

# Scale Setting for Self-consistent Quantum Backgrounds

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- II Improving solutions
- III Self-consistent scale setting
- IV Example: “inoffensive”
- V Examples: “offensive”
- VI Summary and outlook

Based on arXiv:1409.4443 (accepted by PRD) and arXiv:1501.00904



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Introduction



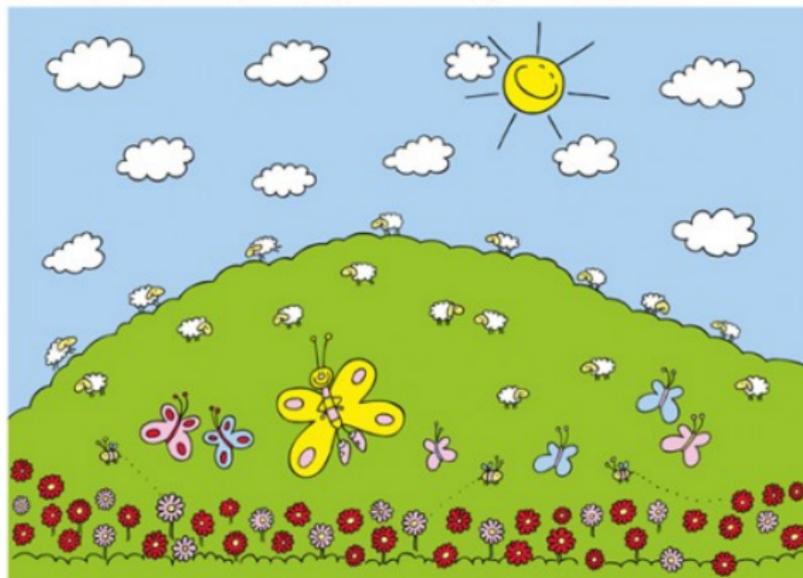
# I Introduction $\Gamma_k$

Ideal case

$$\Gamma_k(\phi)$$

(1)

Ein kleines Stück heile Welt für Dich...



In "Pleasantville"



# I Introduction $\Gamma_k$

## Effective action

Generating functional

$$Z[J] = \int D\varphi \exp \left( -i \int dx (L(\varphi) - \varphi J) \right) \quad (2)$$

Connected diagrams

$$W[J] = \ln Z[J] \quad (3)$$

Effective action (1PI)

$$\Gamma[\phi] = \sup_J \left( \int J\phi - W[J] \right) \quad (4)$$



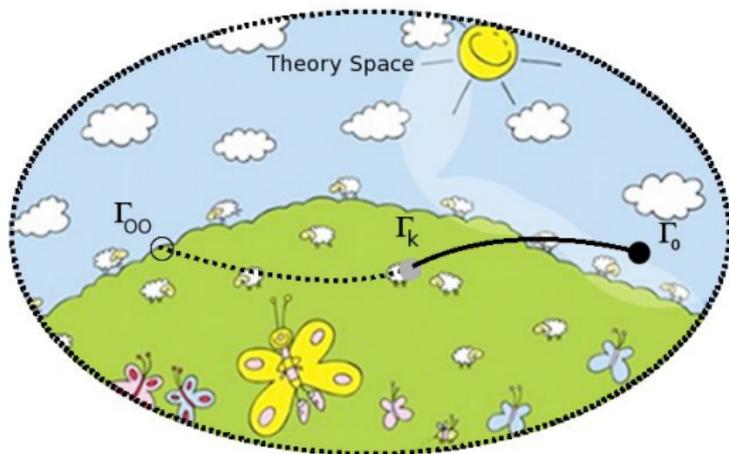
# I Introduction $\Gamma_k$

## Wilson's flow

What happens if one integrates out only certain Momentum shell ?

$$k \rightarrow k + \delta k$$

$\Rightarrow$  RG-Flow



$$\partial_k \Gamma_k = \dots \quad (5)$$

Solve:

Running couplings  $\lambda_k$  and  
effective action

$$\Gamma_k = \Gamma_k(\lambda_k, \phi)$$



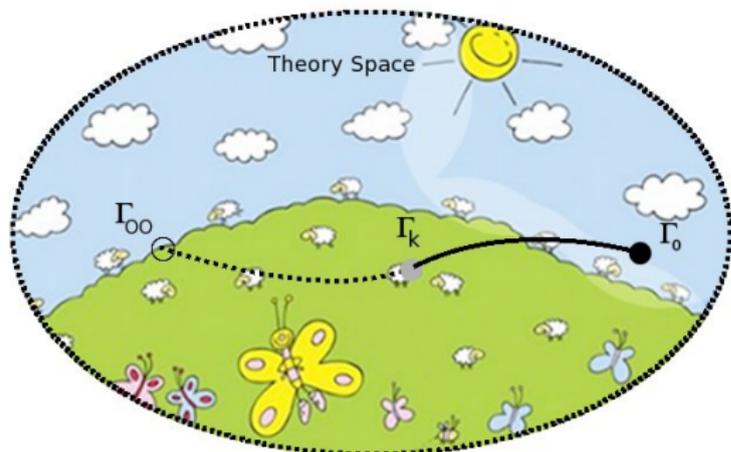
# I Introduction $\Gamma_k$

Wilson's flow dangers

gauge redundancy

$$\phi \rightarrow e^{i\alpha} \phi$$

Faddeev Popov

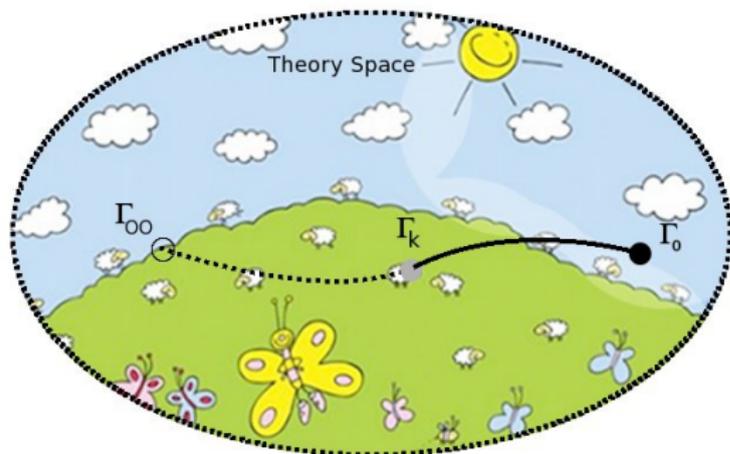


# I Introduction $\Gamma_k$

Wilson's flow dangers

Infrared divergencies

$$k \rightarrow 0, \Gamma_0 = ?$$

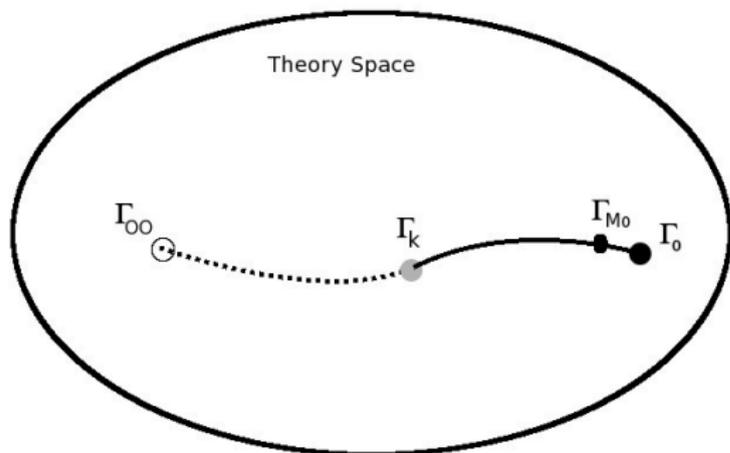


# I Introduction $\Gamma_k$

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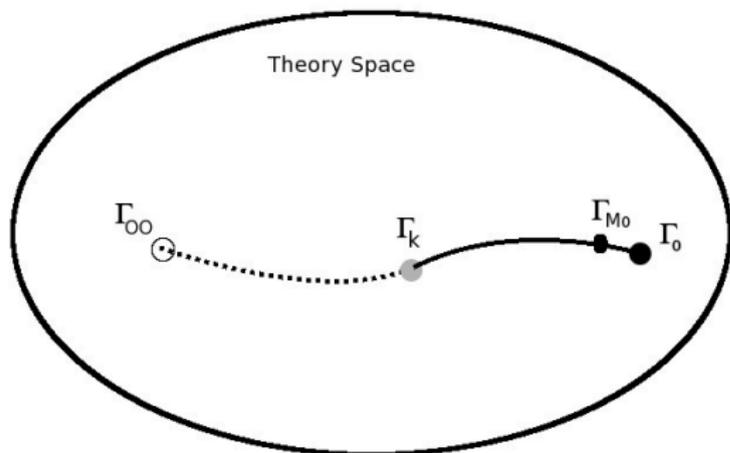


# I Introduction $\Gamma_k$

Wilson's flow dangers

Ultra violet divergencies

$$k \rightarrow \infty, \Gamma_\infty = ?$$

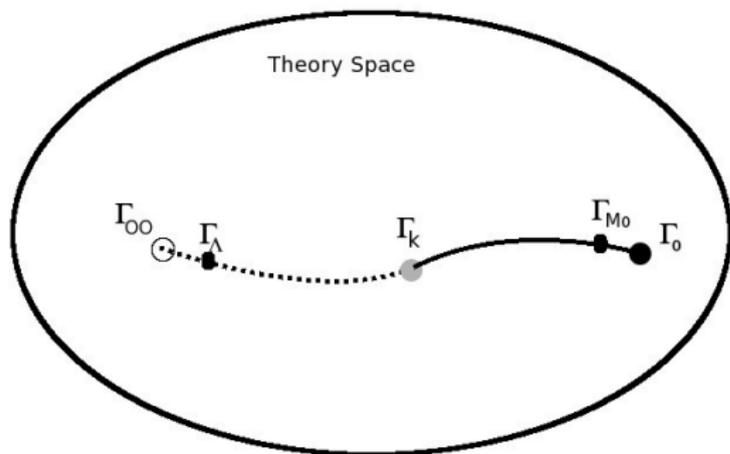


# I Introduction $\Gamma_k$

Wilson's flow dangers

Ultra violet divergencies

$$k \rightarrow \infty, \Gamma_\infty = ?$$



# I Introduction $\Gamma_k$

Wilson's flow dangers

Renormalization:  
Live with the problems



Absorb Infinities  $\sim \Lambda$  in  
definition of couplings  $\lambda$  at  
scale  $M_0$

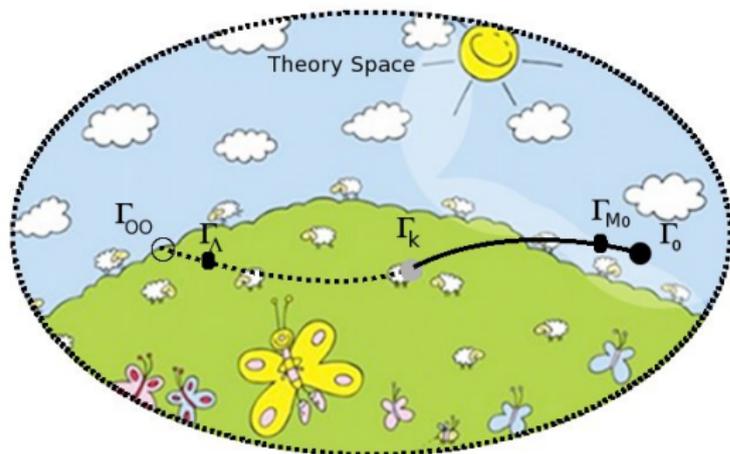
$$\lambda_{i,0} = \lambda_i(M_0)$$



# I Introduction $\Gamma_k$

## Wilson's flow after Renormalization

AFTER renormalization can  
pretend to be back in  
"Pleasantville"



Coupling flow  $\lambda_i(k)$  and  
effective action  $\Gamma_k$



# I Introduction $\Gamma_k$

## Wilson's flow

At the end effective quantum action:

$$\Gamma_k(\phi_i(x), \lambda_n(k)) = \int d^4x \sqrt{-g} \mathcal{L}_k(\phi_i(x), \lambda_n(k)) \quad (8)$$

**Quantum background?**

$$\frac{\delta \Gamma_k}{\delta \phi_i} = 0 \quad . \quad (9)$$

**"GAP EQUATION"**

solutions?

quantum solitons?



# I Introduction $\Gamma_k$

## Wilson's flow

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**Quantum background?**

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**"GAP EQUATION"**

solutions?

quantum solitons?



II

Improving solutions II

II

Improving solutions



# Improving solutions II

## Improving solutions II

$$\frac{\delta \Gamma_k}{\delta \phi_i} = 0 \text{ is too hard}$$

⇒ first solve classical eom

$$\frac{\delta S}{\delta \phi_i} = 0 \quad . \quad (10)$$

and take  $\lambda_0 \rightarrow \lambda_k$  as small correction of those solutions<sup>*i,ii*</sup>

*i)* Uehling potential: in QED textbooks

*ii)* Improved black holes: B.K. and Frank Saueressig; Int.J.Mod.Phys. A29 (2014) 8, 1430011 ...



# Improving solutions II

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# Improving solutions II

More general consideration

“Improved solution” - fully consistent?

No

Problem: Scale setting

$$k \rightarrow^? k(r) \quad (\text{typically } k \sim 1/r) \quad (11)$$

Implies

- Improved classical solution does not solve “gap equations”
- Typically not even  $T^{\mu\nu}$  conserved

⇒ Propose different scale setting



# III

## Self-consistent scale setting III

# III

## Self-consistent scale setting



# Self-consistent scale setting III

## Proposal

Effective action and running couplings

$$\Gamma_k(\phi_i(x), \lambda_n(k)) = \int d^4x \sqrt{-g} \mathcal{L}_k(\phi_i(x), \lambda_n(k)) \quad , \quad (12)$$

“Gap equations” for quantum background

$$\frac{\delta \Gamma_k}{\delta \phi_i} = 0 \quad . \quad (13)$$

Scale setting for this background?



# Self-consistent scale setting III

## Proposal

Scale setting for this background  $k = k(r)$ ?

Idea: Minimal  $k$  sensitivity

(just like Callan-Symanzik equations  $\frac{d}{dk} \langle T \phi_1(x_1) \phi_2(x_2) \dots \rangle_k \Big|_{k=k_{opt}} \equiv 0$ )

Realization: Promote  $k$  to field in  $\Gamma_k$

$$\Gamma_k(\phi_i(x), \lambda_n(k)) \rightarrow \Gamma(\phi_i(x), k(x), \lambda_n(k)) \quad . \quad (14)$$



# Self-consistent scale setting III

## Proposal

### Realization:

“Scale field”  $k(x)$

coupled equations of motion

$$\frac{\delta\Gamma(\phi_i(x), k(x))}{\delta\phi_i} = 0 \quad , \quad \left. \frac{d}{dk} \mathcal{L}(\phi_i(x), k(x), \lambda_n(k)) \right|_{k=k_{opt}} = 0 \quad . \quad (15)$$

The simultaneous solution of (15) assures optimal scale setting



# Self-consistent scale setting III

## Proposal

### First questions on proposal

- Consistent set of equations?
- Solves  $\nabla_\mu T^{\mu\nu} \leftrightarrow k(r)$  problem?
- Examples?



# IV

## Examples IV

# IV

## Examples: “inoffensive”



# IV

## Examples IV

Inoffensive:



# Examples IV

$\phi^4$

Effective action

$$\Gamma = \int d^4x \left( \frac{\alpha_k}{2} (\partial\phi)^2 - \frac{\tilde{m}_k^2}{2} \phi^2 - \frac{\tilde{g}_k}{4!} \phi^4 \right) \quad (16)$$

Eom  $\delta\phi$ :

$$\partial_\mu (\alpha_k \partial^\mu \phi) + \tilde{m}_k^2 \phi + \frac{\tilde{g}_k}{6} \phi^3 = 0 \quad (17)$$

eom  $k$ :

$$\alpha'_k (\partial\phi)^2 - (\tilde{m}_k^2)' \phi^2 - \frac{1}{12} \tilde{g}'_k \phi^4 = 0 \quad , \quad (18)$$

implies conservation

$$\nabla_\mu T^{\mu\nu} = 0$$

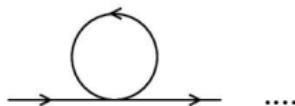


# Examples IV

$\phi^4$  loops

Example: Self-consistent scale setting for  $\phi^4$  at 1 loop:

$$\begin{aligned}\gamma_Z &= \frac{d \ln Z_k}{d \ln k^2} = 0 \quad , \\ \beta_g &= \frac{dg_k}{d \ln k} = \frac{3}{16\pi^2} g_k^2 \quad , \\ \beta_{m^2} &= \frac{dm_k^2}{d \ln k} = \left(-2 + \frac{g_k}{16\pi^2}\right) m_k^2 \quad .\end{aligned}$$



Run proposed machine for spherical symmetry ...

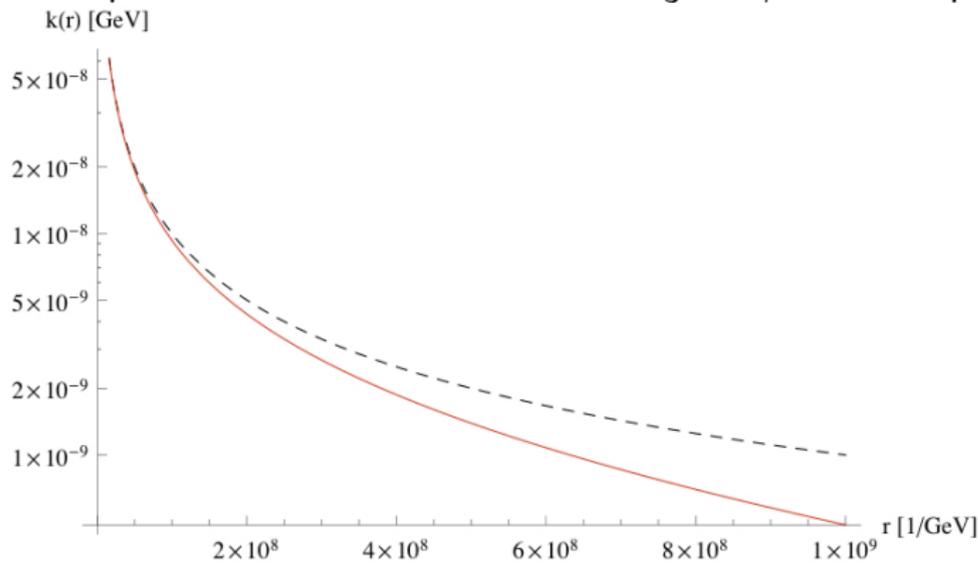
$$k(r) = k_i + \exp\left(-2\sqrt{\frac{14}{3}} \frac{k_0}{k_i} m_0 r\right) \cdot \frac{c_1}{r} \quad .$$



# Examples IV

## $\phi^4$ loops

Example: Self-consistent scale setting for  $\phi^4$  at 1 loop:



- Self-consistent
- Like classical for small  $r$



# IV

## Examples IV

V

Examples: “offensive”



# IV

## Examples IV

Offensive:



# Examples V

## Gravity and gauge

Effective action

$$\Gamma_k[g_{\mu\nu}, A_\alpha] = \int_M d^4x \sqrt{-g} \left( \frac{R - 2\Lambda_k}{16\pi G_k} - \frac{1}{4e_k^2} F_{\mu\nu} F^{\mu\nu} \right) , \quad (21)$$

Einstein-Hilbert-Maxwell



# Examples V

## Gravity and gauge

### Equations of motion

eom  $\delta g_{\mu\nu}$ :

$$G_{\mu\nu} = -g_{\mu\nu}\Lambda_k - \Delta t_{\mu\nu} + 8\pi G_k T_{\mu\nu} \quad , \quad (22)$$

with

$$\Delta t_{\mu\nu} = G_k (g_{\mu\nu}\square - \nabla_\mu \nabla_\nu) \frac{1}{G_k} \quad . \quad (23)$$

and

$$T_{\mu\nu} = F_\nu^\alpha F_{\mu\alpha} - \frac{1}{4} g_{\mu\nu} F_{\mu\nu} F^{\mu\nu} \quad . \quad (24)$$

eom  $\delta A_\mu$ :

$$D_\mu \left( \frac{1}{e_k^2} F^{\mu\nu} \right) = 0 \quad ,$$



# Examples V

## Gravity and gauge

### Conservation and symmetry

Symmetry coordintates:

$$\nabla^\mu G_{\mu\nu} = 0 \quad (26)$$

Symmetry  $U(1)$ :

$$\nabla_{[\mu} F_{\alpha\beta]} = 0 \quad . \quad (27)$$



# Examples V

## Gravity and gauge

Consistency:

Using everybody ... one can actually show that

Scale setting eom  $\delta k$ :

$$\left[ R \nabla_{\mu} \left( \frac{1}{G_k} \right) - 2 \nabla_{\mu} \left( \frac{\Lambda(k)}{G_k} \right) - \nabla_{\mu} \left( \frac{4\pi}{e_k^2} \right) F_{\alpha\beta} F^{\alpha\beta} \right] \cdot (\partial^{\mu} k) = 0 \quad . \quad (28)$$

is actually self-consistent consequence of eoms and conservation laws



# Examples V

## Gravity and gauge

Background solutions?

Actually, yes:



# Examples V

## Gravity and gauge

Background solutions?

Actually, yes:



# Examples V

## de Sitter black hole

De Sitter case:

eom  $g_{\mu\nu}$ :

$$G_{\mu\nu} = -g_{\mu\nu}\Lambda_k - \Delta t_{\mu\nu} \quad , \quad (29)$$

eom  $k$ :

$$R\nabla_\mu \left( \frac{1}{G_k} \right) - 2\nabla_\mu \left( \frac{\Lambda_k}{G_k} \right) = 0 \quad . \quad (30)$$

unknown functions:  $g_{00}(r)$ ,  $g_{11}(r)$ , and  $k(r)$

known (with **caveat**):  $\Lambda_k$ , and  $G_k$

Don't like **caveat** ...



# Examples V

## de Sitter black hole

Trade (trick):

$$\begin{aligned} g_{00}(r), g_{11}(r), k(r), \Lambda_k, \text{ and } G_k \\ \Rightarrow \\ g_{00}(r), g_{11}(r), \Lambda(r), \text{ and } G(r) \end{aligned}$$

Too many unknowns: Schwarzschild ansatz:  $g_{00} = 1/g_{11} \equiv f$

$$\begin{aligned} \Rightarrow \\ f(r), \Lambda(r), \text{ and } G(r) \end{aligned}$$

Can be solved



# Examples V

## de Sitter black hole

generalized de Sitter solution:

$$G(r) = \frac{G_0}{\epsilon r + 1} \quad (31)$$

$$f(r) = 1 + 3G_0 M_0 \epsilon - \frac{2G_0 M_0}{r} - (1 + 6\epsilon G_0 M_0)\epsilon r - \frac{\Lambda_0 r^2}{3} + r^2 \epsilon^2 (6\epsilon G_0 M_0 + 1) \ln \left( \frac{c_4(\epsilon r + 1)}{r} \right) \quad (32)$$

$$\Lambda(r) = \frac{-72\epsilon^2 r(\epsilon r + 1) \left(\epsilon r + \frac{1}{2}\right) \left(G_0 M_0 \epsilon + \frac{1}{6}\right) \ln \left( \frac{c_4(\epsilon r + 1)}{r} \right) + 4r^3 \Lambda_0 \epsilon^2 + (12\epsilon^3 + 6\Lambda_0 \epsilon + 72\epsilon^4 G_0 M_0) r^2}{2r(\epsilon r + 1)^2} \quad (33)$$
$$+ \frac{(72\epsilon^3 G_0 M_0 + 11\epsilon^2 + 2\Lambda_0) r + 6\epsilon^2 G_0 M_0}{2r(\epsilon r + 1)^2} .$$

Constants of integration:  $G_0$ ,  $M_0$ ,  $\Lambda_0$ ,  $\epsilon$ ,  $c_4$



# Examples V

## de Sitter black hole

Classical limit:

$$G(r)|_{\epsilon \rightarrow 0} = G_0 \quad (34)$$

$$\Lambda(r)|_{\epsilon \rightarrow 0} = \Lambda_0 \quad . \quad (35)$$

$$f(r)|_{\epsilon \rightarrow 0} = -\frac{\Lambda_0 r^2}{3} - \frac{2G_0 M_0}{r} + 1 \quad , \quad (36)$$

$\epsilon$  parametrizes scale dependence of the couplings



# Examples V

## de Sitter black hole

Asymptotics and horizons:

+  $r \rightarrow 0$

$$R_{\mu\nu\alpha\beta}R^{\mu\nu\alpha\beta} = \frac{48G_0^2M_0^2}{r^6} - \frac{48G_0^2M_0^2\epsilon}{r^5} + \mathcal{O}(r^{-4}) \quad . \quad (37)$$

$\Rightarrow$  Singularity persists

+  $r \rightarrow \infty$

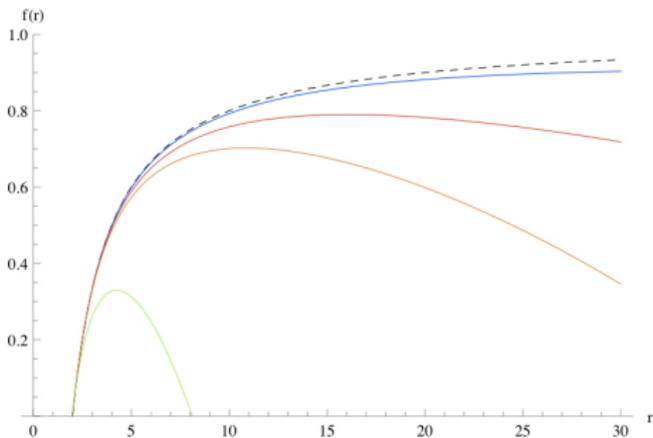
$$f(r) = -r^2 \frac{1}{3} (\Lambda_0 - 3\epsilon^2(6\epsilon G_0 M_0 + 1) \ln(c_4\epsilon)) + \mathcal{O}(r) \quad , \quad (38)$$

$\Rightarrow$  shift of effective cosmological constant

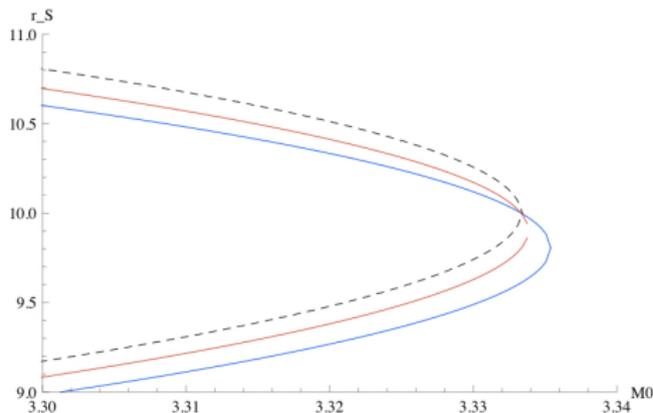


# Examples V

## de Sitter black hole



$f(r)$  for var.  $\epsilon$



Horizons for var.  $\epsilon$



# Examples V

## Reissner Nordstrom black hole

Reissner Nordstrom case:

eom  $g_{\mu\nu}$ :

$$G_{\mu\nu} = -\Delta t_{\mu\nu} + 8\pi \frac{G_k}{e_k^2} T_{\mu\nu} \quad , \quad (39)$$

eom  $A^\mu$

$$D_\mu \left( \frac{1}{e_k^2} F^{\mu\nu} \right) = 0 \quad . \quad (40)$$

eom  $k$ :

$$\left[ R \nabla_\mu \left( \frac{1}{G_k} \right) - \nabla_\mu \left( \frac{4\pi}{e_k^2} F_{\alpha\beta} F^{\alpha\beta} \right) \right] \cdot (\partial^\mu k) = 0 \quad . \quad (41)$$



# Examples V

## Reissner Nordstrom black hole

Unknown functions:

$$g_{00}(r) \equiv f(r), g_{11}(r), k(r), F_{01} = q(r),$$

Known

$$G_k, e_k$$

Same trade as before  $g_{11}(r) \equiv 1/f(r)$  and  $k(r)$  versus  $G(r), e(r)$ ,  
 $\Rightarrow$

$$f(r), q(r), G(r), e(r)$$

Can be solved ...



# Examples V

## Reissner Nordstrom black hole

Solution:

$$G(r) = \frac{G_0}{\epsilon r + 1} \quad (42)$$

$$f(r) = \frac{r^4 \epsilon^2 e_0^2 + 4\epsilon r^3 e_0^2 + 4(1 - G_0 M_0 \epsilon) e_0^2 r^2 - 8r G_0 M_0 e_0^2 + 16\pi G_0 Q_0^2}{4r^2 (\epsilon r + 1)^2 e_0^2}$$

$$e^2(r) = \frac{(r^6 \epsilon^4 e_0^2 + 3r^5 \epsilon^3 e_0^2 + (3r^4 e_0^2 - 4r^3 e_0^2 G_0 M_0 + 48r^2 \pi G_0 Q_0^2) \epsilon^2 + 48\epsilon r \pi G_0 Q_0^2 + 16\pi G_0 Q_0^2) e_0^2 \pi}{Q_0^2 G_0 (\epsilon r + 1)^3}$$

$$q(r) = \frac{Q_0}{4\pi e_0^2} \frac{e^2(r)}{r^2}$$

Integration constants:  $G_0$ ,  $Q_0$ ,  $e_0$ , and  $\epsilon$



# Examples V

## Reissner Nordstrom black hole

Classical limit:

$$f(r)|_{\epsilon \rightarrow 0} = 1 - \frac{2G_0 M_0}{r} + \frac{4\pi G_0 Q_0^2}{r^2 e_0^2} \quad , \quad (43)$$

$$e^2(r)|_{\epsilon \rightarrow 0} = (4\pi)^2 e_0^2 \quad (44)$$

$$q(r)|_{\epsilon \rightarrow 0} = \frac{4\pi Q_0}{r^2} \quad . \quad (45)$$

$\epsilon$  parametrizes scale dependence of the couplings (again)



# Examples V

## Reissner Nordstrom black hole

Asymptotics:

+  $r \rightarrow 0$

$$R_{\alpha\beta\gamma\delta}R^{\alpha\beta\gamma\delta} = 2^7 7 \frac{G_0^2 \pi^2 Q_0^4}{e_0^4 r^8} + \mathcal{O}(1/r^7) \quad . \quad (46)$$

$\Rightarrow$  Singularity persists

+  $r \rightarrow \infty$

line element

$$ds_\infty^2 = -\frac{1}{4} dt^2 + 4 dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \quad . \quad (47)$$

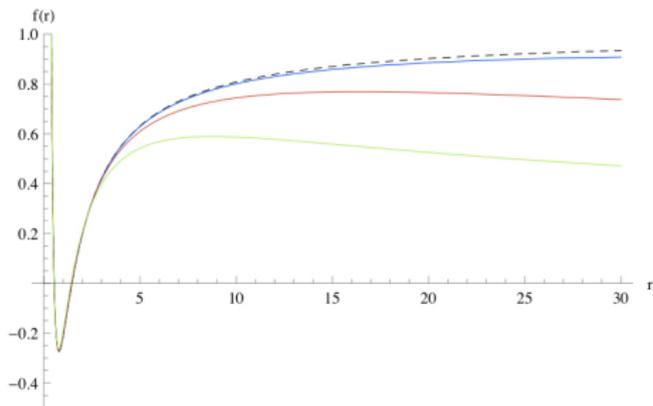
$\Rightarrow$  like global monopole (still charge is unchanged)

$$Q = \int_{\partial\Sigma} d^2z \sqrt{\gamma^{S_2}} n_\mu \sigma_\nu \frac{F^{\mu\nu}}{e^2} = \frac{Q_0}{e_0^2} \quad ,$$



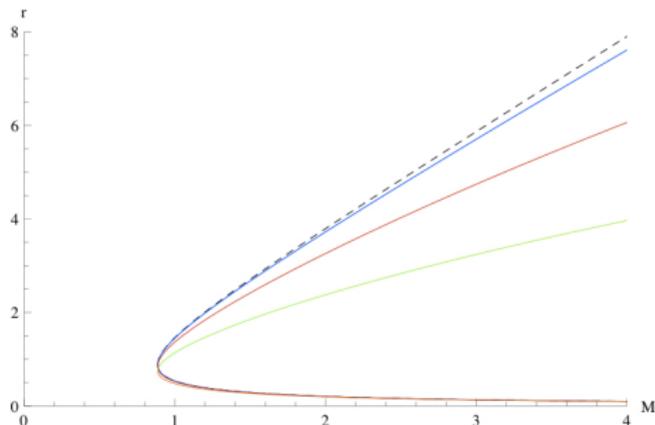
# Examples V

## Reissner Nordstrom black hole



$f(r)$  for var.  $\epsilon$

Observe that cosmic censorship unchanged



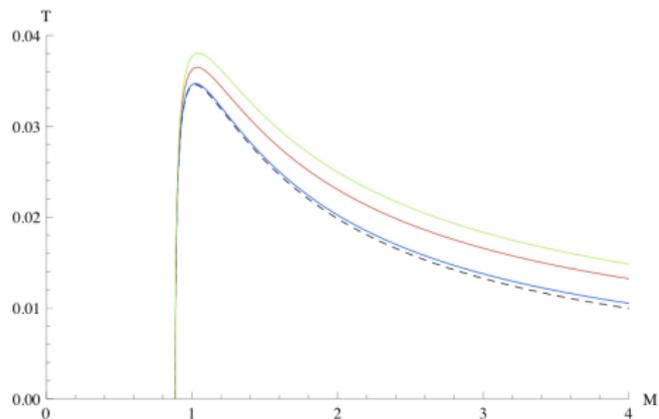
Horizons  $r_{\pm}(M_0)$  for var.  $\epsilon$

$$M_0 = 2\sqrt{\pi} \frac{Q_0}{e_0 G_0} ,$$



# Examples V

## Reissner Nordstrom black hole



Temperature slightly increased for  $\epsilon \neq 0 \sim \epsilon^2$

$$T = T_0 + \epsilon^2 \frac{G_0 \left( G_0 M_0^2 + M_0 \sqrt{\frac{G_0 (e_0^2 G_0 M_0^2 - 4\pi Q_0^2)}{e_0^2}} - 4\pi \frac{Q_0^2}{e_0^2} \right)}{8\pi \left( \sqrt{\frac{G_0 (e_0^2 G_0 M_0^2 - 4\pi Q_0^2)}{e_0^2}} + G_0 M_0 \right)} + O(\epsilon^3) .$$



V

Summary

# V Summary



# Summary

- **Context:** Self-consistent quantum backgrounds  $\delta\Gamma_k/\delta\phi_i = 0$
- **Problem:** Scale setting and consistency
- **Proposal:** Self-consistent scale setting  $k \rightarrow k(x)$   
 $\delta\Gamma_k/\delta\phi_i = 0$  and  $\delta\Gamma_k/\delta k = 0$
- **Examples:** Studied  $\phi^4$ , and de Sitter-, RN black holes
- **Outlook:** Hopefully to be applied in many more contexts ;-)



Thank you

