# SOFT WALLS: STABILIZATION AND EWSB

Gero von Gersdorff (École Polytechnique) University of Sussex, October 2010

Collaboration with J.A.Cabrer and M.Quirós

arXiv:1005.5134,0907.5361, w.i.p.

# OUTLINE

- Features of Warped Extra Dimensions
- Stabilizing Models with 2 branes
- Soft Wall models (Models with I brane)
- Stabilizing the Soft Wall
- Simulating Soft Walls with IR branes
- EWSB and EWPT in warped spaces

#### OPEN QUESTIONS IN THE SM (AND BEYOND)

- What is the origin of Electroweak Symmetry Breaking?
- Why is the scale of the Z and W bosons 10<sup>17</sup> times smaller than the Planck mass? (Hierarchy Problem)
- Why is there such a huge hierarchy in the masses of the Standard Model fermions?
- What is the origin of neutrino masses?
- If there is Supersymmetry, how is it broken?
- If there is a Grand Unified Theory, how is it broken to the SM, and why are there no colored Higgses?

#### All these issues can be addressed in models with Extra Dimensions

# RS MODELS

Randall & Sundrum '99

#### RS MODELS

Randall & Sundrum '99



Fifth Dimension



### RS MODELS

Randall & Sundrum '99







#### Fundamental cutoff scale is redshifted



#### Fundamental cutoff scale is redshifted



#### Fundamental cutoff scale is redshifted

Arkani-Hamed & Schmalz '00, Huber & Shafi '00, Shifman & Dvali '00, Gherghetta & Pomarol '00













# SUMMARY RS MODELS

# SUMMARY RS MODELS

- Extra Dimensions with WARPED background successful for
  - Explaining electroweak hierarchy
  - Explaining fermion mass hierarchy

# SUMMARY RS MODELS

- Extra Dimensions with WARPED background successful for
  - Explaining electroweak hierarchy
  - Explaining fermion mass hierarchy
- Other features
  - Distinctive Collider Signature (KK gravitons)
  - Dual to strongly coupled gauge theories in 4D
  - "Modelling QCD"

# PROBLEMS

- Pure 5D Gravity with negative Cosmological constant (and appropriate brane tensions) has RS as a solution.
- BUT: Interbrane distance is UNDETERMINED
- There is an extra massless mode (RADION)

$$g_{MN} = g_{MN}^{RS} + \begin{pmatrix} h_{\mu\nu} & \\ & h_{55} \end{pmatrix}$$

Both brane tensions need to be fine tuned

- The question of Radius stabilization
  - What determines **DISTANCE** between UV and IR brane?
  - How can I generate a POTENTIAL and a MASS for the Radion?
  - What ensures that the 4D metric is FLAT?
  - How NATURAL is it?

# SUPERPOTENTIAL METHOD

# SUPERPOTENTIAL METHOD

 $\lambda_0(\phi$ 

 $V(\phi)$ 

 $\lambda_1(\phi)$ 

- $ds^{2} = e^{-2A(y)} dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dy^{2}$ • Gravity + scalar field with bulk and brane potential
- Solve Einstein equations coupled to scalar

# SUPERPOTENTIAL METHOD

- Gravity + scalar field  $ds^2 = e^{-2A(y)} dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dy^2$ with bulk and brane potential  $\lambda_0(\phi)$
- Solve Einstein equations coupled to scalar

 $\lambda_0(\varphi$ 

- Define a "Superpotential"  $V(\phi) = 3W'(\phi)^2 12W^2(\phi)$  NO SUSY INVOLVED
- Einstein equations become  $\phi'(y) = W'(\phi)$   $A'(y) = W(\phi)$
- Boundary values from minimizing the 4D potentials

 $V_i(\phi) = \lambda_i(\phi) - 6 W(\phi)$ 

DeWolfe et al '99, Brandhuber & Sfetsos '99

- Solve to get bulk profiles  $\phi'(y) = W'(\phi)$   $A'(y) = W(\phi)$ - Minimize to get brane values  $V_i(\phi_i) = \lambda_i(\phi_i) - 6W(\phi_i)$ 

- Solve to get bulk profiles  $\phi'(y) = W'(\phi)$   $A'(y) = W(\phi)$ - Minimize to get brane values  $V_i(\phi_i) = \lambda_i(\phi_i) - 6W(\phi_i)$ 
  - Notice that  $e^{k y_1} = 10^{16} \implies k y_1 \approx 37$ - Choose some suitable W such that  $k y_1 = \int_{\phi_0}^{\phi_1} \frac{1}{W'} \approx 37$ - Now shift superpotential  $W \rightarrow W + k$   $A(y) \rightarrow A(y) + k y$ - Adds warping without changing the value of  $k y_1$

DeWolfe et al '00, Cabrer, GG & Quirós '09

## GOLDBERGER WISE

#### Goldberger & Wise '99



#### Do we need two branes?

# GAUGE/GRAVITY DUALITY

# GAUGE/GRAVITY DUALITY

- Gravity/Gauge theory correspondence stipulates that the 5D theory is dual to a strongly coupled 4D gauge theory that
  - is approximately conformal in the UV
  - has large number of colors
  - describes the same physics as 5D theory

# GAUGE/GRAVITY DUALITY

- Gravity/Gauge theory correspondence stipulates that the 5D theory is dual to a strongly coupled 4D gauge theory that
  - is approximately conformal in the UV
  - has large number of colors
  - describes the same physics as 5D theory
- KK modes correspond to resonances of gauge theory
  - RS with two branes: KK spectrum is roughly  $m_n^2 \sim n^2$
  - 4D strongly coupled gauge theories have many more possibilities.



#### IR brane can be replaced by SOFT WALL
Soft Walls models only possess a single (UV) brane, but nevertheless exhibit a finite length in the 5th dimension. The IR brane is replaced by a curvature singularity at which the metric vanishes.

Soft Walls models only possess a single (UV) brane, but nevertheless exhibit a finite length in the 5th dimension. The IR brane is replaced by a curvature singularity at which the metric vanishes.



Soft Walls models only possess a single (UV) brane, but nevertheless exhibit a finite length in the 5th dimension. The IR brane is replaced by a curvature singularity at which the metric vanishes.



Profiles diverge at finite y if  $W(\phi) \sim \phi^2$  or faster!

#### APPLICATIONS

# APPLICATIONS

- Things that CAN be done with Soft Walls
  - Electroweak Breaking
  - Strong interactions (AdS/QCD)
  - Flavour physics

Batell, Gherghetta & Sword '08, Falkowski & Perez-V. '08, Cabrer, GG & Quiros (in progr.)

Karch et al '06, Gursoy et al '07, Batell & Gherghetta '08,

Atkins & Huber '10

# APPLICATIONS

- Things that CAN be done with Soft Walls
  - Electroweak Breaking
  - Strong interactions (AdS/QCD)
  - Flavour physics

Batell, Gherghetta & Sword '08, Falkowski & Perez-V. '08, Cabrer, GG & Quiros (in progr.)

Karch et al '06, Gursoy et al '07, Batell & Gherghetta '08,

Atkins & Huber '10

- Things that CANNOT be done with Soft Walls
  - Solve Cosmological Constant problem

Arkani-Hamed et al '00, Kachru, Schulz & Silverstein '00 , Csaki et al '00

Forste et al '00, Cabrer, GG & Quirós '09

## SPECTRA WITH SOFT WALLS

# SPECTRA WITH SOFT WALLS

• Even though the physical length is finite, the conformal length might be either finite or infinite:

 $\begin{array}{ll} \mbox{Proper Length coordinates} & \mbox{Conformally flat coordinates} \\ ds^2 &= e^{-2A(y)} dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dy^2 & ds^2 &= e^{-2A(z)} (dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dz^2) \\ y_s &< \infty, \qquad z_s &= z(y_s) & \mbox{can be finite or infinite} \end{array}$ 

# SPECTRA WITH SOFT WALLS

• Even though the physical length is finite, the conformal length might be either finite or infinite:

Proper Length coordinatesConformally flat coordinates $ds^2 = e^{-2A(y)} dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dy^2$  $ds^2 = e^{-2A(z)} (dx^{\mu} dx^{\nu} \eta_{\mu\nu} + dz^2)$  $y_s < \infty$ , $z_s = z(y_s)$ can be finite or infinite

 In the conformally flat frame, the KK spectrum of any bulk field follows a Schrödinger Equation

$$-\psi''(z) + \hat{V}(z)\psi(z) = m^2\psi(z)$$
  
Depends on the background

#### NON CONFINING POTENTIALS Conformal length infinite



$W(\phi)$	$\leq \phi^2$	$> \phi^2 \\ < e^{\phi}$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$ $0 < \beta \le \frac{1}{2}$	$> e^{\phi} \phi^{\frac{1}{2}}$ $< e^{2\phi}$	$\geq e^{2\phi}$
$y_s$	$\infty$		A State of the second second			
$z_s$			$\infty$	finite		
mass	continuous		continuous	discrete		
spectrum			w/ mass gap	$m_n \sim n^{2\beta}$ $m_n \sim r$		$\sim n$
consistent			VOS			no
solution	yes					110

$W(\phi)$	$\leq \phi^2$	$> \phi^2$	$e^{\phi}$	$e^{\phi}\phi^{eta}$	$> e^{\phi}\phi^{\frac{1}{2}}$	$\geq e^{2\phi}$		
$vv(\phi)$		$< e^{\phi}$		$0 < \beta \le \frac{1}{2}$	$< e^{2\phi}$			
$y_s$	$\infty$		A State of the second second					
$z_s$			$\infty$	finite				
mass	conti	0110119	continuous	discrete				
spectrum	continuous		w/ mass gap	$ m_n \sim n^{2\beta} $ $m_n \sim r$		$\sim n$		
consistent				no				
solution		yes						

-Asymptotic behaviour of W

Γ	$W(\phi)$	$\leq \phi^2$	$> \phi^2 \\ < e^{\phi}$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$ $0 < \beta \le \frac{1}{2}$	$> e^{\phi} \phi^{\frac{1}{2}}$ $< e^{2\phi}$	$\geq e^{2\phi}$		
+	$y_s$	$\infty$		finite					
	$z_s$			$\infty$	finite				
	mass	continuous		continuous	discrete				
	spectrum			w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$		
	consistent solution			yes			no		

Asymptotic behaviour of W
 Singularity in "proper distance"

Γ	$W(\phi)$	$\leq \phi^2$	$> \phi^2 \\ < e^{\phi}$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$ $0 < \beta \le \frac{1}{2}$	$> e^{\phi} \phi^{\frac{1}{2}}$ $< e^{2\phi}$	$\geq e^{2\phi}$	
	$y_s$	$\infty$		A State of the second s	finite			
	$z_s$			$\infty$		finite		
	mass	continuous		continuous	discrete			
	spectrum		liuous	w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$	
	consistent solution			yes			no	

Asymptotic behaviour of W
Singularity in "proper distance"
Singularity in "conformal distance"

Г	$W(\phi)$	$\leq \phi^2$	$> \phi^2 \\ < e^{\phi}$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$ $0 < \beta \le \frac{1}{2}$	$> e^{\phi} \phi^{\frac{1}{2}}$ $< e^{2\phi}$	$\geq e^{2\phi}$
	$y_s$	$\infty$			finite		
	$z_s$			$\infty$	finite		
	mass	continuous		continuous	discrete		
	spectrum		luous	w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$
	consistent solution			yes			no

Asymptotic behaviour of W
 Singularity in "proper distance"
 Singularity in "conformal distance"
 Asymptotic form of the spectrum

	$W(\phi)$	$\leq \phi^2$	$> \phi^2$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$	$> e^{\phi} \phi^{\frac{1}{2}}$	$\geq e^{2\phi}$	
			$< e^{\varphi}$		$0 < \beta \le \frac{1}{2}$	$ $ < $e^{2\phi}$		
	$y_s$	$\infty$			finite			
	$z_s$			$\infty$		fini	te	
	mass	conti	2110110	continuous	d	liscrete		
	spectrum	continuous		w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$	
	consistent			VOS			no	
	solution			ycs				
Asymptotic behaviour of W — Finite Length								
——Singularity in "conformal distance"								
——Asymptotic form of the spectrum Gursoy et al '07.								

Cabrer, GG & Quirós '09

		$W(\phi)$	$\leq \phi^2$	$> \phi^2 \\ < e^{\phi}$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$ $0 < \beta \le \frac{1}{2}$	$ \begin{array}{ c c c } > e^{\phi} \phi^{\frac{1}{2}} \\ < e^{2\phi} \end{array} $	$\geq e^{2\phi}$				
		$y_s$	$\infty$			finite						
+		$Z_S$			$\infty$		fini	te				
		mass	conti	0110115	continuous	discrete						
		spectrum	continuous		w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$				
		consistent solution			yes		no					
Asymptotic behaviour of W — Finite Leng								Finite Length Mass gap appears				

		$W(\phi)$	$\leq \phi^2$	$> \phi^2$	$e^{\phi}$	$e^{\phi}\phi^{\beta}$	$> e^{\phi}\phi^{\frac{1}{2}}$	$\geq e^{2\phi}$	
				$< e^{\varphi}$		$0 < \beta \leq \frac{1}{2}$	$< e^{-\varphi}$		
		$y_s$	$\infty$			finite			
		$Z_S$			$\infty$		fini	te	
		mass	conti	110115	continuous	d	liscrete		
		spectrum	continuous		w/ mass gap	$m_n \sim n^{2\beta}$	$m_n \sim$	$\sim n$	
		consistent						12.0	
		solution			yes			IIO	
						<b>1</b>			
	L	Asymptoti	c beha	viour	ofW	- Finite Length			
L		Singularity	in "pro	istance''	– Mass gap appears				
		Singularity	in "cor	Spectrum discrete					
		Asymptoti	c form	Gursoy et al '07, Cabrer GG & Quirós '09					

## SOFT WALL STABILIZATION

# SOFT WALL STABILIZATION

- Stabilization works similar as before
- Choose some suitable W such that

$$ky_s = \int_{\phi_0}^{\infty} \frac{1}{W'(\phi)} \approx 37$$

- Now shift superpotential  $W \to W + k$  $A(y) \to A(y) + k y$ 

- Shift does not change position of singularity

# SOFT WALL STABILIZATION

- Stabilization works similar as before - Choose some suitable W such that  $ky_s = \int_{\phi_0}^{\infty} \frac{1}{W'(\phi)} \approx 37$ - Now shift superpotential  $W \to W + k$   $A(y) \to A(y) + k y$ - Shift does not change position of singularity

The Warping affects the Mass scale: - The Unparticle mass gap - The level spacing in the discrete case



Cabrer, GG & Quirós '09

## PARTICULAR MODELS

Consider the class of models  $W(\phi) = k(1 + e^{\nu\phi})$  $ky_s = \frac{1}{\nu^2} e^{-\nu\phi_0} \approx 37$  for O(1) negative values for  $\phi_0$ 







gap + continuous



gap+very densely spaced discretuum



Discrete, hard-wall like

#### WAVE FUNCTIONS



 $z/z_s$ 

## THE RADION SPECTRUM

#### THE RADION SPECTRUM



m/
ho(
u)

#### THE RADION SPECTRUM



m/
ho(
u)

 $\mathcal{V}$ 

#### NO ZERO MODE AND NO TACHYON (STABLE)

### EFFECTIVE IR BRANES

SOFT WALLS more general than HARD WALLS (IR branes)
 IR branes are far simpler to deal with analytically
 Can one "simulate" soft wall effects with IR branes?



#### PLAN

Integrate over the region near the singularity

$$S = \int_{0}^{y_{s}} d^{5}x \,\mathcal{L}_{\text{bulk}} \longrightarrow S = \int_{0}^{y_{1}} d^{5}x \,[\mathcal{L}_{\text{bulk}} + \mathcal{L}_{\text{SW}} \,\delta(y - y_{1})]$$
  
Equivalent description of SW in terms of IR brane



## PI AN

Integrate over the region near the singularity

$$S = \int_{0}^{y_{s}} d^{5}x \,\mathcal{L}_{\text{bulk}} \longrightarrow S = \int_{0}^{y_{1}} d^{5}x \,[\mathcal{L}_{\text{bulk}} + \mathcal{L}_{\text{SW}} \,\delta(y - y_{1})]$$
  
Equivalent description of SW in terms of IR brane



• For  $y_1$  close to  $y_s$  close IR Lagrangian "universal" Facilitates comparison with standard 2 brane compactif. IR Lagrangian makes sense even for p>TeV Useful approximation scheme: Approximate "new" bulk by RS metric

()
#### EFFECTIVE IR LAGRANGIANS

## EFFECTIVE IR LAGRANGIANS

→ Generates IR kinetic terms with form factors

 $\mathcal{L}_{SW} \supset p^2 \mathcal{F}(p) \psi(p)^2$ 

→ RS with IR form factors included reproduce SW spectrum

# EFFECTIVE IR LAGRANGIANS

→ Generates IR kinetic terms with form factors

 $\mathcal{L}_{SW} \supset p^2 \mathcal{F}(p) \psi(p)^2$ 

→ RS with IR form factors included reproduce SW spectrum

→ Bulk potentials generate IR brane potentials

 $\mathcal{L}_{bulk} \supset V(\psi) \longrightarrow \mathcal{L}_{SW} \supset \lambda_1(\psi)$ 

→ Can lead to symmetry breaking BC

→ Electroweak precision observables concern in RS

Huber & Shafi '00, Delgado et al '04

★ Require KK masses of gauge bosons to be 3-10 TeV

→ Electroweak precision observables concern in RS

Huber & Shafi '00, Delgado et al '04

★ Require KK masses of gauge bosons to be 3-10 TeV

→ Before turning on EWSB typical modes look like this:



→ Electroweak precision observables concern in RS

Huber & Shafi '00, Delgado et al '04

★ Require KK masses of gauge bosons to be 3-10 TeV

→ Before turning on EWSB typical modes look like this:

Breaking localized in UV
 little mixing between
 zero and KK modes.
 hierarchy problem

→ Electroweak precision observables concern in RS

Huber & Shafi '00, Delgado et al '04

★ Require KK masses of gauge bosons to be 3-10 TeV

→ Before turning on EWSB typical modes look like this:

Breaking localized in UV
 Iittle mixing between
 zero and KK modes.
 hierarchy problem



#### SANDT PARAMETERS

SW or HW, For generic backgrounds 
$$A(y)$$
,  $h(y)$   

$$\alpha T = s_W^2 m_Z^2 y_1 \int e^{2A(y)} \left(\Omega(y) - \frac{y}{y_1}\right)^2$$

$$\alpha S = m_Z^2 y_1 \int e^{2A(y)} \left(\Omega(y) - \frac{y}{y_1}\right) \left(\frac{y}{y_1} - 1\right)$$

$$\Omega(y) = \int h^2(y) e^{-2A(y)} \quad \text{Higgs on UV brane: } \Omega = 1$$

$$\text{Higgs on IR brane: } \Omega = 0$$

Bounds from T go away with Bulk Custodial Symmetry

**X** Bounds from S remain at  $m_{KK} = 3 \, {
m TeV}$  Agashe et al '04

Bounds from T go away with Bulk Custodial Symmetry

**X** Bounds from S remain at  $m_{KK} = 3 \text{ TeV}$  Agashe et al '04

Sounds from S can improve slightly for bulk Higgs  $m_{KK} \to 2 \,\, {\rm TeV}$ Agashe et al '07

Bounds from T go away with Bulk Custodial Symmetry

**X** Bounds from S remain at  $m_{KK} = 3 \text{ TeV}$  Agashe et al '04

Bounds from S can improve slightly for bulk Higgs  $m_{KK} \to 2 \,\, {\rm TeV}$ Agashe et al '07

 Bounds on S and T do not improve significantly for generic
 metric (IR localized Higgs)
 Delgado & Falkowski '07, Archer & Huber '10

Bounds from T go away with Bulk Custodial Symmetry

**X** Bounds from S remain at  $m_{KK} = 3 \text{ TeV}$  Agashe et al '04

Bounds from S can improve slightly for bulk Higgs  $m_{KK} \to 2 \,\, {\rm TeV}$ Agashe et al '07

 Bounds on S and T do not improve significantly for generic
 metric (IR localized Higgs)
 Delgado & Falkowski '07, Archer & Huber '10

✓ Bounds from S and T can improve by changing metric (SW) and putting Higgs in the bulk. (gravitational fine-tuning?)

Cabrer, GG & Quiros w.i.p.

# CONCLUSIONS

- RS models provide neat way of obtaining EW and fermion mass hierarchy
- IR brane can be consistently replaced by Soft Walls
- Spectra of Soft Wall models richer than in usual RS (gapped continuum, gapped, discretuum, Regge-like, etc.)
- Stabilization can be achieved without ANY fine tuning
- Soft Walls effects can be "simulated" by IR branes with appropriate brane Lagrangians (spectrum, symmetry breaking...)
- EW precision parameters can improve with SW's