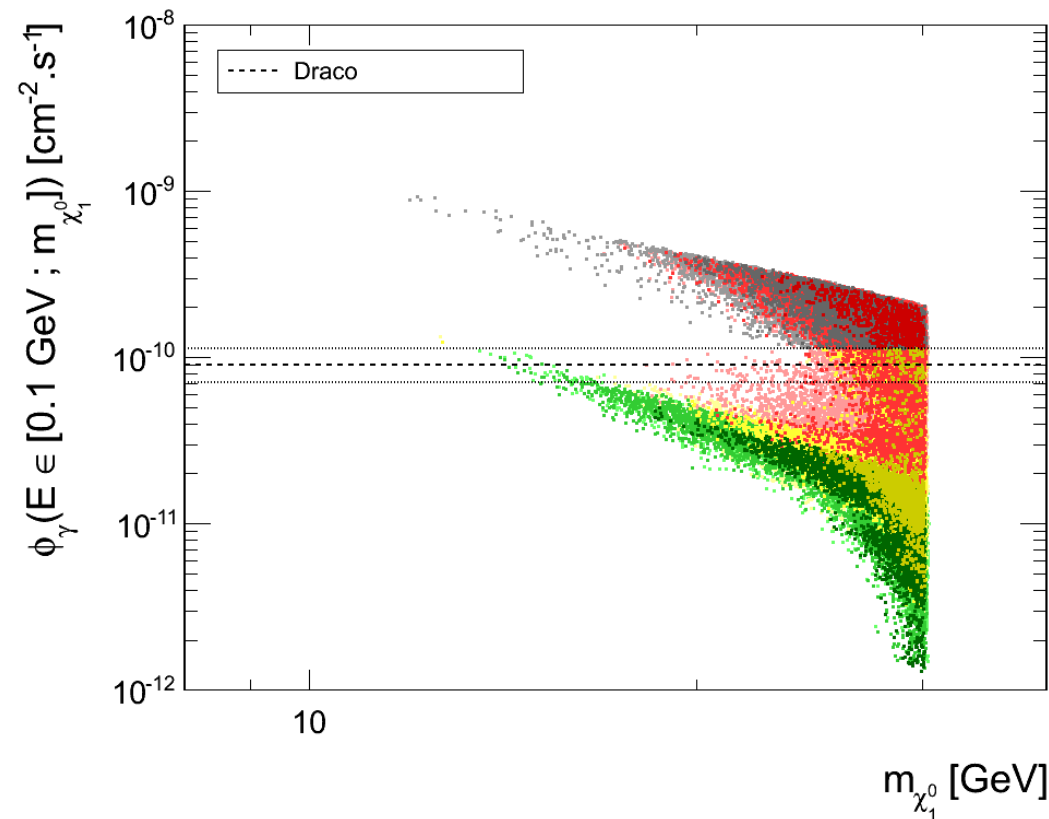
$$100\% \Omega_{WMAP} h^2 > \Omega_\chi h^2 > 10\% \Omega_{WMAP} h^2,$$

constraint	value/range	tolerance	applied
S_{masses}	-	none	both
$\Omega_{WMAP} h^2$	0.01131 - 0.1131	0.0034	both
$(g-2)_\mu$	$25.5 \cdot 10^{-10}$	stat: $6.3 \cdot 10^{-10}$ sys: $4.9 \cdot 10^{-10}$	both
$\Delta\rho$	≤ 0.002	0.0001	MSSM
$b \rightarrow s\gamma$	$3.52 \cdot 10^{-4}$ [38, 39]	th: $0.24 \cdot 10^{-4}$ exp: $0.23 \cdot 10^{-4}$	both
$B_s \rightarrow \mu^+ \mu^-$	$\leq 4.7 \cdot 10^{-8}$	$4.7 \cdot 10^{-10}$	both
$R(B \rightarrow \tau\nu)$	1.28 [38]	0.38	both
m_H	≥ 114.4	1%	MSSM
$Z \rightarrow \chi_1 \chi_1$	≤ 1.7 MeV	0.3 MeV none	MSSM NMSSM
$e^+ e^- \rightarrow \chi_1 \chi_{2,3}$	≤ 0.1 pb [40]	0.001 pb none	MSSM NMSSM
ΔM_s	$117.0 \cdot 10^{-13}$ GeV	th: $21.1 \cdot 10^{-13}$ GeV exp: $0.8 \cdot 10^{-13}$ GeV	NMSSM
ΔM_d	$3.337 \cdot 10^{-13}$ GeV	th: $1.251 \cdot 10^{-13}$ GeV exp: $0.033 \cdot 10^{-13}$ GeV	NMSSM

Theoretically: yes,
Experimentally: no (?)

Very light MSSM neutralinos are out but that is not the end of the story (...yet!...)

1108.1338



Parameter	Minimum	Maximum	Tolerance
M_1	1	1000	3
M_2	100	2000	30
M_3	500	6500	10
μ	0.5	1000	0.1
$\tan\beta$	1	75	0.01
M_A	1	2000	4
A_t	-3000	3000	100
$M_{\tilde{t}_R}$	70	2000	15
$M_{\tilde{t}_L}$	70	2000	15
$M_{\tilde{q}_{1,2}}$	300	2000	14
$M_{\tilde{q}_3}$	300	2000	14

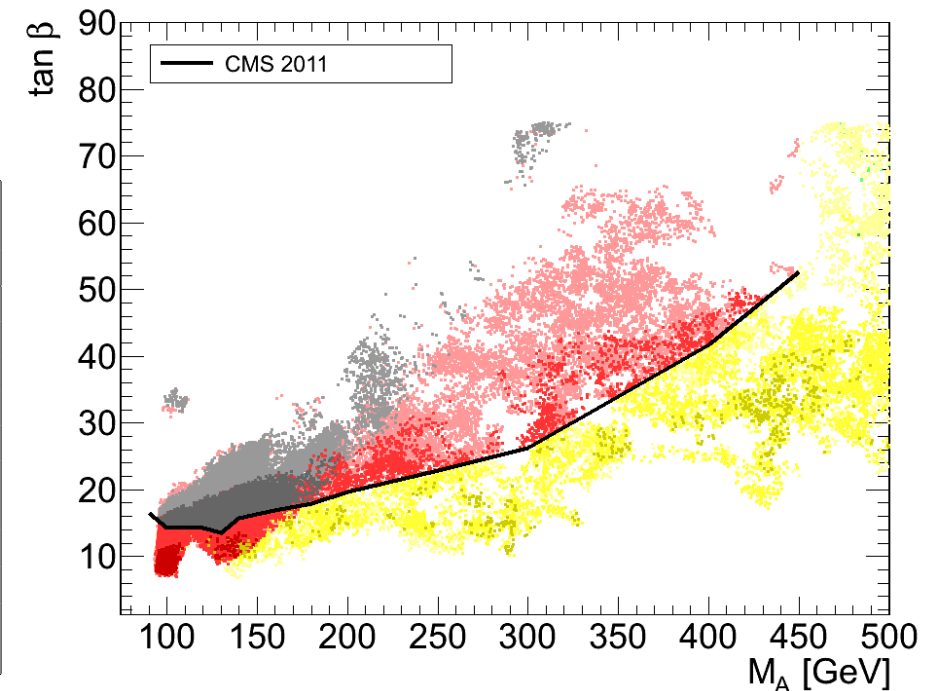
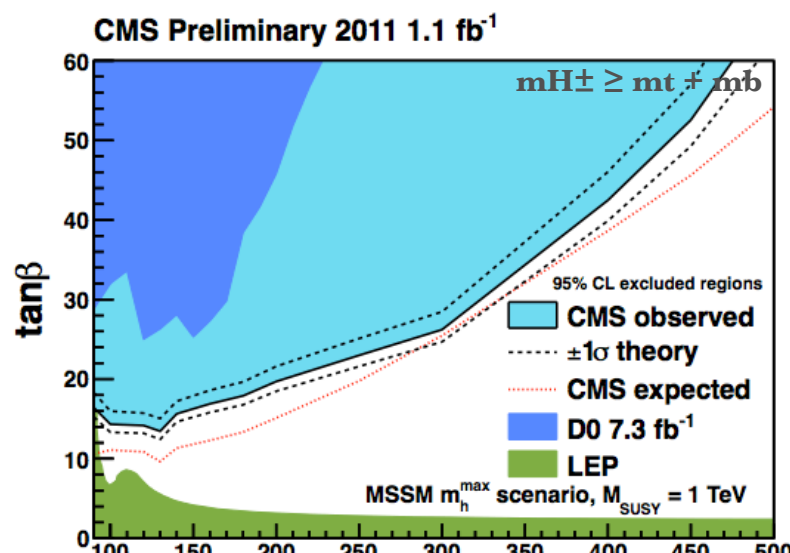


FIG. 2: Allowed points in the $\tan\beta$ vs. M_A plane in the $m_{\chi_1^0} < 30$ GeV search. We show only the region where $M_A < 500$ GeV. The exclusion limit from CMS is also displayed. In yellow (red), points excluded by one (two) constraint and in black those excluded by three constraints (CMS, XENON100 and dSph as described in section III A). The shading represents Q : weights of darker points are at most at 1σ from $Q_{\max}$ while the lighter points are at most at 2σ and 3σ .



Detectability at the remaining points?

(Can light neutralino escape constraints?)

Complementarity of the 3 types of experiments

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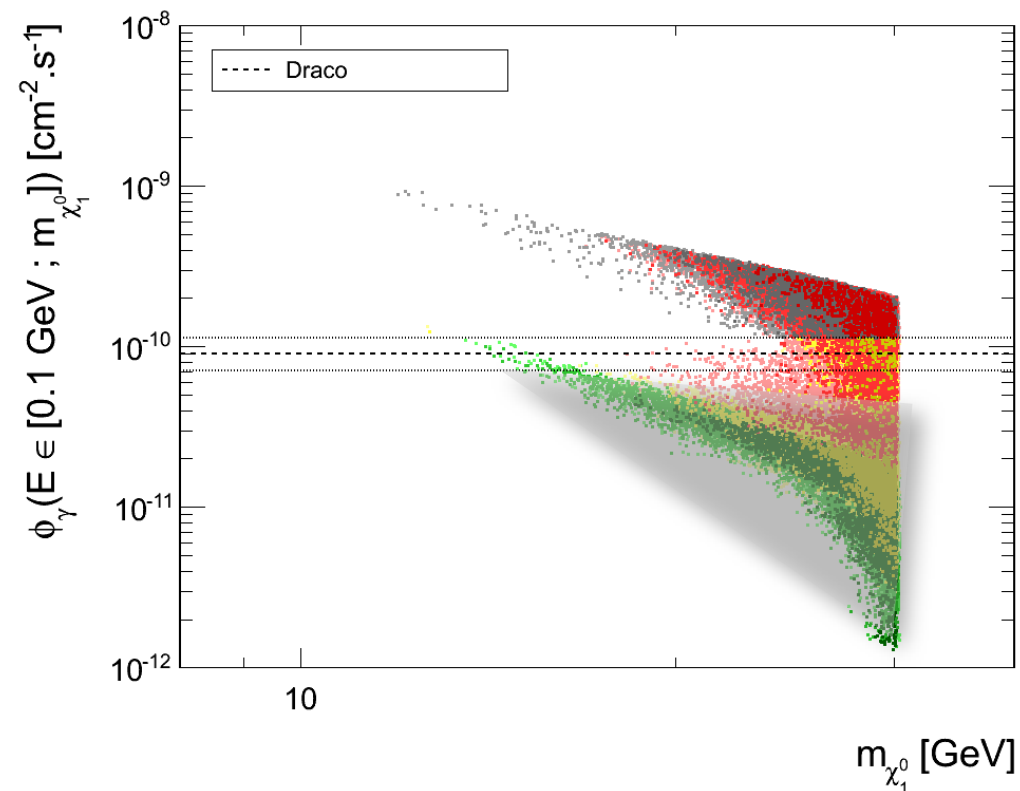


FIG. 5: Integrated γ -ray flux from the Draco dwarf spheroidal galaxy as a function of the neutralino mass in the $m_{\chi_1^0} < 30$ GeV search. We show limits from Fermi-LAT. Same color code as Fig. 4.

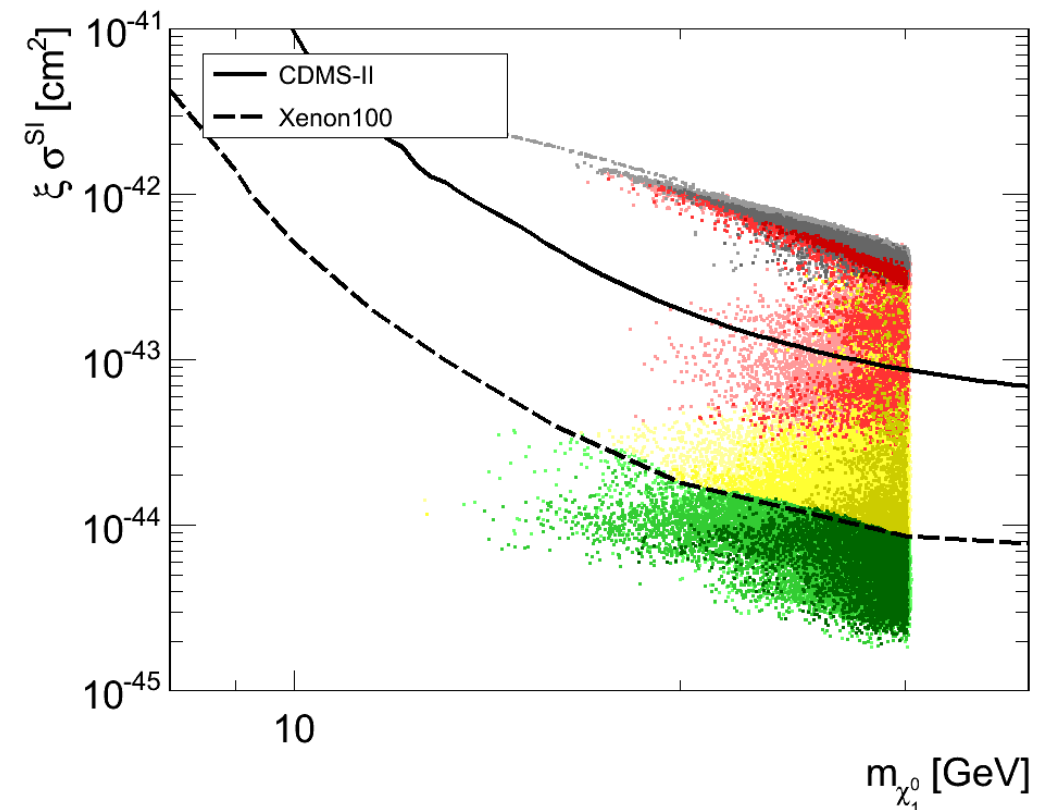
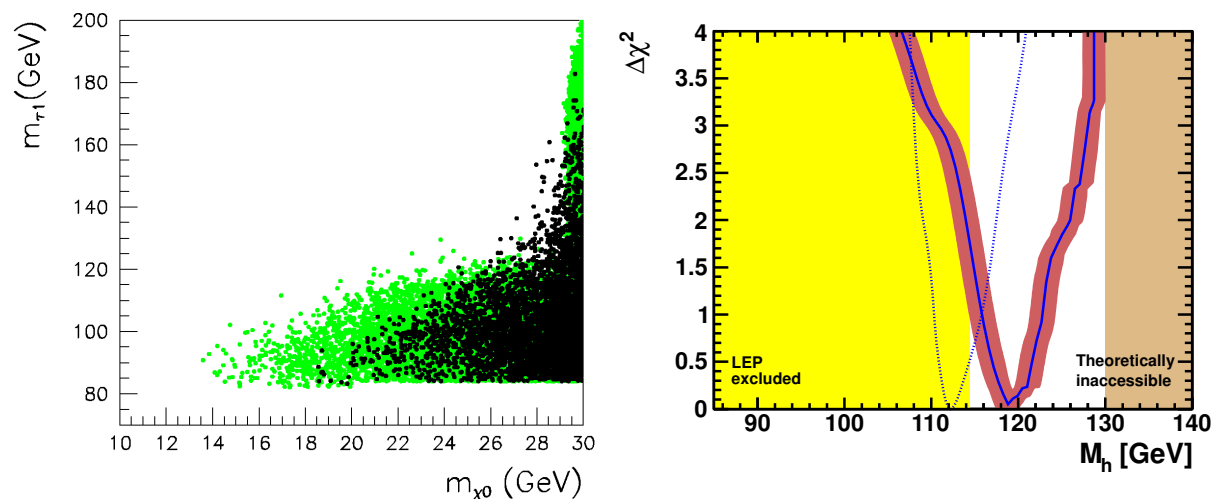
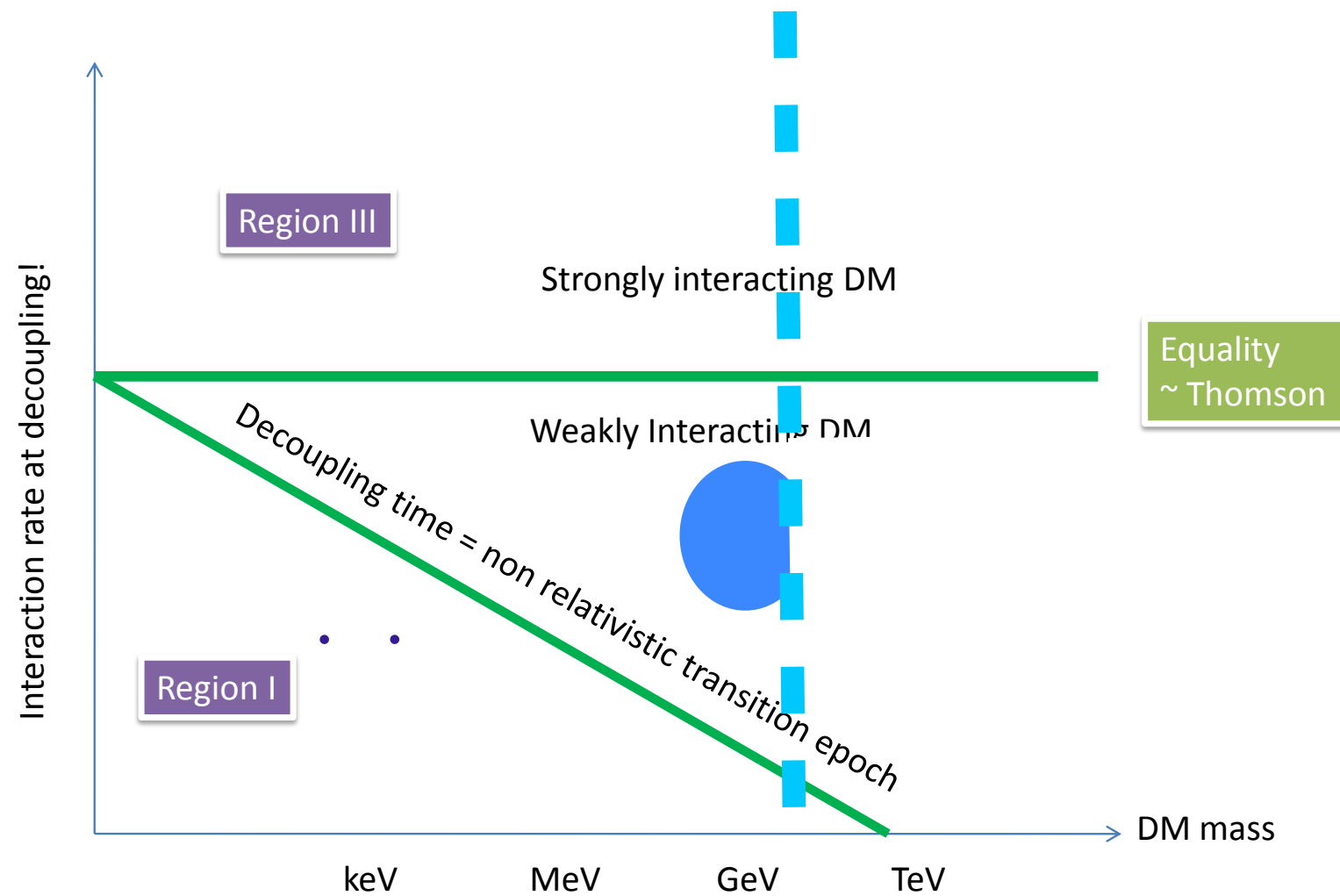


FIG. 4: Points of the $m_{\chi_1^0} < 30$ GeV search represented in the $\xi\sigma^{SI}$ vs. neutralino mass plane. Exclusion limits from CDMS-II [29] and XENON100 are shown. The color code is the same as in Fig. 2, green points are allowed.

**Light MSSM neutralinos should
not resist a very long time...**

**LHC + Direct Detection+Indirect detection:
the perfect “gun” machine?**

Light pMSSM neutralinos? not so sure!



Another SUSY example: NMSSM

1107.1614

fixed at the weak scale. The free parameters are taken to be the gaugino masses $M_1, M_2 = M_3/3$, the Higgs sector parameters $\mu, \tan\beta, \lambda, \kappa, A_\lambda, A_\kappa$, a common mass for the sleptons $m_{\tilde{l}}$ and the squarks $m_{\tilde{q}}$ as well as only one non-zero trilinear coupling, A_t , for more details see [4].

$$W = \lambda S H_u H_d + \frac{1}{3} \kappa S^3 \quad \mu = \lambda \langle S \rangle$$

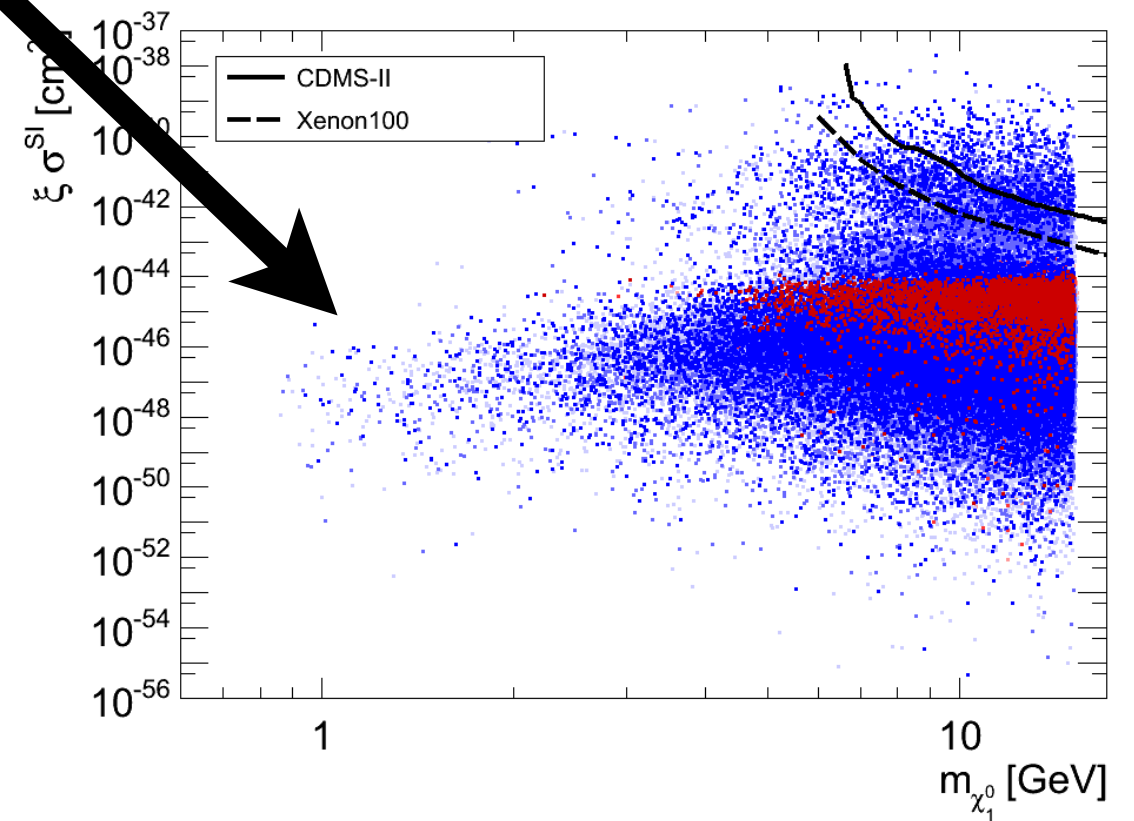
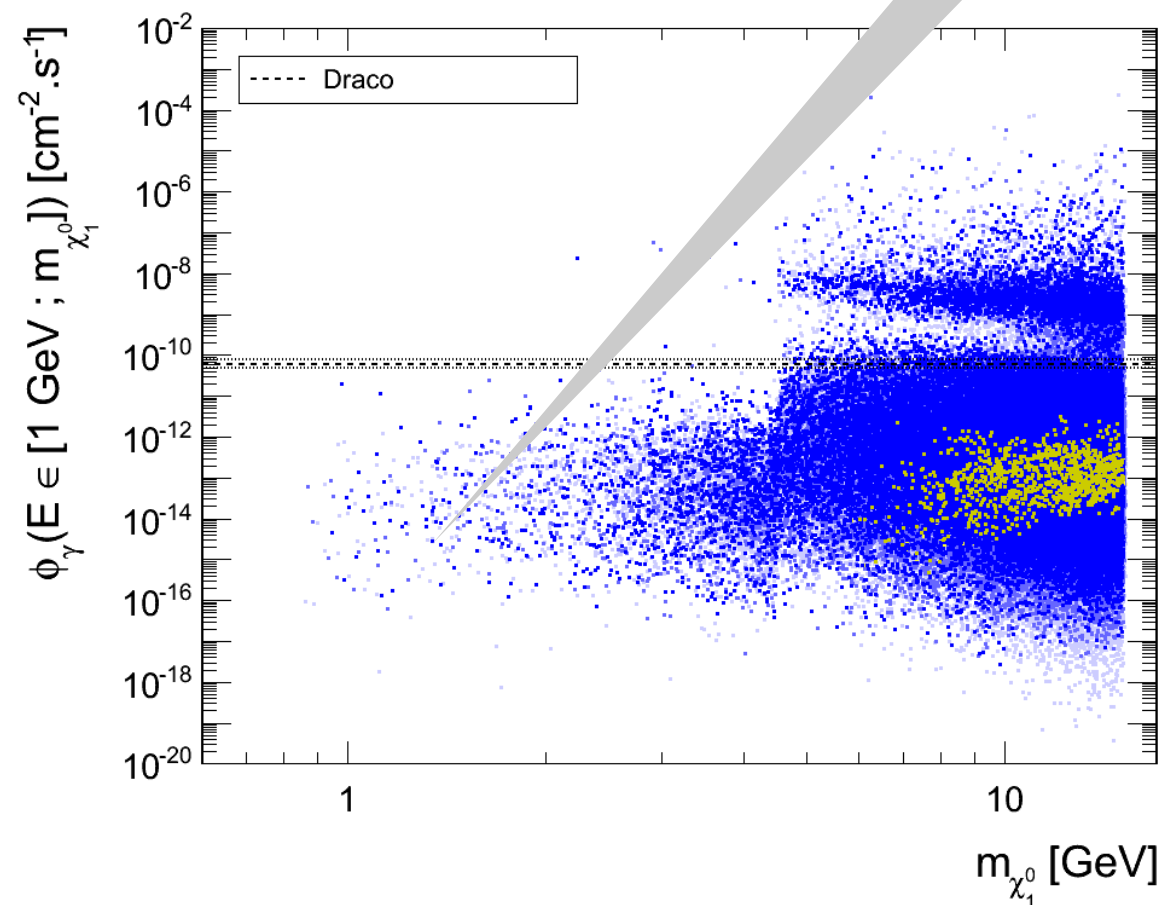
$$\mathcal{L}_{\text{soft}} = m_{H_u}^2 |H_u|^2 + m_{H_d}^2 |H_d|^2 + m_S^2 |S|^2$$

$$+ (\lambda A_\lambda H_u H_d S + \frac{1}{3} \kappa A_\kappa S^3 + h.c.)$$

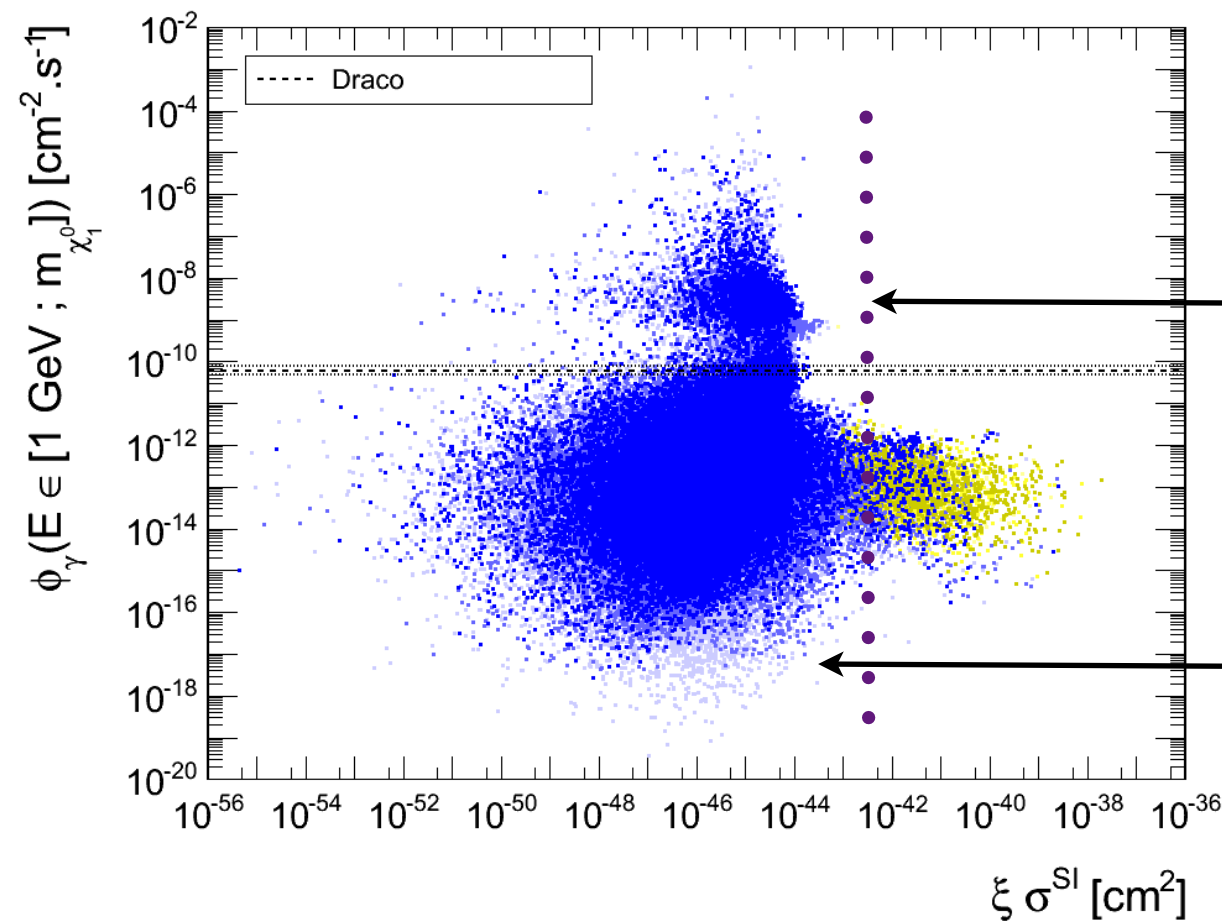
$\mu, \tan\beta$ as well as $\lambda, \kappa, A_\lambda, A_\kappa$.

Here we are, light neutralinos...

But...



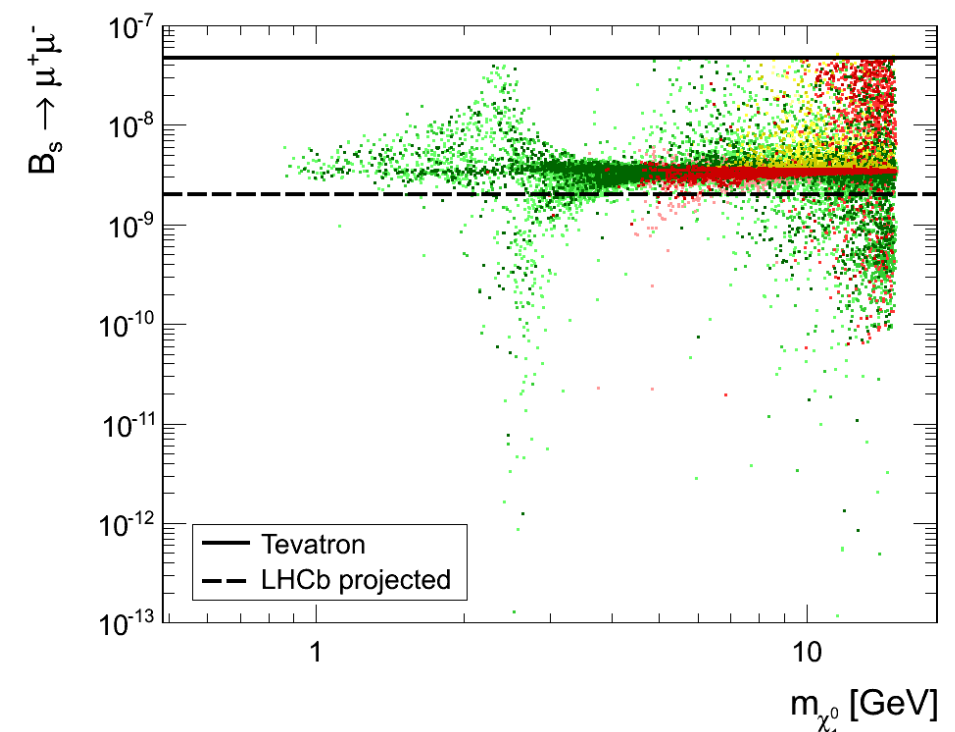
Light NMSSM neutralinos may be hard to discover and could be a solution if no new physics at LHC but easy to exclude if no Higgs!



Direct Detection

But LHC might help...

Higgs...



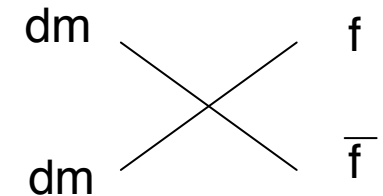
What (unconventional) else in the thermal region?

(Did we miss something important?)

The mass limit dogma

Hut-Lee&Weinberg (massive neutrinos):

$$\sigma v \propto \frac{m_{dm}^2}{m_w^4}$$



general assumptions

❖ DM is a fermion

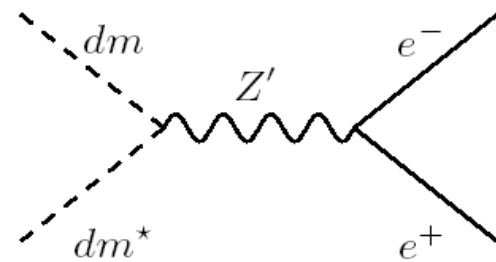
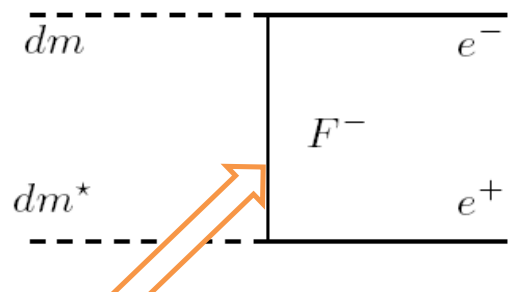
❖ A heavy particle is exchanged

$$\sigma v \propto \frac{m_{dm}^2}{m_F^4} = 3 \cdot 10^{-26} \text{ cm}^3/\text{s}$$

Hence

$$m_{dm} > O(1) \text{ GeV}$$

The mass limit dogma



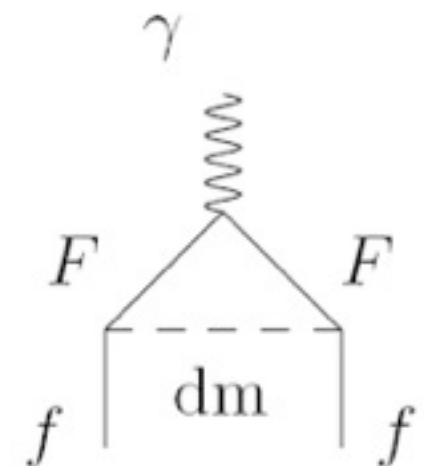
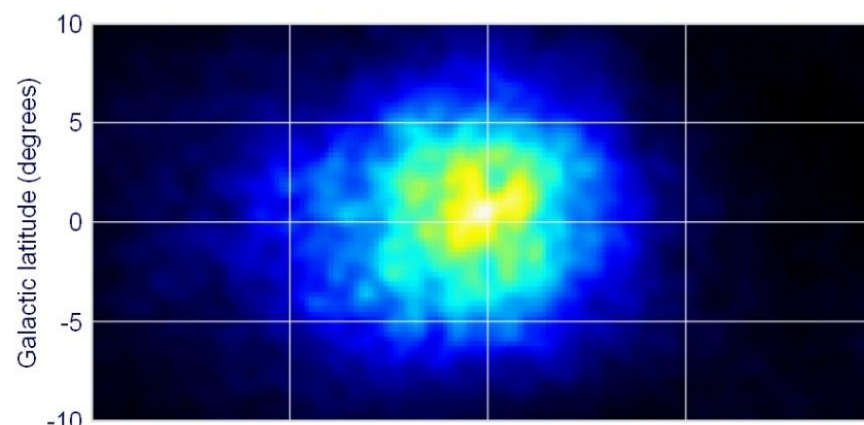
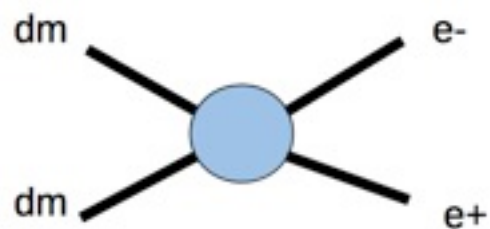
$$\sigma v \propto \frac{1}{m_F^2} = 3 \cdot 10^{-26} \text{ cm}^3/\text{s}$$

m_{dm} is not constrained!

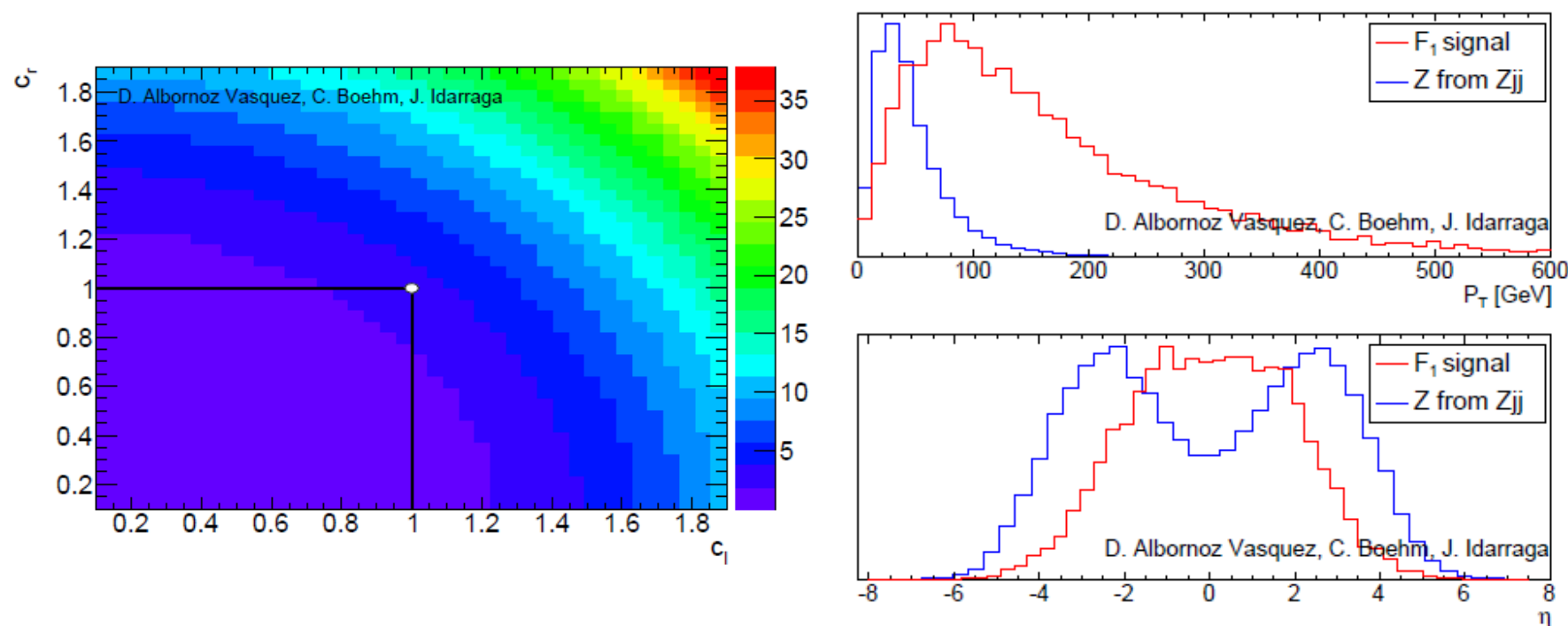
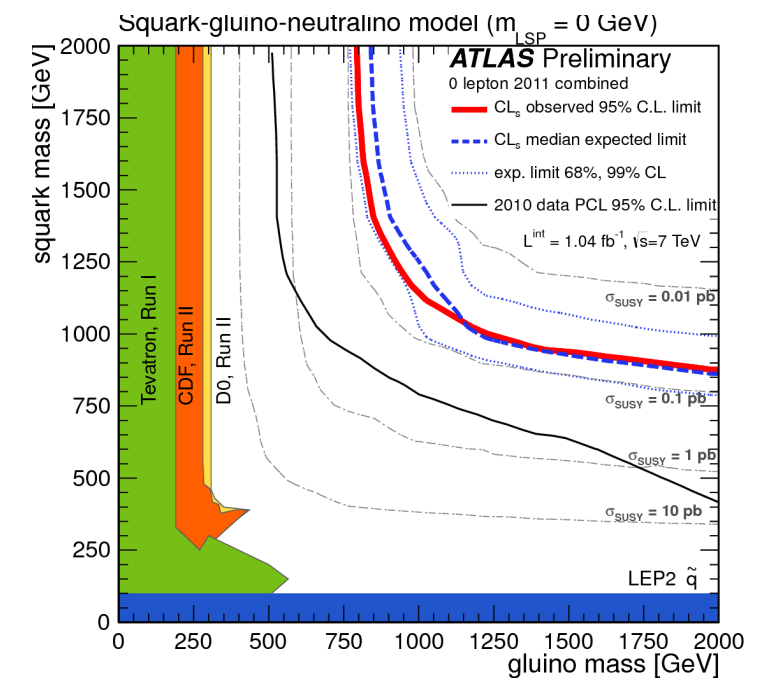
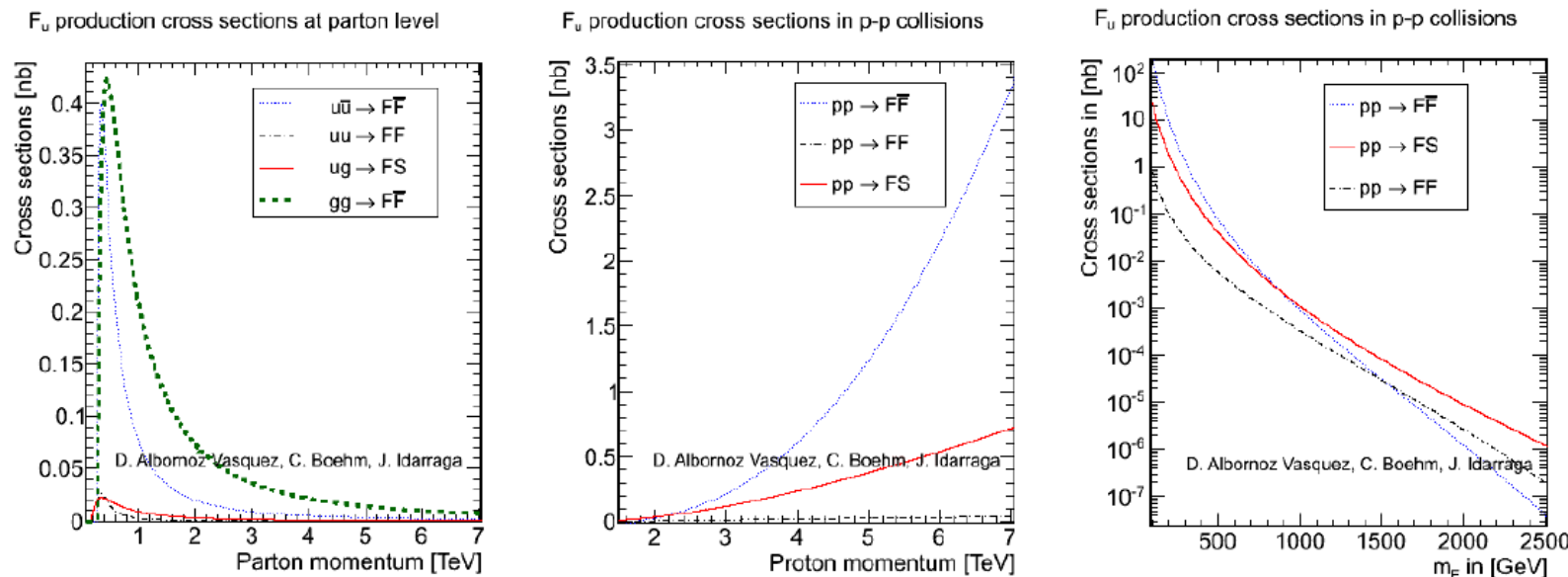
Exceptions to Hut-Lee&Weinberg are possible
Annihilating DM can be lighter than 10 GeV
 In fact, it can be as light as a few keV...

But direct detection experiments
 cannot easily go down 10 GeV

Note that MeV DM could also be interesting (cf 511 keV line)!

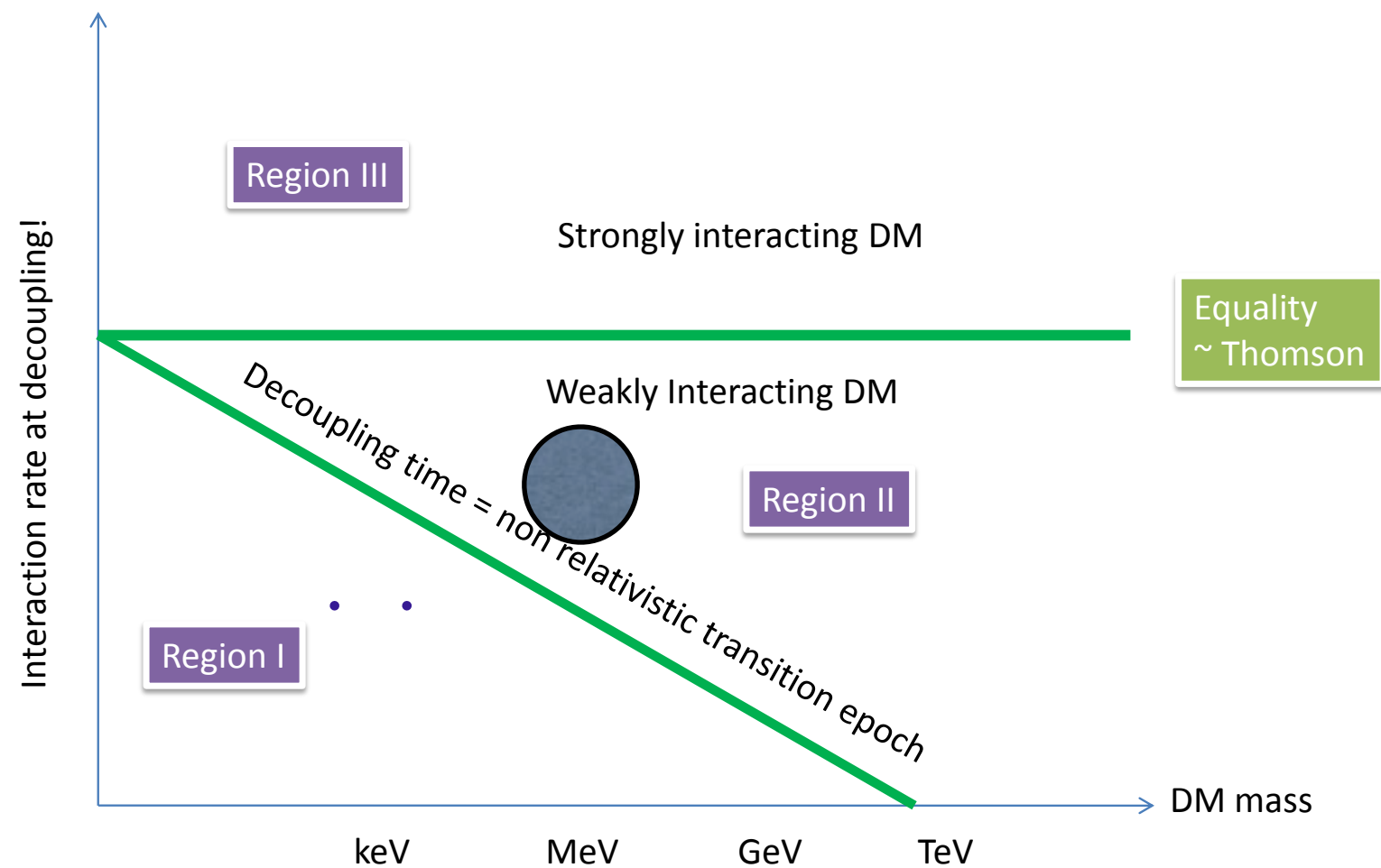


Anyes Taffard



Exclude at 95% C.L.
 $m_{\tilde{g}} \leq 800$ GeV $m_{\tilde{q}} \leq 850$ GeV
 If $m_{\tilde{g}} = m_{\tilde{q}}$, masses < 1075 GeV

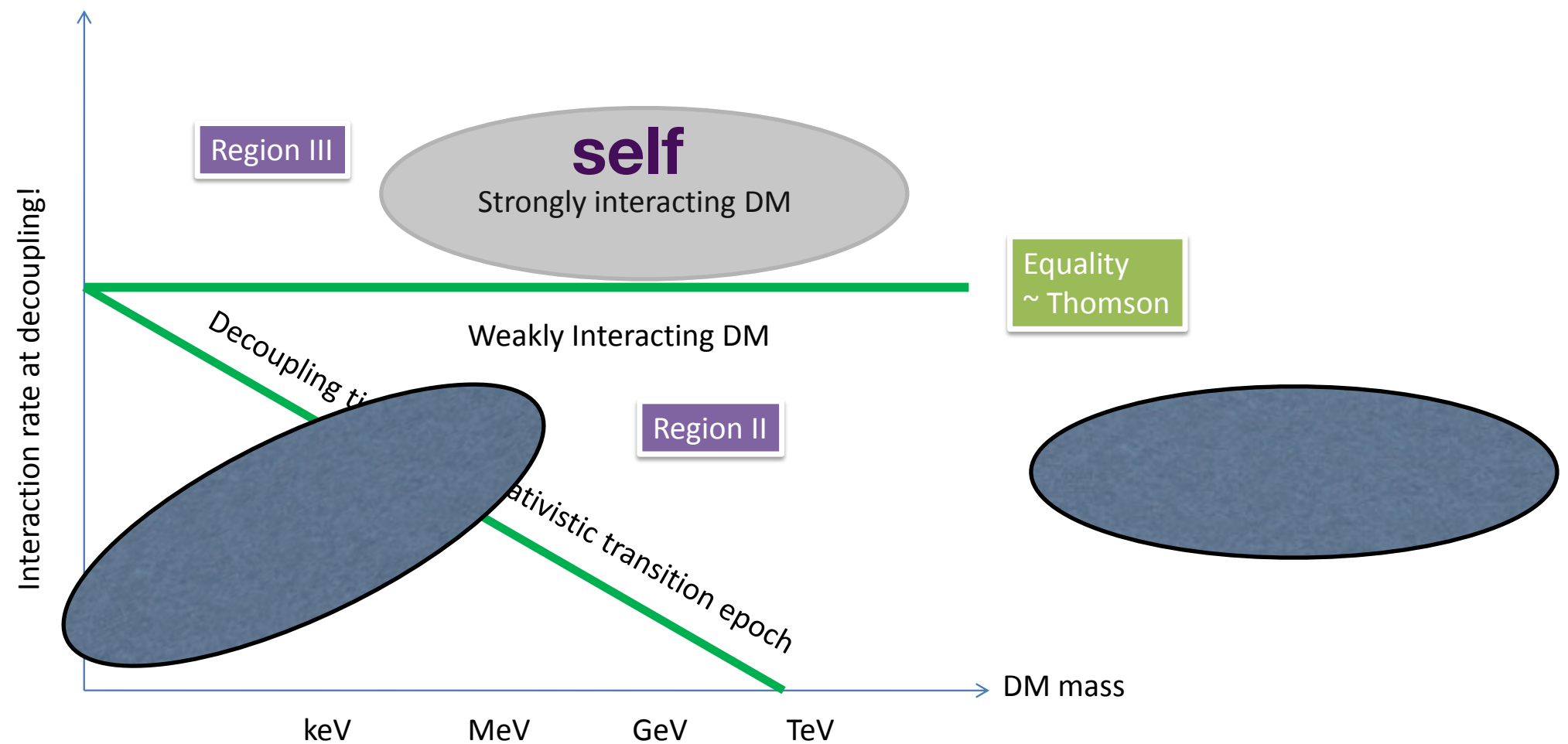
MeV DM possible but either no charged coloured particles or very heavy ones...



So what do we do if nothing is found at
LHC?

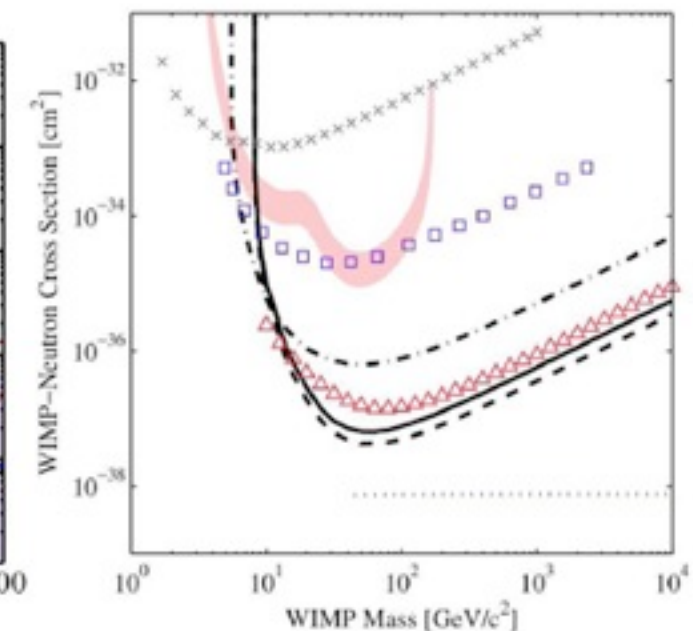
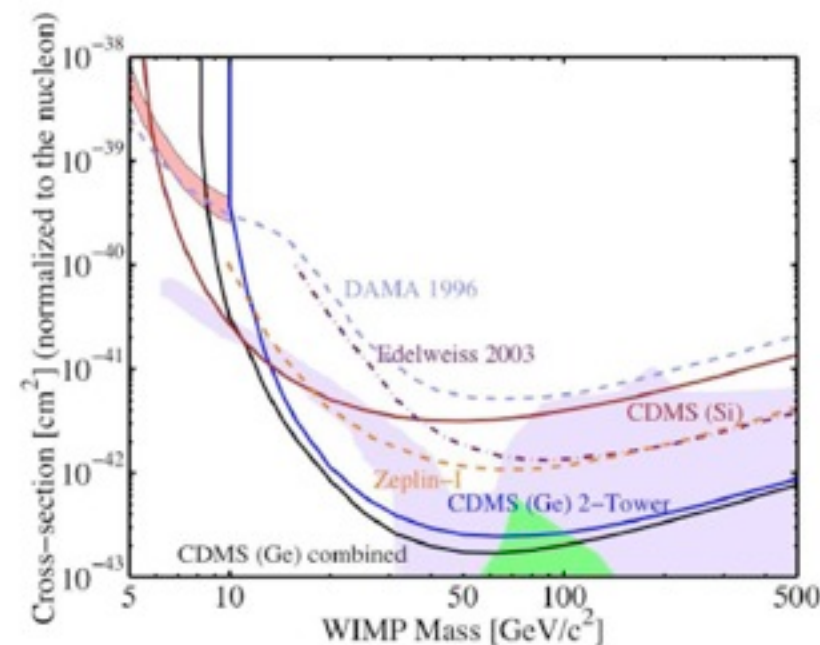
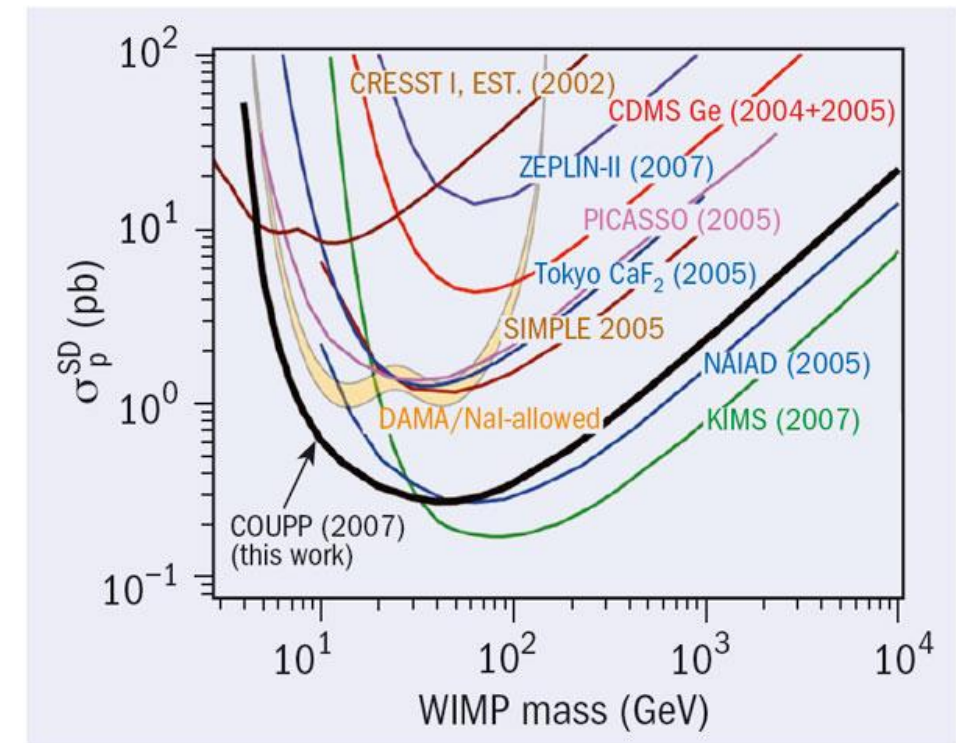
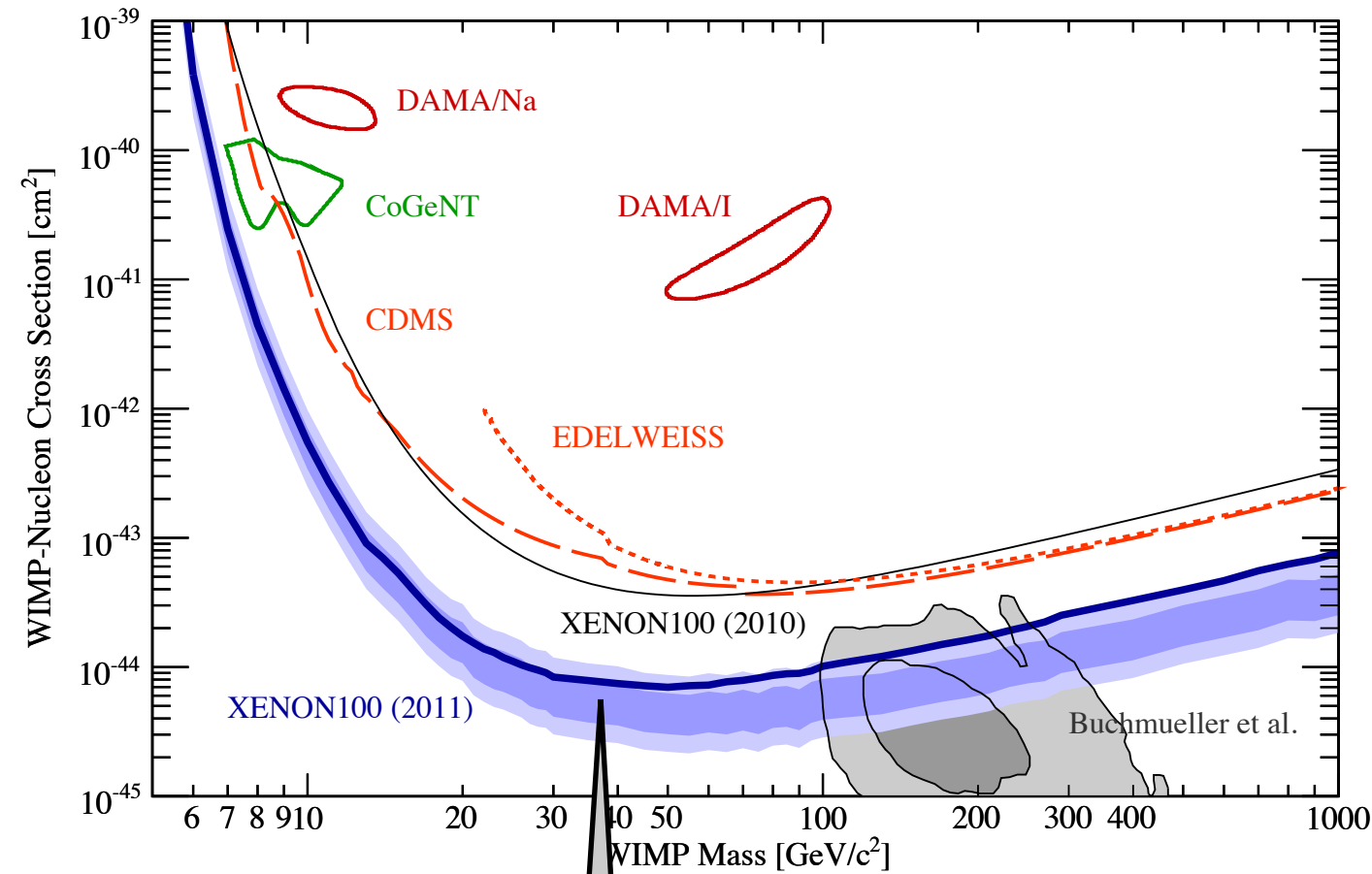
Implication for DM

Start to explore new (and challenging) horizons



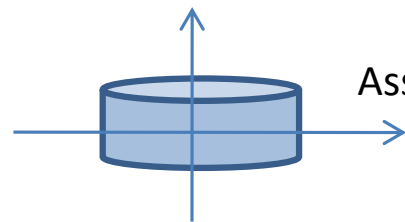
Dark matter direct detection will become useless...unless...

Plethora of experiments but ...



One will have to explore in detail cosmic ray background!

Deianaye et al 2007



Assume that diffusion zone is well approximated by a cylinder

$$\tilde{I}(\lambda_D) = \sum_{i=1}^{\infty} \sum_{n=1}^{\infty} J_0(\alpha_i r / R_{\text{gal}}) \varphi_n(z) \exp \left\{ -\tilde{C}_{i,n} (\tilde{t} - \tilde{t}_S) \right\} R_{i,n}$$

Annotations for the equation:

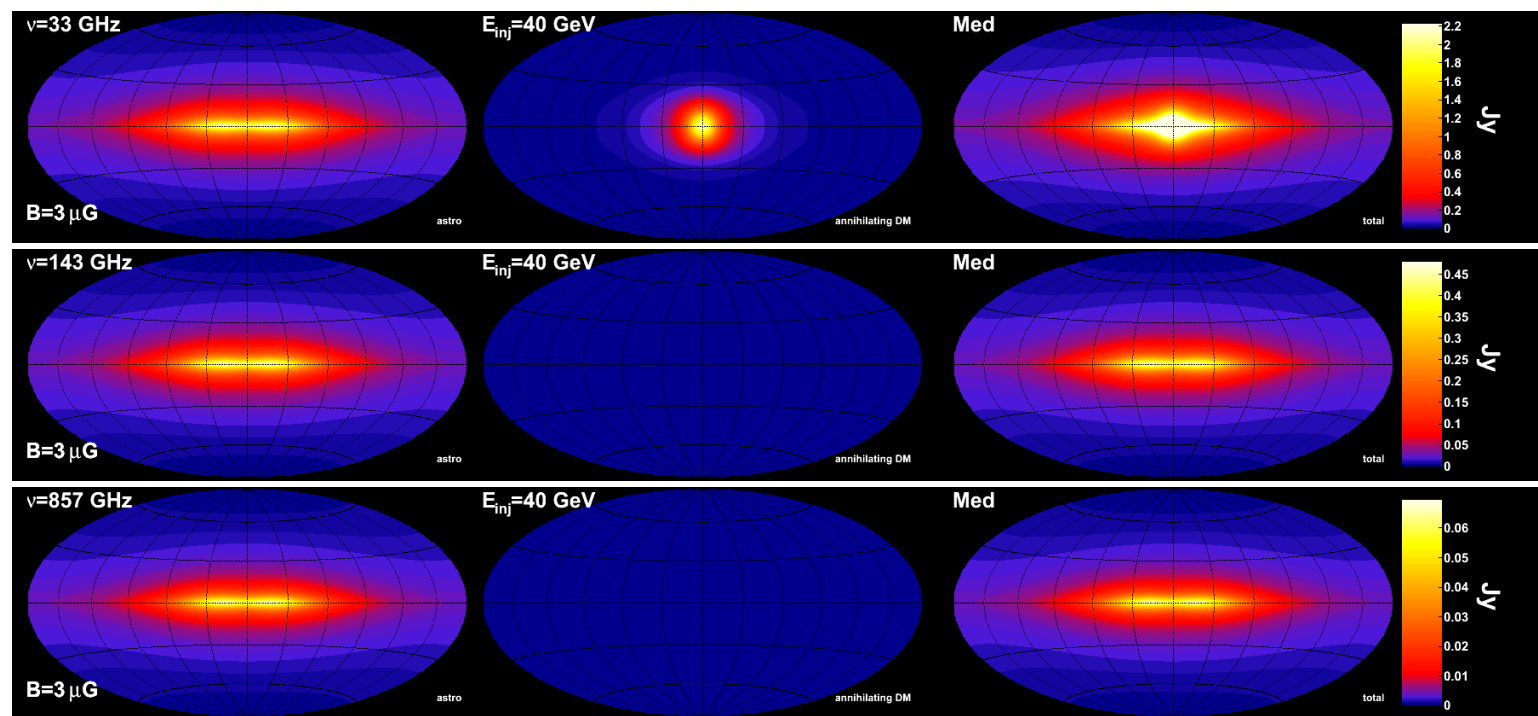
- Radius of the cylinder (points to R_{gal})
- Bessel in r_{cyl} (points to J_0)
- Fourier in z_{cyl} (points to φ_n)
- Propagation length (points to $\tilde{C}_{i,n}$)
- Source term (points to $R_{i,n}$)

$$\lambda_D^2 = 4K_0 \tau_E \left\{ \frac{\epsilon^{\delta-1} - \epsilon_S^{\delta-1}}{1 - \delta} \right\}$$

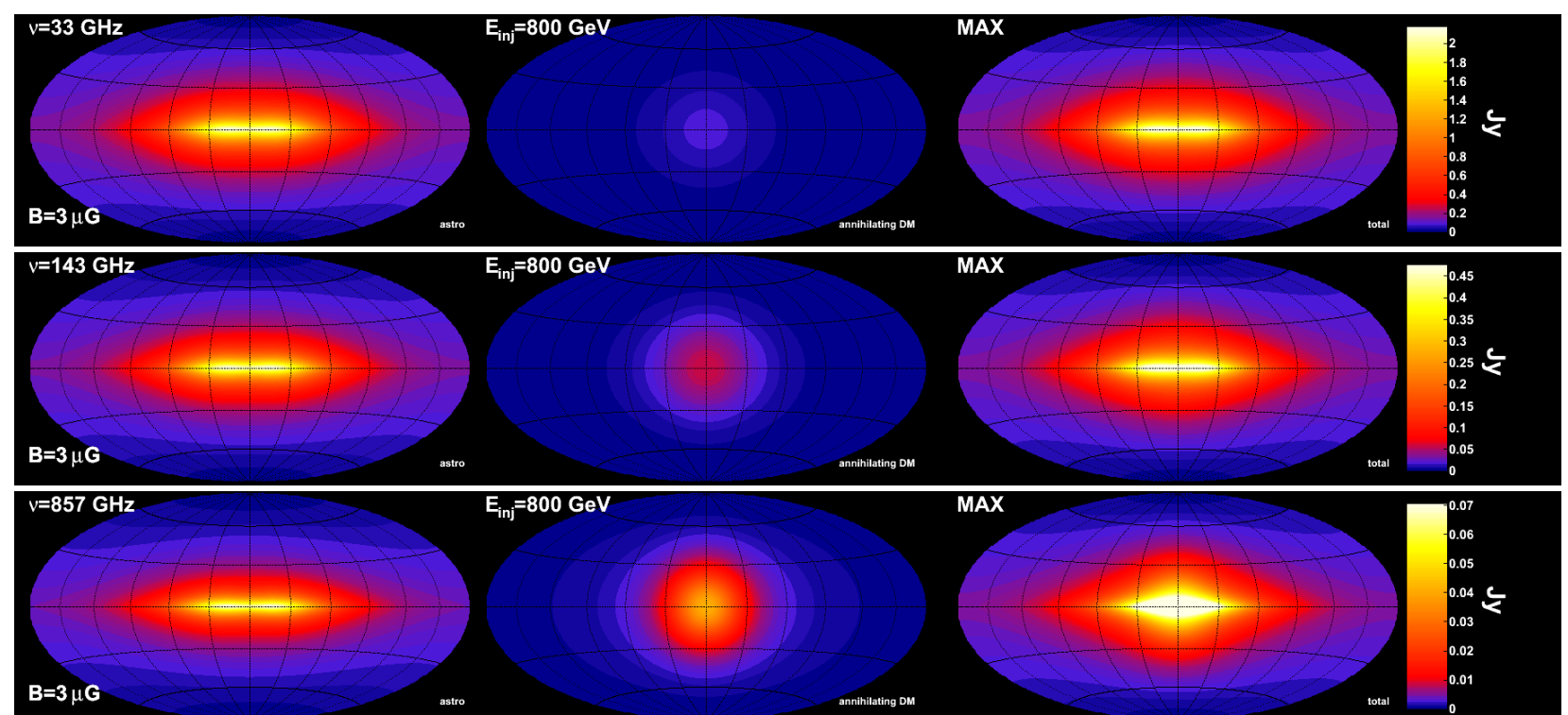
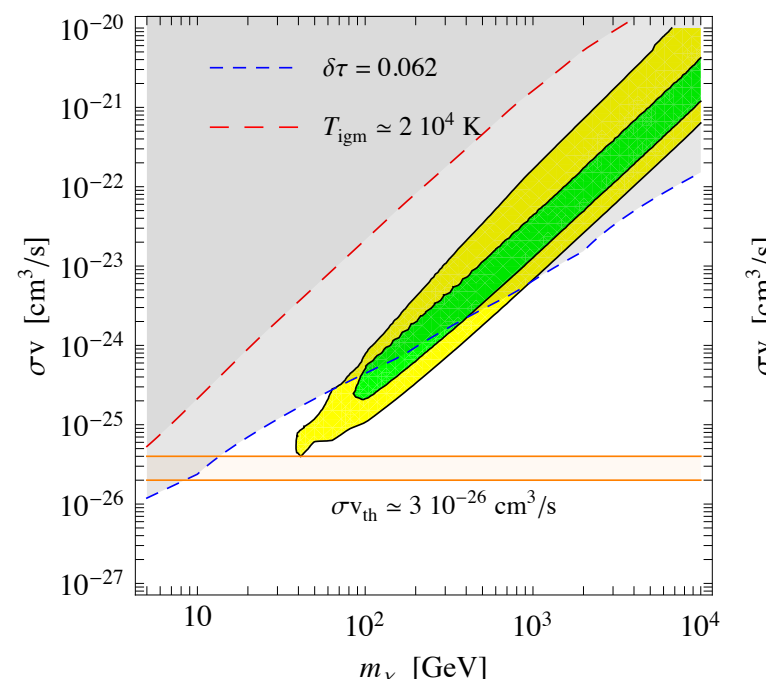
$$K(x, E) = K_0 \epsilon^{\delta} \text{ where } \epsilon = E/E_0$$

Finally flux is proportional to the integral over l.o.s (and dE) of $\tilde{I}(\lambda_D)$

And we can still use the EM spectrum
(although difficult with Planck foregrounds)

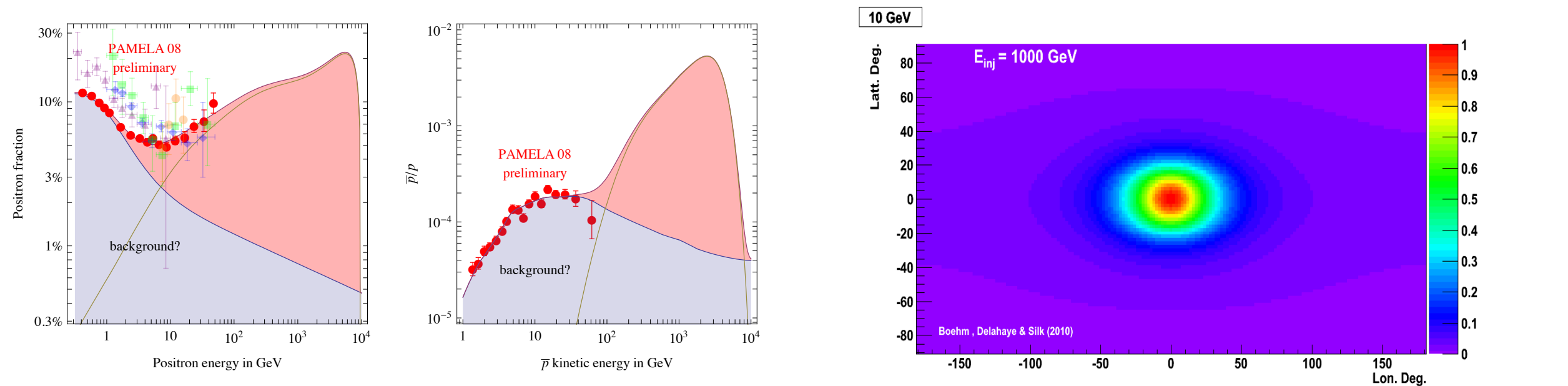


at submillimetre frequencies
but predictions at
radio, X and Gamma-ray too..

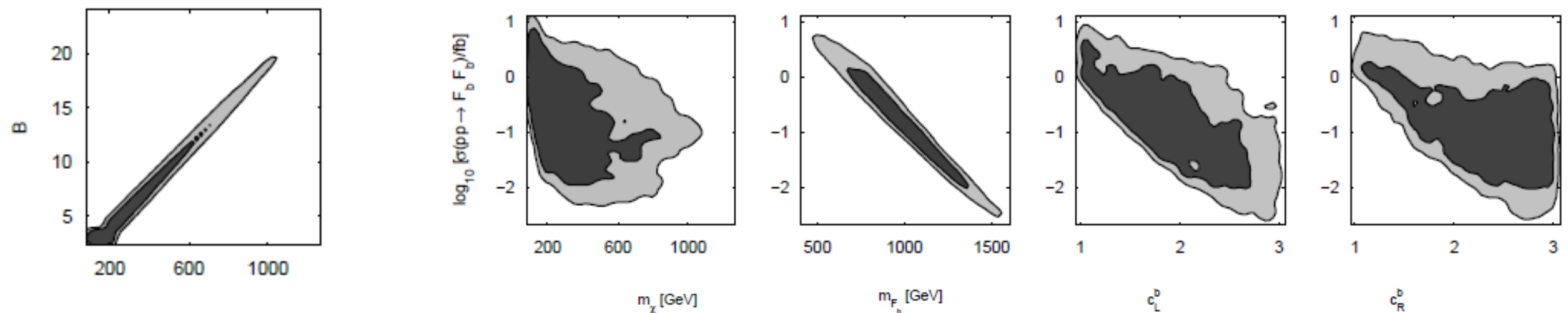


The final states of DM annihilations thought again

- Direct Detection based on DM interactions with quarks: **Leptophilic DM** could explain the absence of signal!
- Still a chance to explain PAMELA but FERMI experiment is very constraining!



- Leptophilic maybe ... but not necessarily quarkophobic!



What else can we do in absence of
positive results at LHC?

Explore WDM/CDM more!

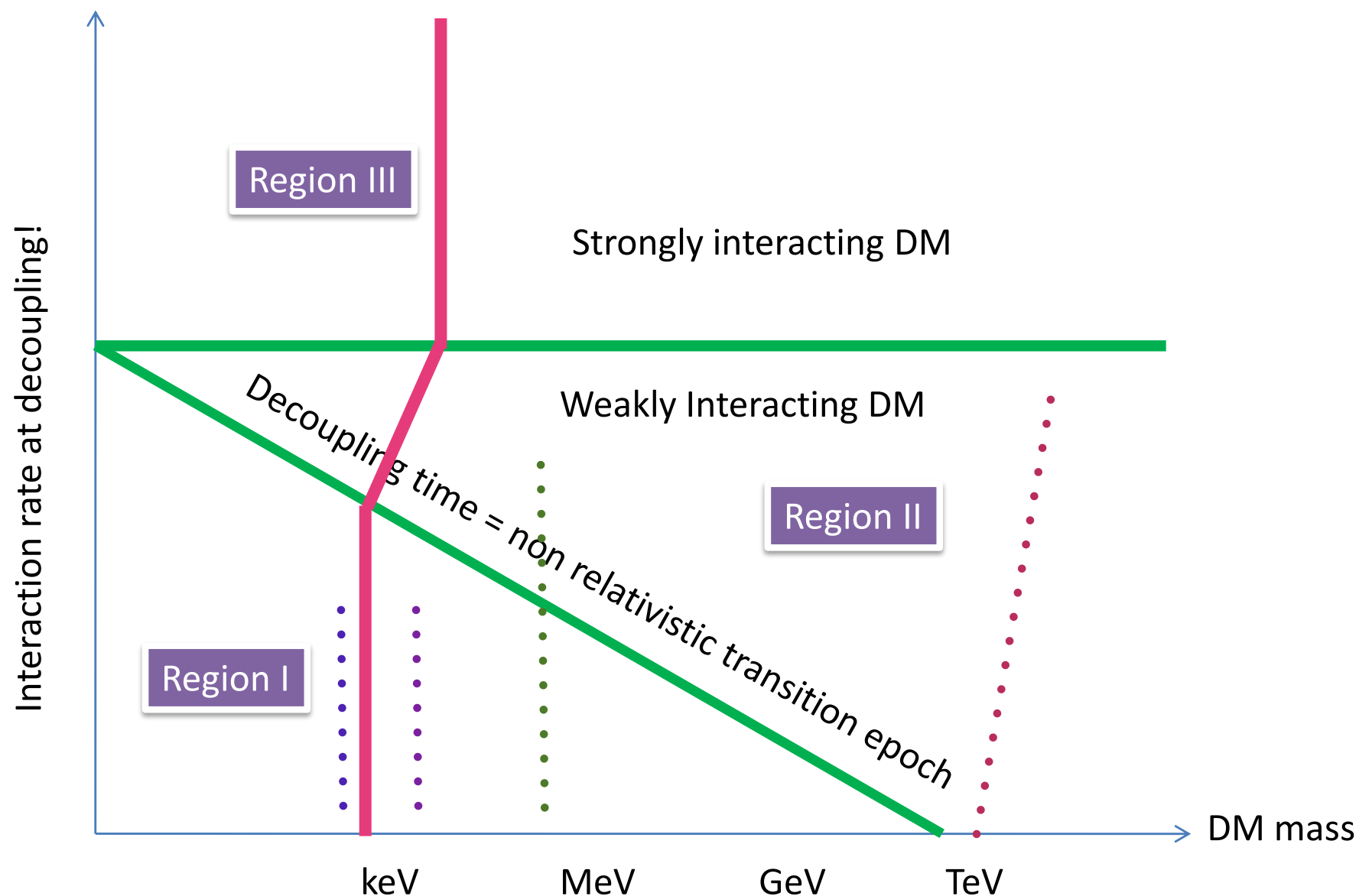
LSS formation can reduce the parameter space quite a bit!

The sole requirement that the free-streaming length does not exceed the length of the smallest fluctuations needed to survive to explain the formation of objects down to $\sim 10^6$ Msol cuts part of the parameter space.

In particular, it leads to $m_{\text{dm}} > \text{keV}$,
not only in region I [that everybody knows] but also for region II and III with
an interesting property for region II

$$l_{fs(i)} \propto \int_{t_{dec(i)}}^t \frac{v_i dt}{a}.$$

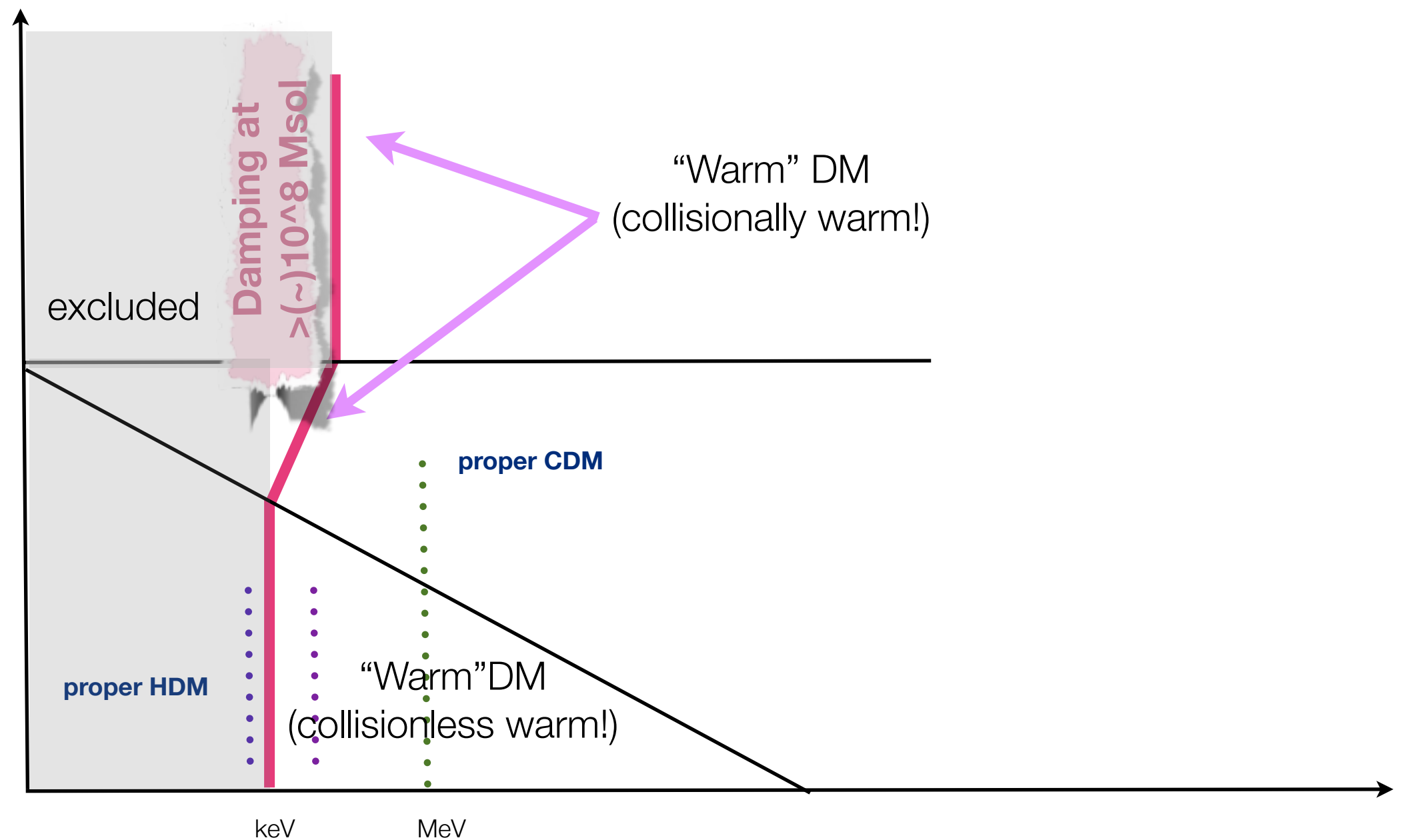
The decoupling time is important!!!



Primordial DM fluctuations are erased if DM freely propagate while still relativistic or so

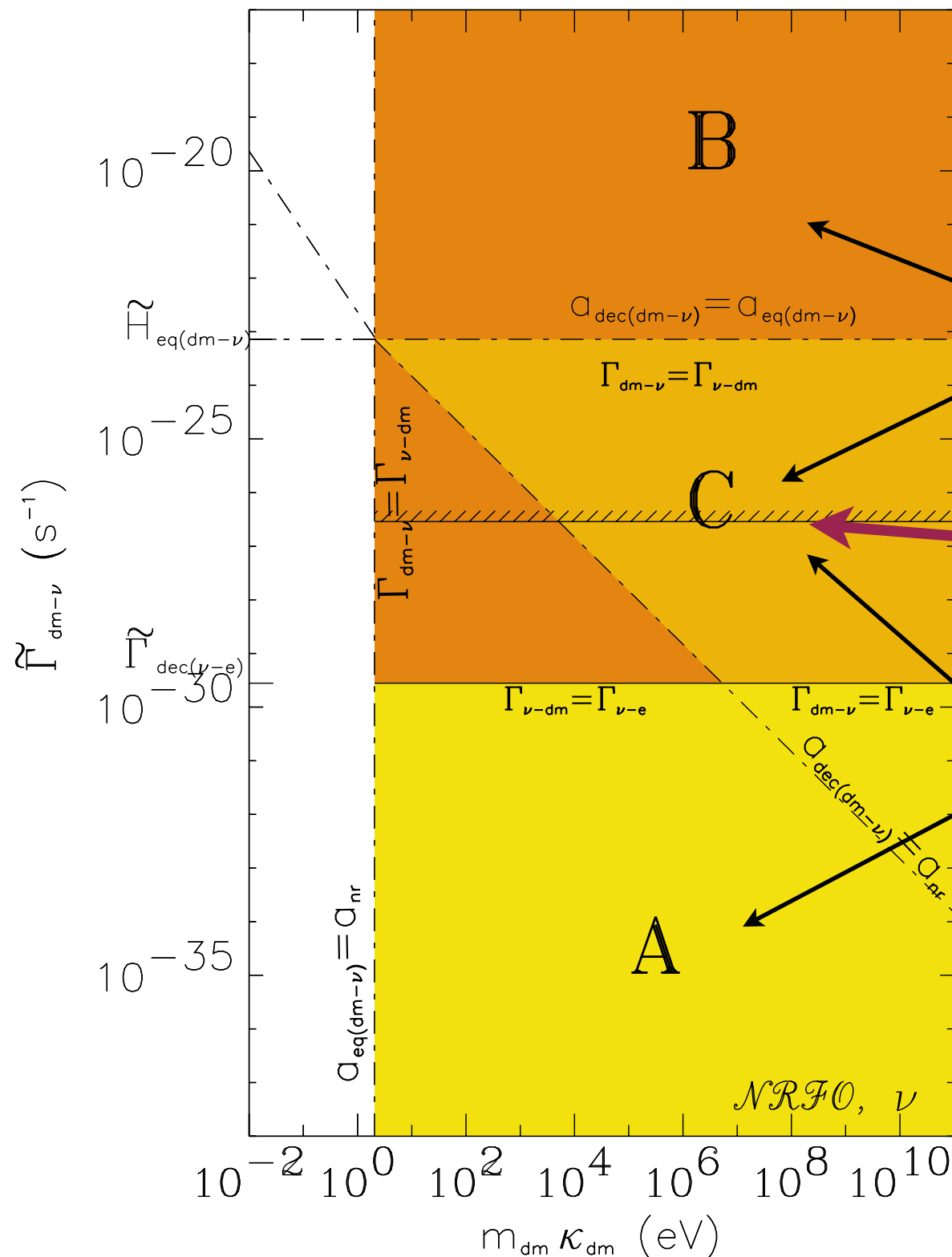
Definition of WDM revisited....

[astro-ph/0012504](#); [astro-ph/0410591](#)



DM-neutrino interactions: a new damping effect

astro-ph/0012504; astro-ph/0410591



If DM interacts with neutrinos at ν -e decoupling, the neutrino free-streaming can be transmitted to the DM

Forbidden regions

Collisional WDM!

Allowed regions

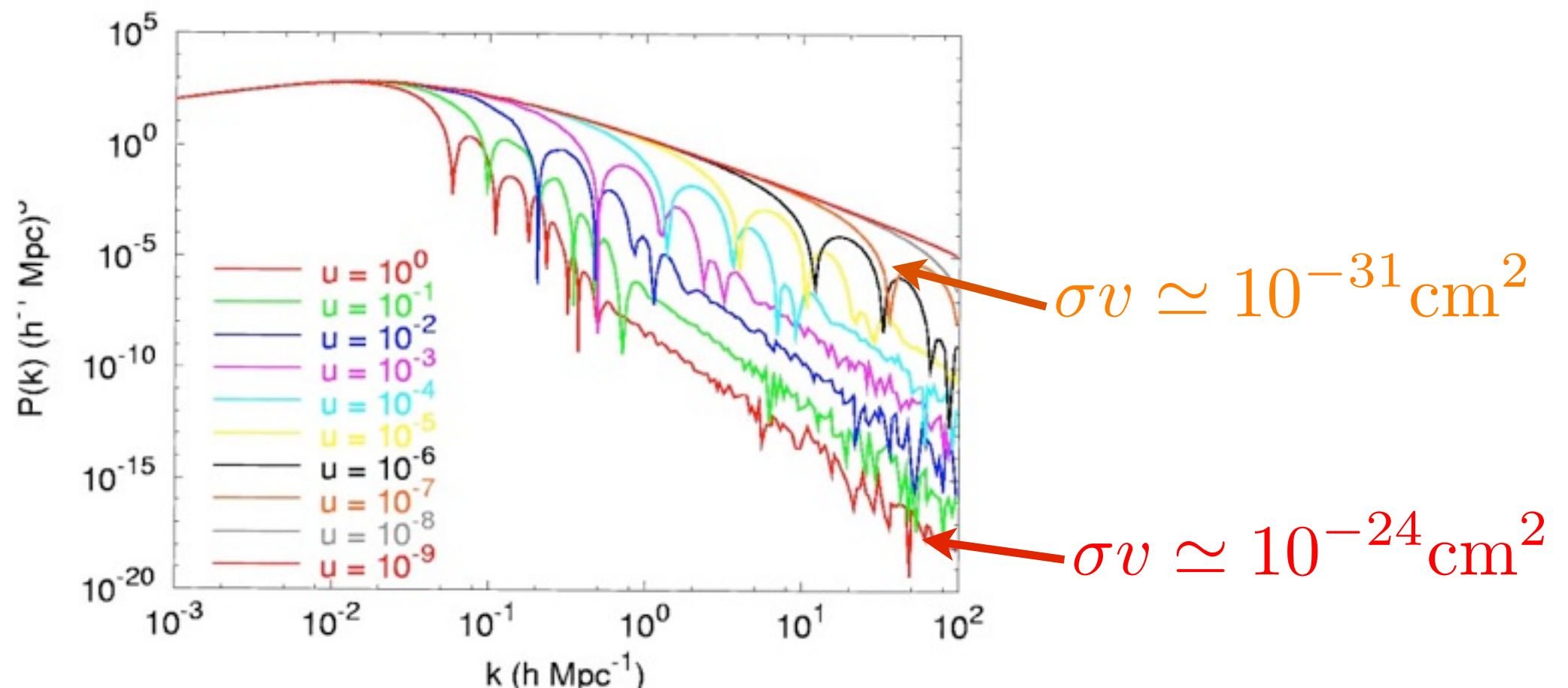
$n_{DM} \ll n_{\text{neutrinos}}$
neutrinos are almost decoupled

(Coll)WDM with photons and baryons: Let us now use the Power Spectrum!

[astro-ph/0112522](#)

DM interactions cannot be arbitrarily large (even with baryonic matter or themselves, cf self-interactions)

$P(k)$ for DM interactions = $u \cdot \text{Thomson cross section}$

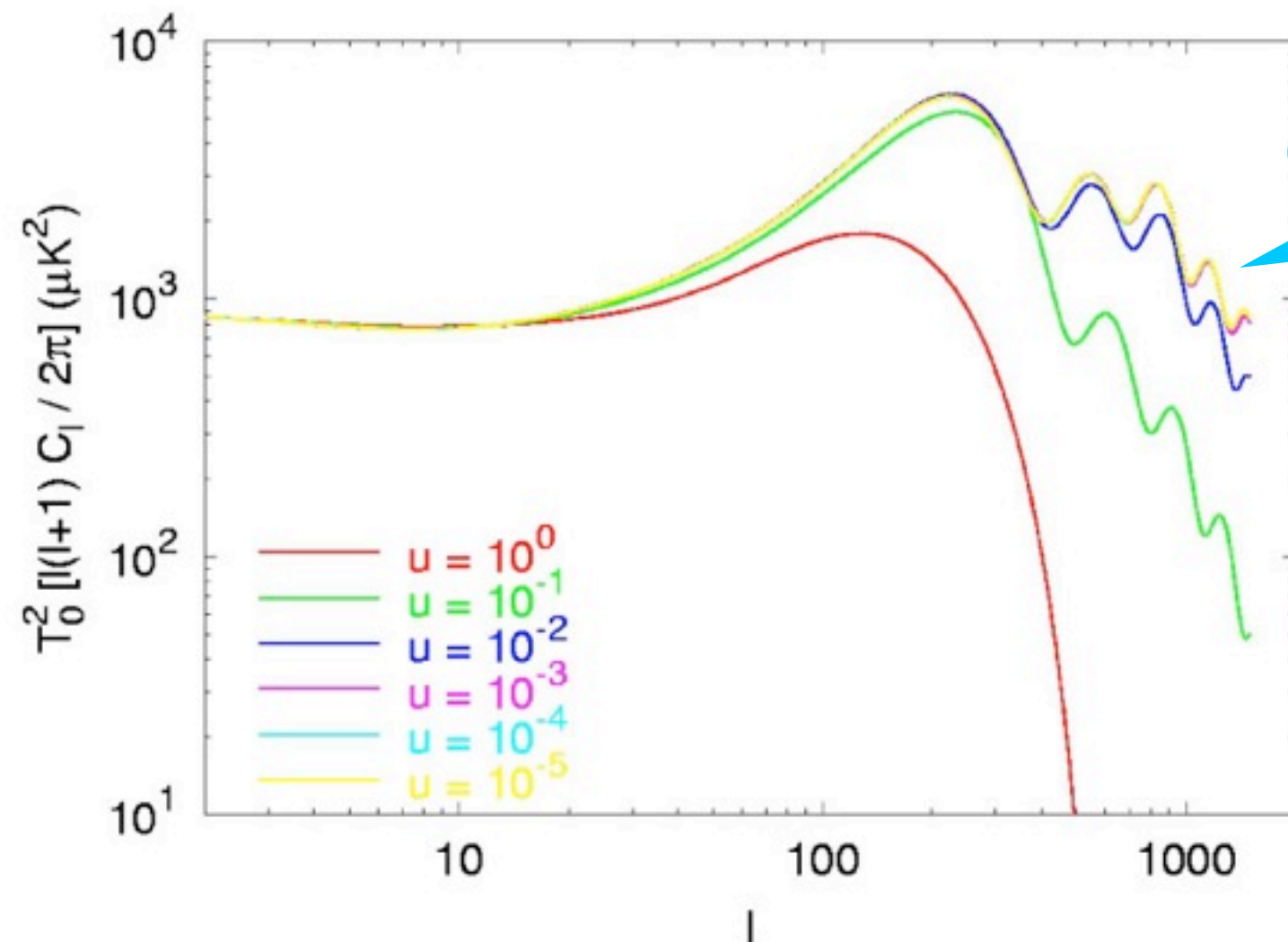


The stronger the interactions, the more pronounced the oscillations

CMB physics/PLANCK can be used to constrain the DM interactions

[astro-ph/0112522](#)

Oscillations are a distinctive sign
They should lead to a signature
in the Planck Cl



Now time to zoom in!

Conclusion

**Should we be worried if no Higgs
and new physics is seen at LHC?**

Conclusion of 10-20-30-40 years of DM hunting

