

Lieb-Robinson bounds & nonequilibrium dynamics in trapped ions

Phys. Rev. Lett. **111**, 230404 (2013)

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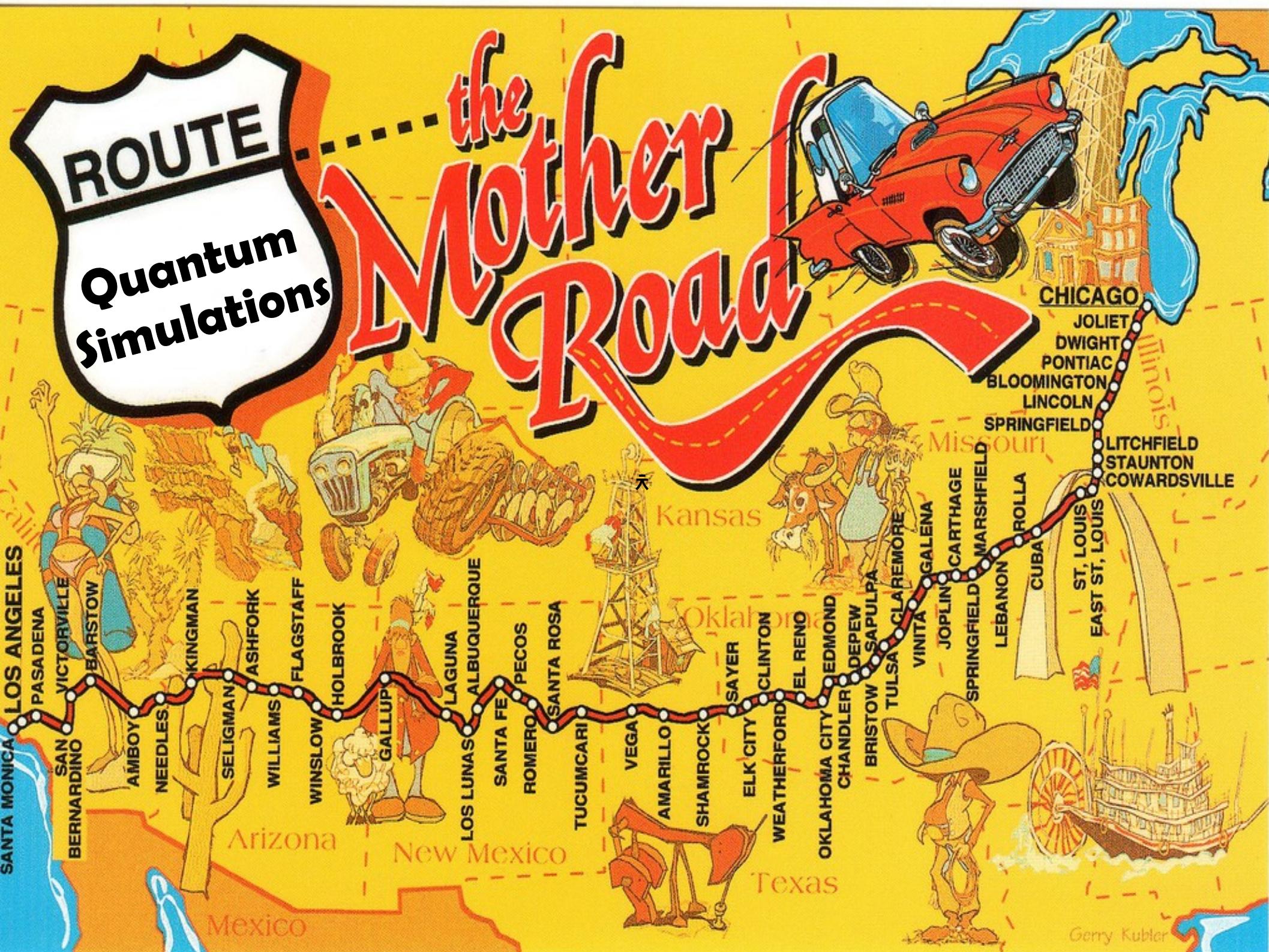
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ROUTE

Quantum
Simulations

the Mother Road



To-do

- Experimental tools
- Theoretical tools
- Fundamental questions
- Concrete applications

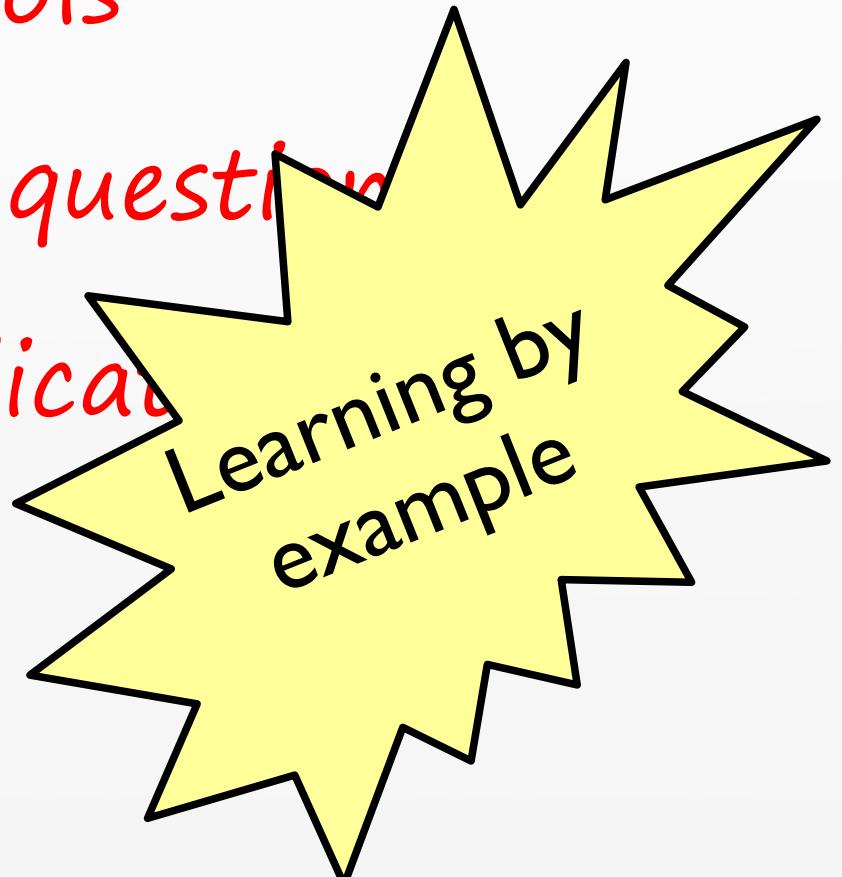
To-do

Experimental tools

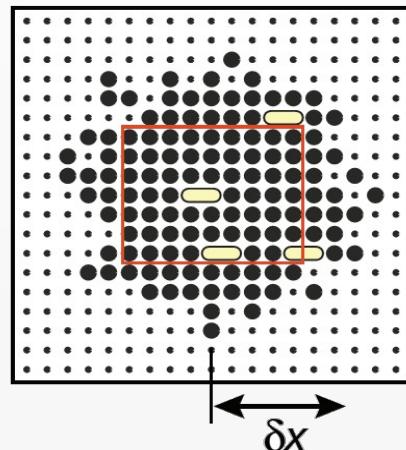
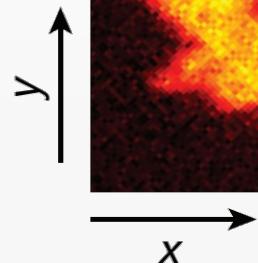
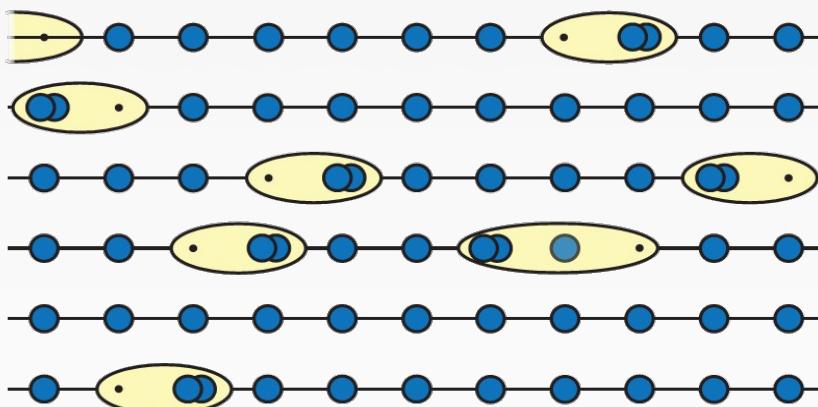
Theoretical tools

Fundamental questions

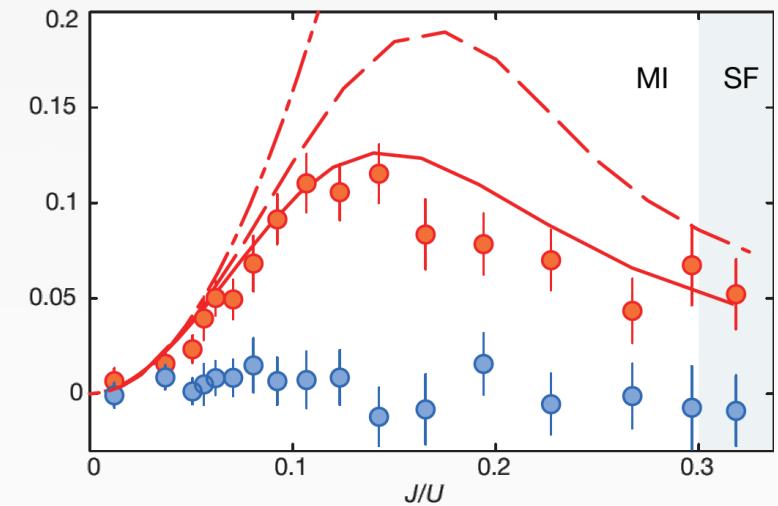
Concrete applications



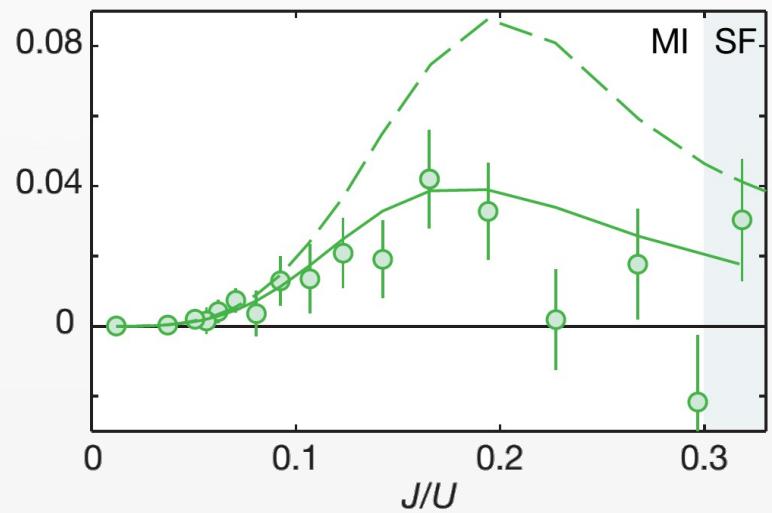
Static correlations & order



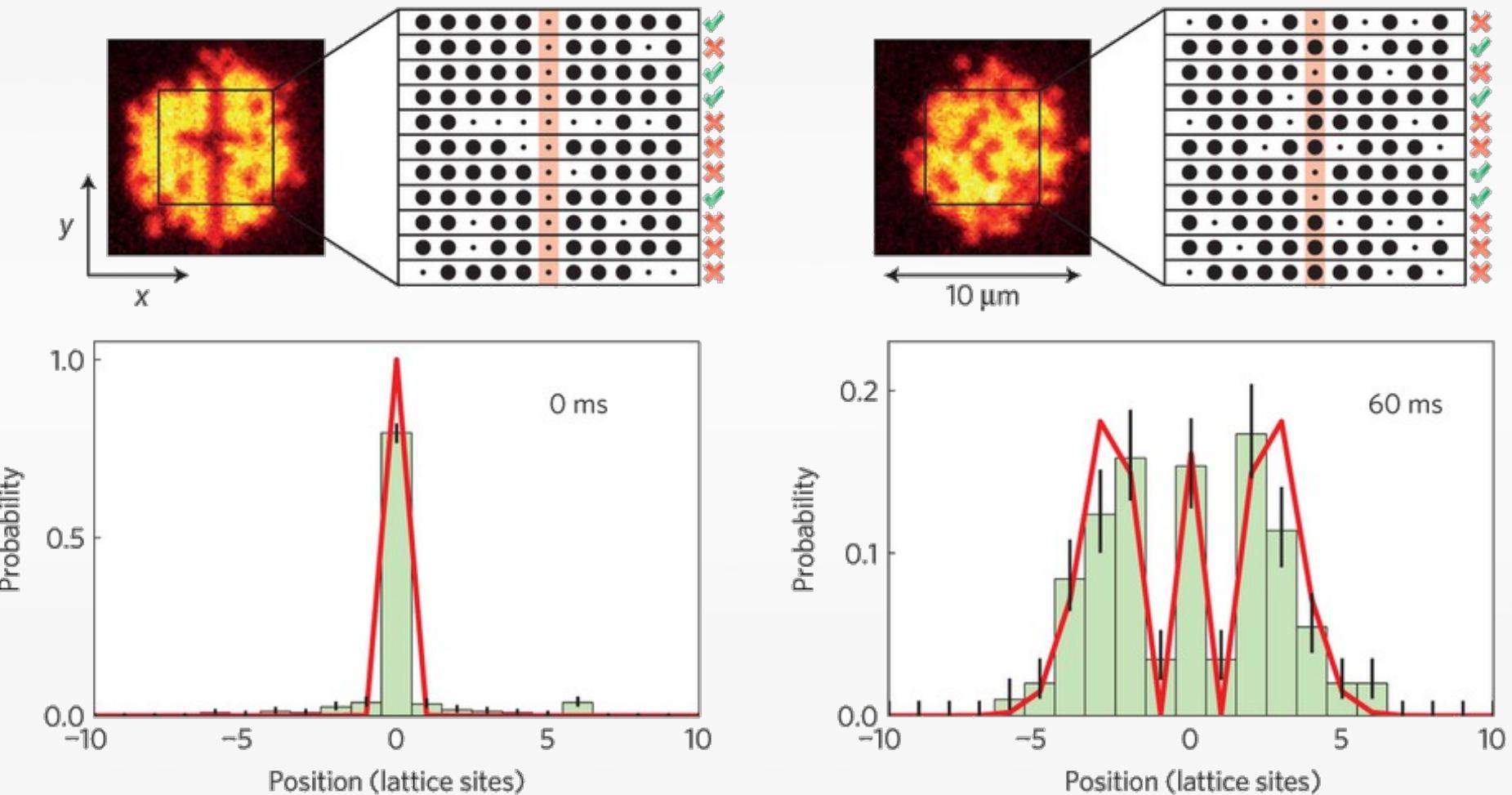
$$\langle S_i^z S_{i+1}^z \rangle$$



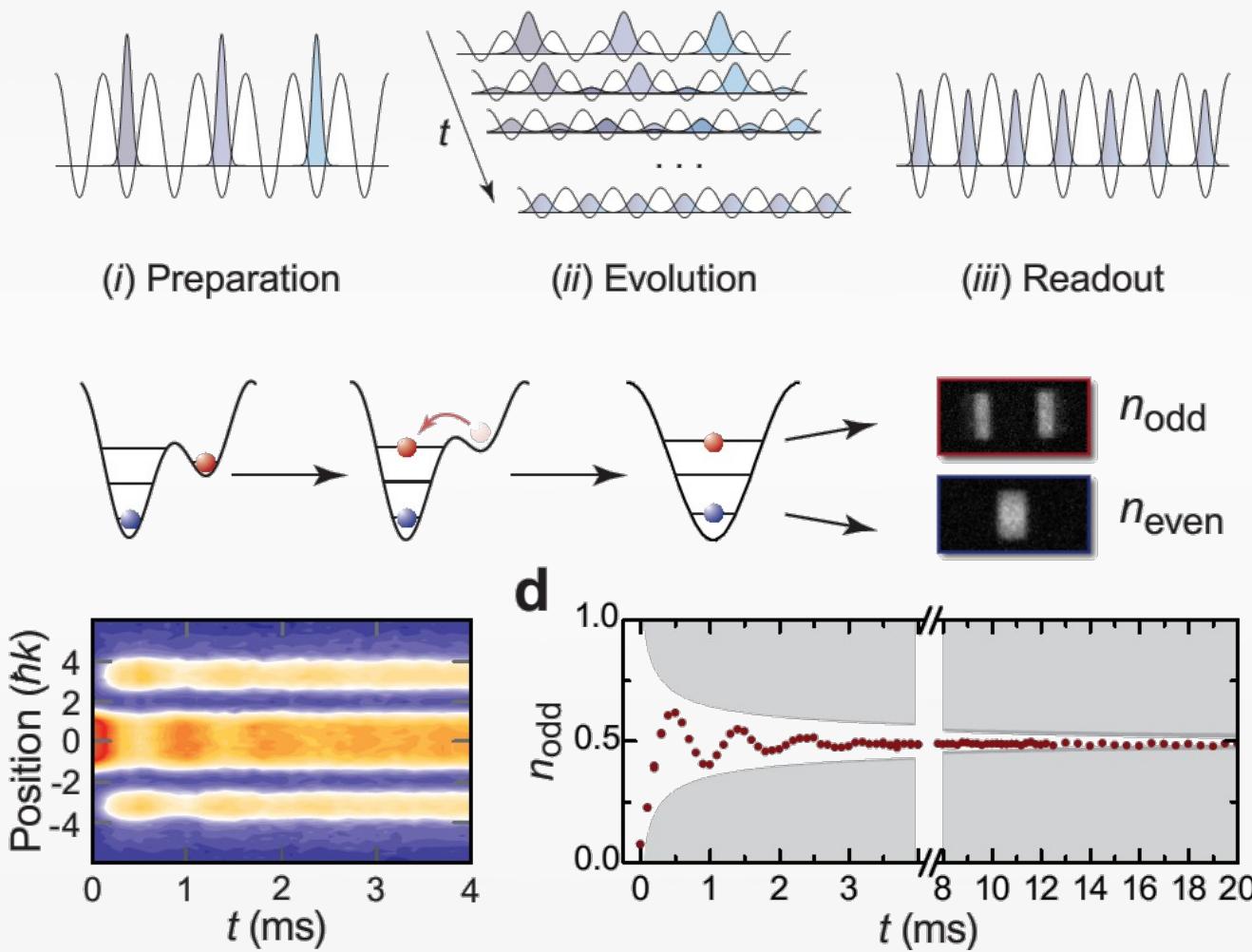
$$\langle S_1^z S_2^z S_3^z \rangle$$



Quasiparticle dynamics



Equilibration



CMP Problems

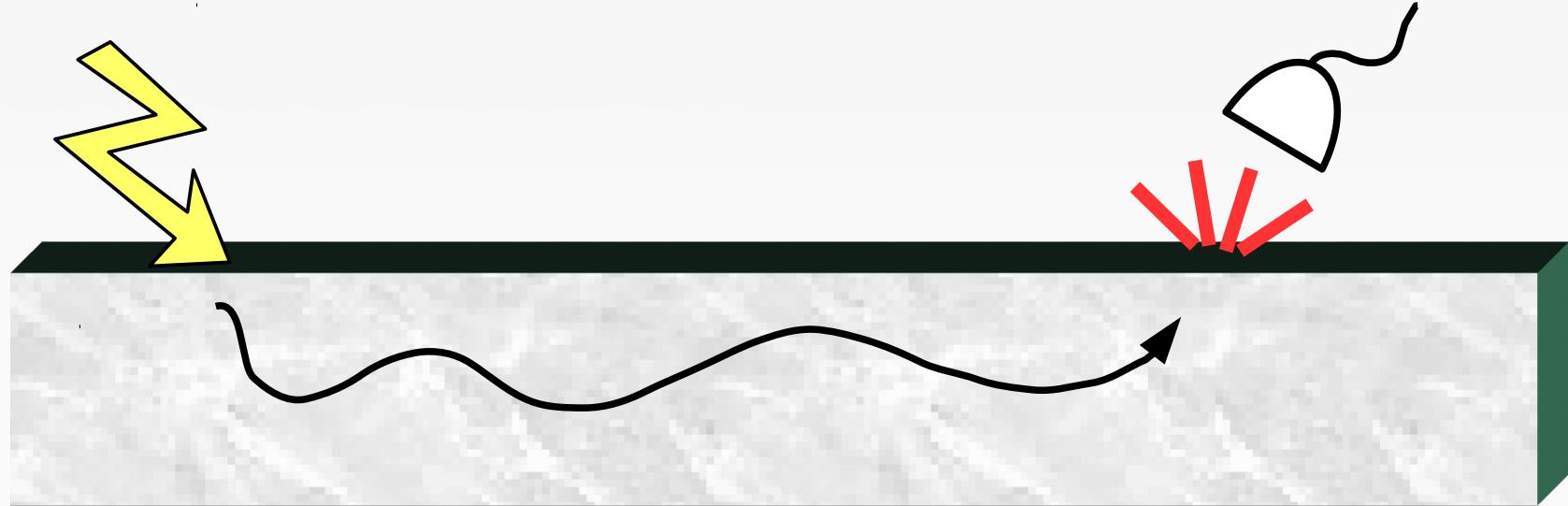
Equilibrium

{ Correlations
Entanglement
Topological order
Phase transitions

Nonequilibrium

{ Energy gaps
Excitations / quasiparticles
Equilibration / thermalization
Correlation propagation
Causality

Linear response

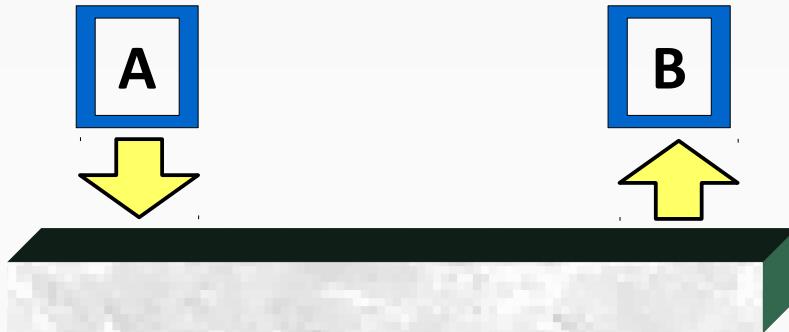


$$|\Psi_0\rangle \rightarrow \left(1 - i \frac{\varepsilon}{\hbar} A(0)\right) |\Psi_0\rangle + O(\varepsilon^2)$$

$$\langle B(t) \rangle \rightarrow \langle B(t) \rangle + \frac{\varepsilon}{\hbar} \langle [B(t), A(0)] \rangle + O(\varepsilon^2)$$

Entanglement

Mutual influence

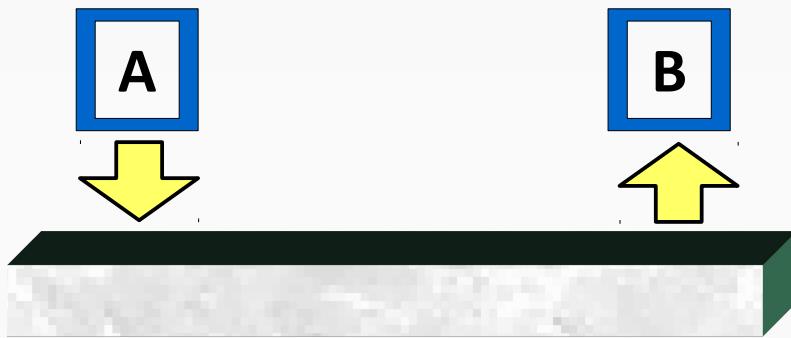


Vacuum fluctuations

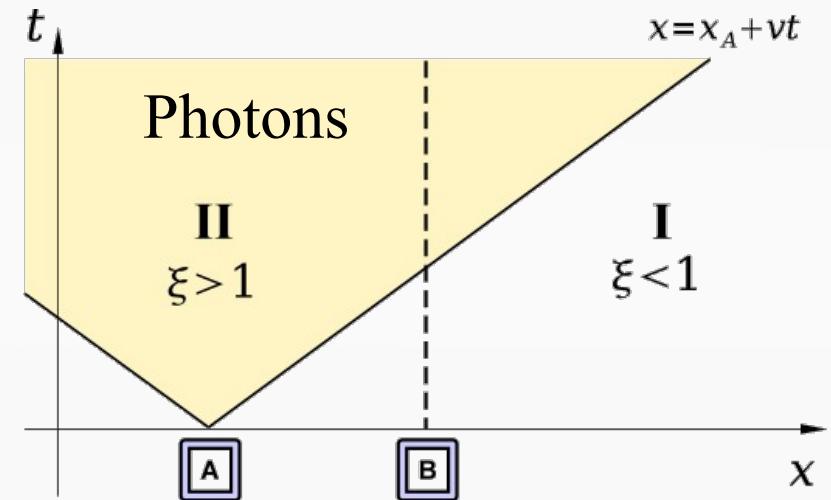


Entanglement

Mutual influence



Causality



$$\begin{aligned} \|[A(x,t), B(0,0)]\| &= 0 \\ |x| &> v|t| \end{aligned}$$

Effective causality

Approximate

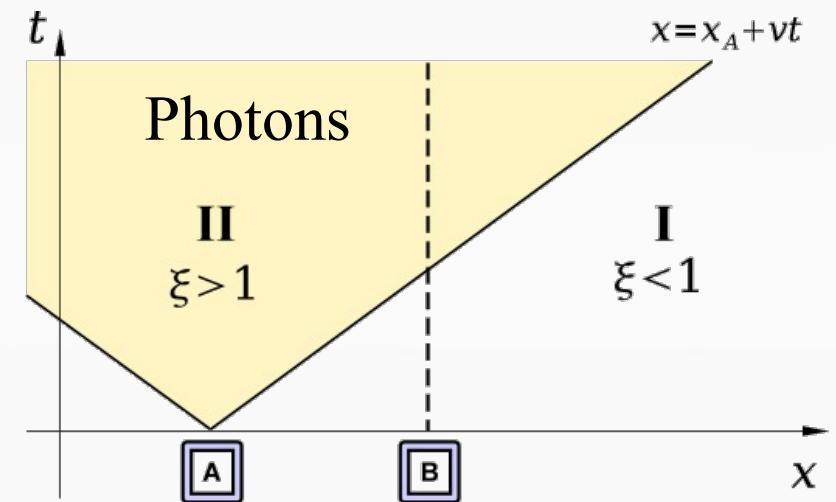
Non-relativistic, many-body system

$$H = \sum_{\langle i, j \rangle} \sigma_i \sigma_j$$



$$\| [A(x, t), B(0, 0)] \| \leq c \exp(\mu(v)t) \quad \forall |x| > v|t|$$

Causality

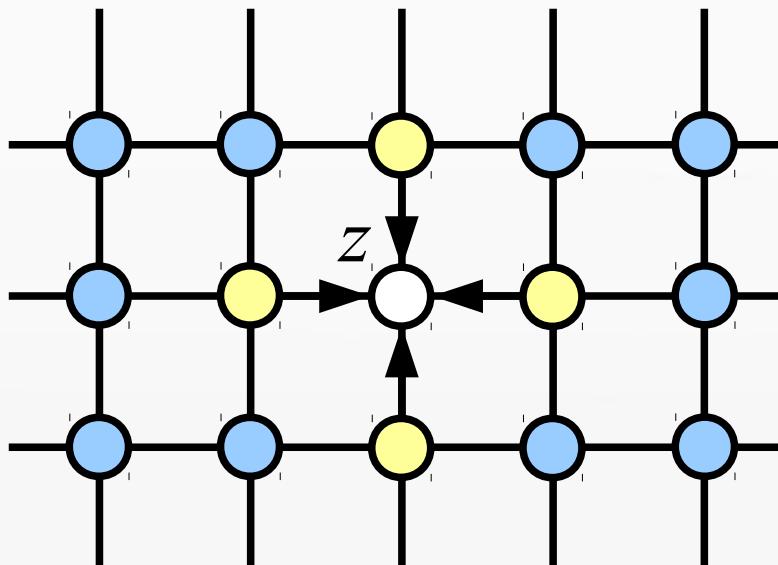


$$\| [A(x, t), B(0, 0)] \| = 0 \quad |x| > v|t|$$

Spin models

Hopping of particles:

$$H = \sum_{\langle i, j \rangle} J_{ij} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+)$$



Hopping from the z nearest neighbors is slower than

$$z \times \max_{\langle i, j \rangle} |J|$$

Thus the group velocity is smaller than that.

Basic idea

Starting point

$$\frac{d}{dt} [B_X(t), A_Y] = [[B_X(t), H], A_Y]$$

with a local model

$$H = \sum_{\langle x, y \rangle} H_{xy}$$

*M. W. Hastings, “Locality in Quantum Systems”, Les Houches Notes
E. H. Lieb, D. W. Robinson, Commun. Math. Phys. 28, 251-257 (1972)*

Basic idea

Starting point

$$\frac{d}{dt} [B_X(t), A_Y] = [[B_X(t), H], A_Y]$$

Quantity to bound

$$C(t, X, Y) = \|[B_X(t), A_Y(t)]\|$$

Recurrence relation

$$C(t, X, Y) \leq \sum_{\langle Z, X \rangle} \|H_{XZ}\| \|B_X\| \int C(t, Z, Y) dt$$

*M. W. Hastings, “Locality in Quantum Systems”, Les Houches Notes
E. H. Lieb, D. W. Robinson, Commun. Math. Phys. 28, 251-257 (1972)*

Basic idea

Starting point

$$\frac{d}{dt} [B_X(t), A_Y]$$

Quantity to bound

$$C(t, X, Y)$$



Recurrence relation

$$C(t, X, Y) \leq \sum_{\langle Z, X \rangle} \|H_{XZ}\| \|B_X\| \int C(t, Z, Y) dt$$

Beyond here

What has been done

- Harmonic oscillators
 - Bounded nonlinearity
 - Generalized oscillators
- Dissipative models
- Bounds on selected observables

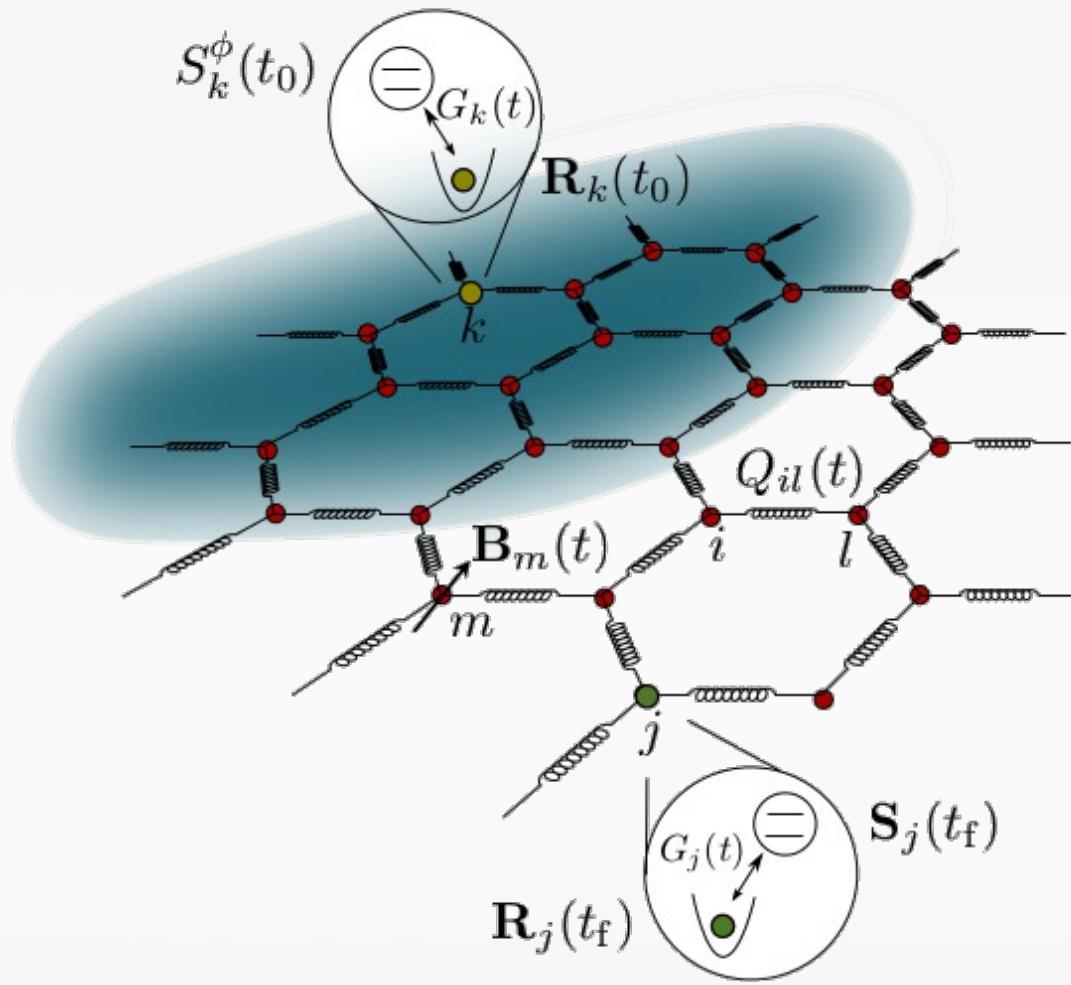
What for

- Condensed Matter
- Proof of exponential clustering of correlations
 - Area law
 - Correlation length
- Proof of MPS simulability
- Generalized adiabatic theorem.
- Proof of robustnes of topological order

Lieb-Robinson bounds for spin-boson lattice models and trapped ions

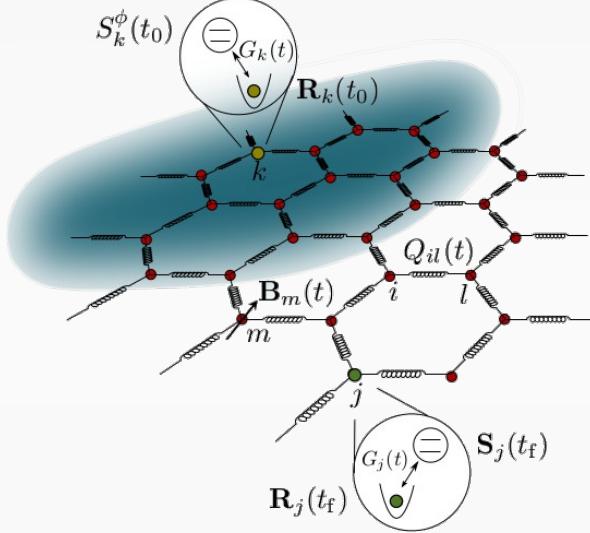
*J. Jünemann, A. Cadarso, D. Pérez-García, A. Bermúdez, JJGR
Phys. Rev. Lett. 111, 230404 (2013)
arXiv:1307.1992*

Spin-boson model



$$H = \sum_m \frac{\Delta_m}{2} \sigma_m^z + \sum_k \omega_k a_k^+ a_k + \sum_m g_m(t) \sigma_m^z x_m$$

Spin-boson model

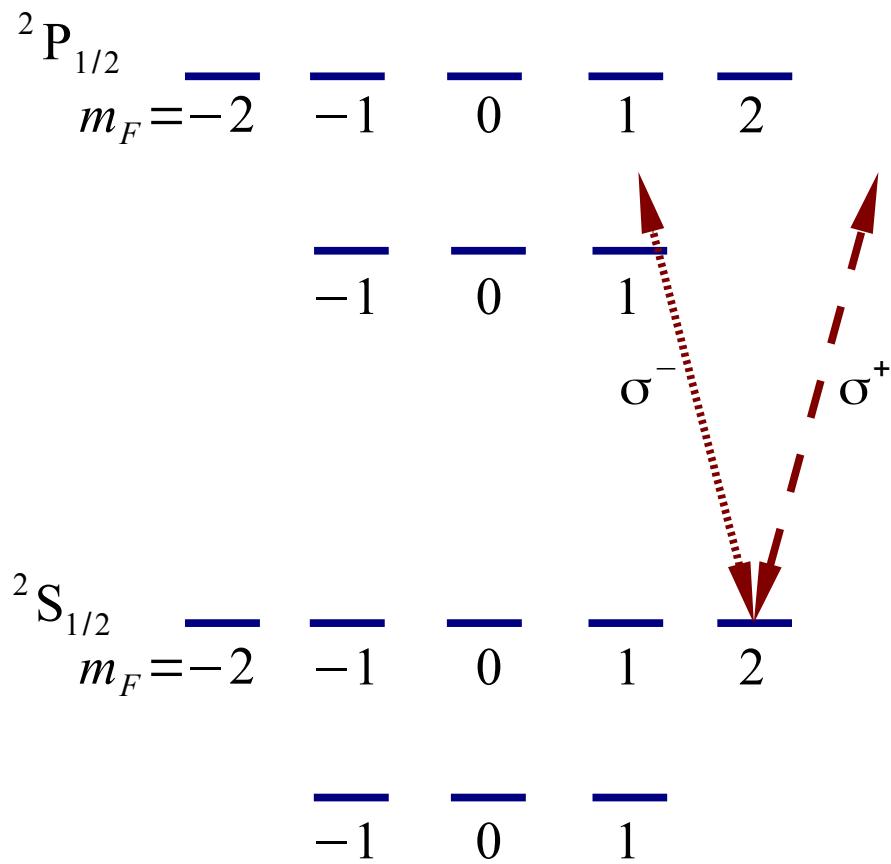


$$H = \sum_m \frac{\Delta_m}{2} \sigma_m^z + \sum_k \omega_k a_k^+ a_k + \sum_m g_m(t) \sigma_m^z x_m$$

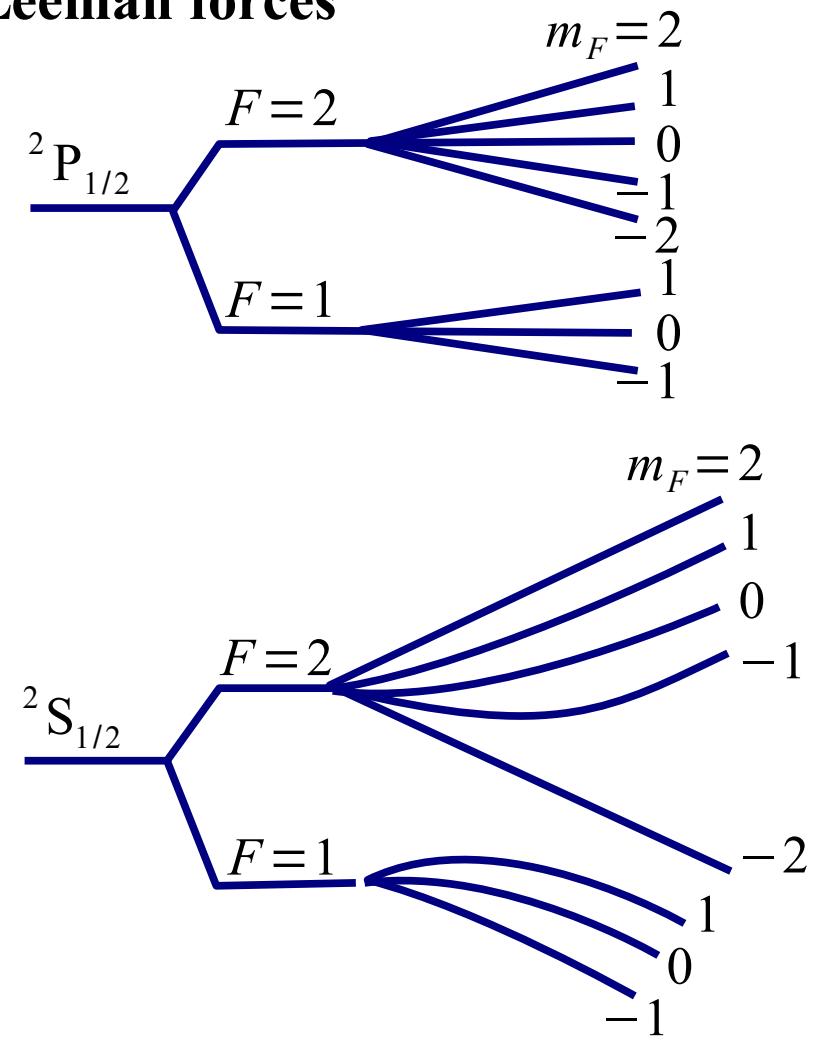
- Two models with LR bounds
- **Infinite-dimensional**
- **Highly nonlinear**
- Non-integrable
- Multiple experimental applications
 - Trapped ions, c-QED, nanophotonics, etc.

State dependent forces

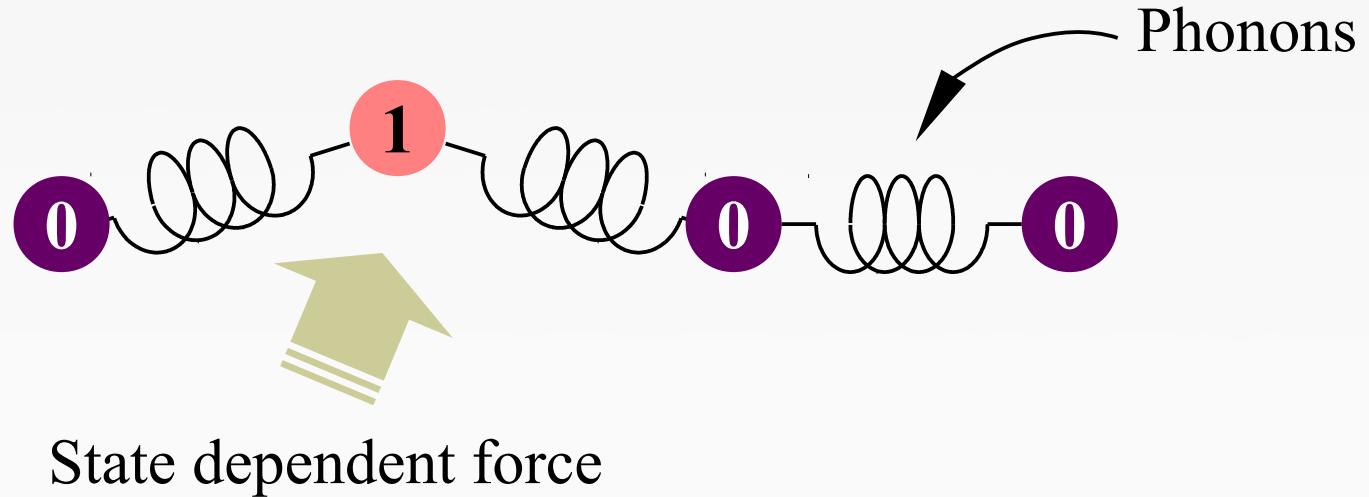
Optical (AC-Stark) forces



Zeeman forces



Spin-boson model



$$H = \sum_k \omega_k a_k^+ a_k + \sum_m \frac{\Delta_m}{2} \sigma_m^z + \sum_m F_m(t) \sigma_m^z x_m$$

Balance between Coulomb force and trap creates a crystal

Light pushes the atoms depending on the state

Effective interactions

Accumulated phase can be computed

$$\varphi = \sum_{i,j} \sigma_i^z \sigma_j^z \times \int_0^T \int_0^{t_1} F_i(t_1) G_{ij}(t_1 - t_2) F_j(t_2) dt_1 dt_2$$

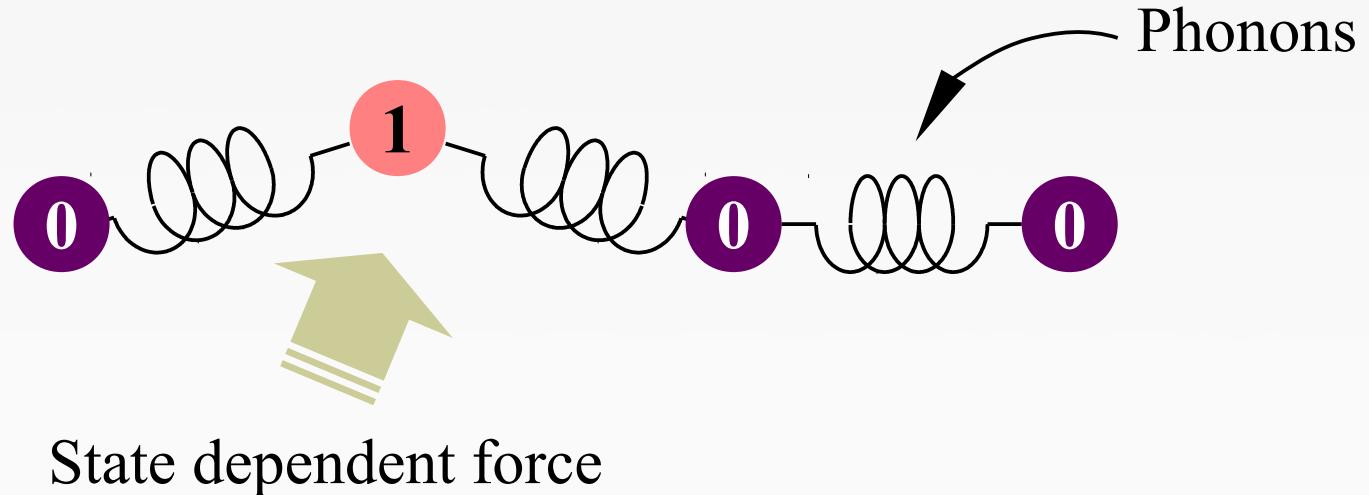
$$G_{ij}(t) = \sum_{ij} \frac{M_{ik} M_{jk}}{m \omega_k \hbar} \sin(\omega_k t)$$

This may be recasted as an spin-spin interaction

$$H_{eff} \sim \sum_{i,j} \sigma_i^z J_{ij} \sigma_j^z$$

which can be used for quantum simulation, QIPC, entanglement

Quantum simulation



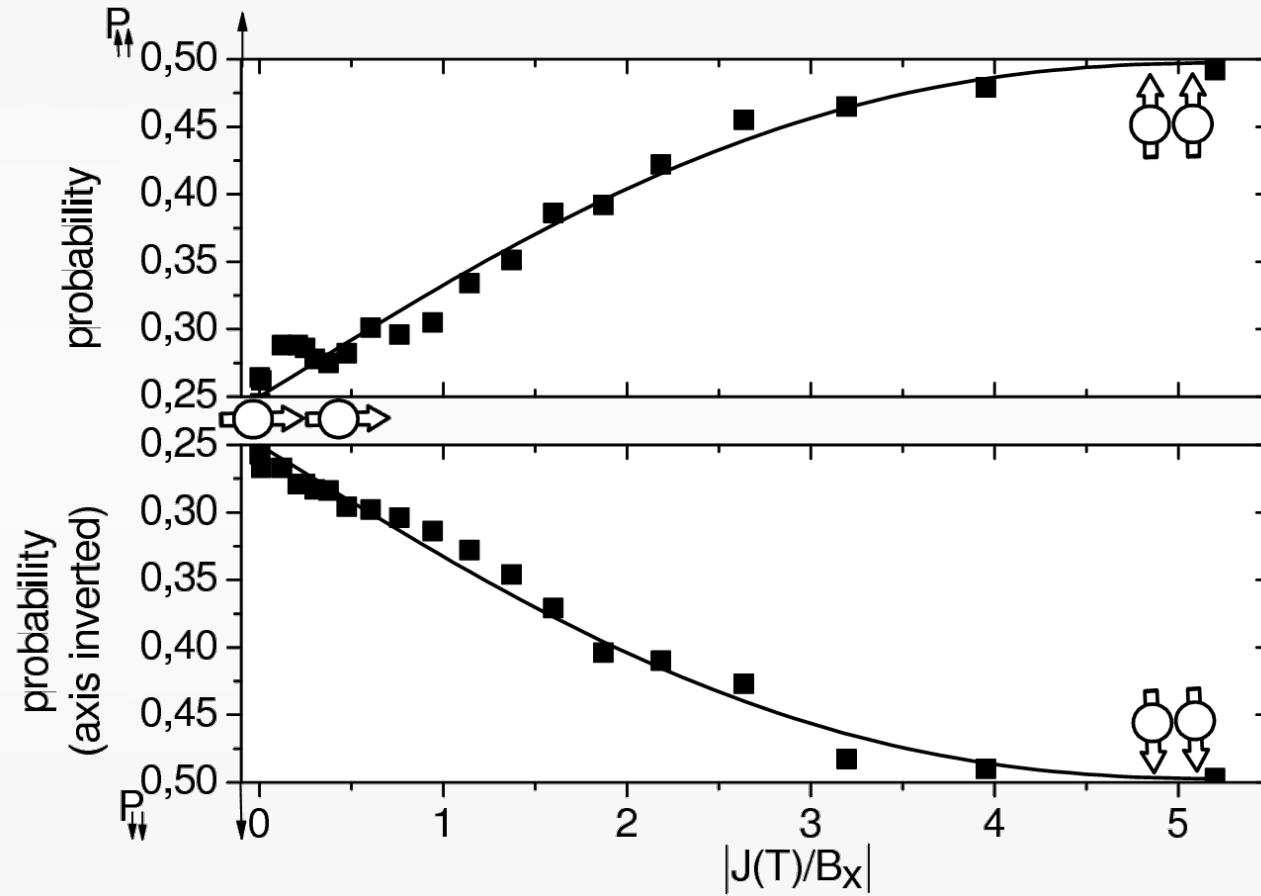
$$H = \sum_k \omega_k a_k^+ a_k + \sum_m \frac{\Delta_m}{2} \sigma_m^z + \sum_m F_m(t) \sigma_m^z x_m$$



$$F, \dot{F} \ll \omega, \Delta; \langle n_i \rangle \approx 0$$

$$H_{eff} \sim \sum_{i,j} \sigma_i^z J_{ij} \sigma_j^z$$

Quantum simulation

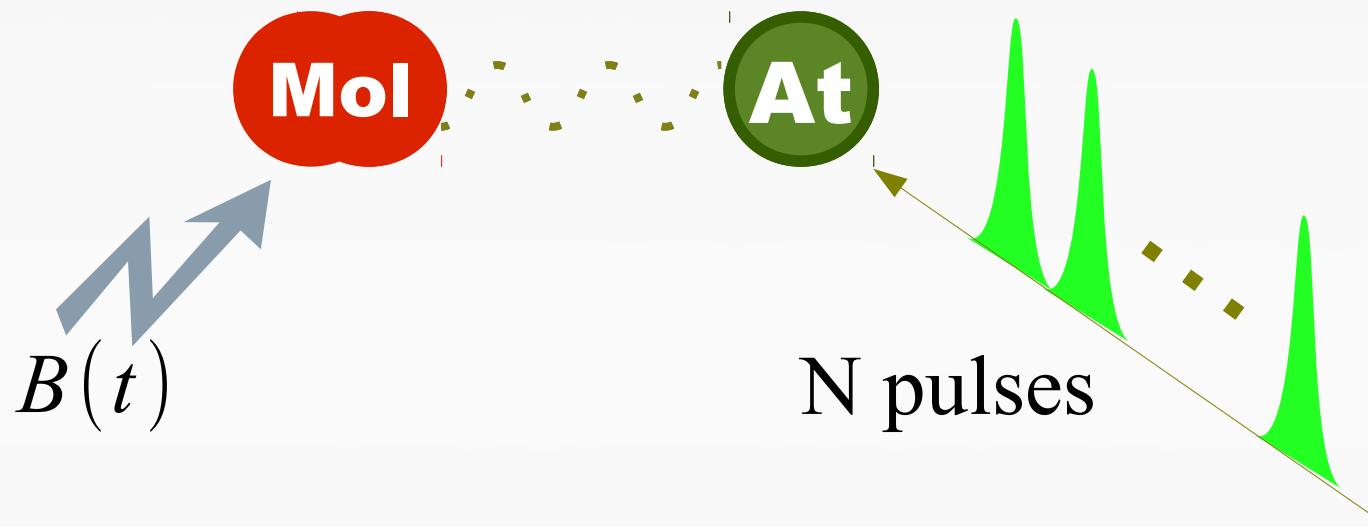


A. Friedenauer et al, *Nature Physics* 4, 757 - 761 (2008)

R. Islam et al, *Nature Comm.* 2, 377 (2011); *ibid, Science* 340 (2013)

J. W. Britton et al, *Nature* 484, 489 (2012)

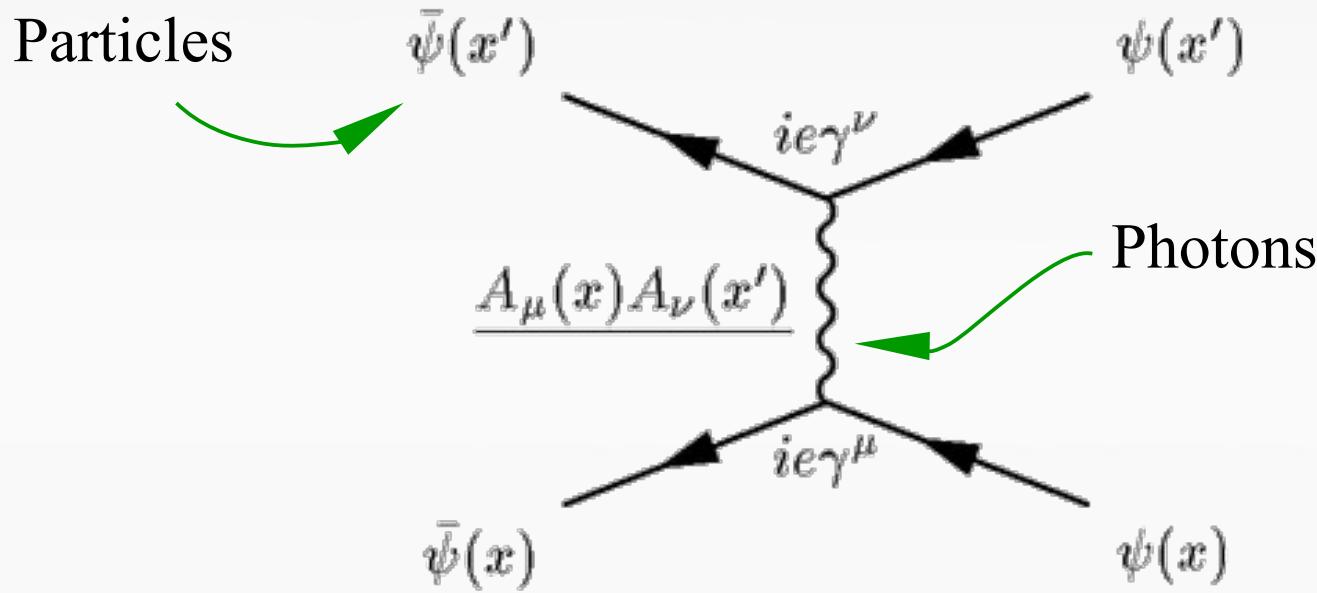
Signal intensification



Different ions may experience different forces

$$\varphi = \text{Im} \int_0^T \int_0^t G(|t - t'|) f_1(t) f_2(t') dt dt'$$
$$\propto B \times N$$

Boson-mediated interactions



A very general framework: c-QED, ions, photonics,...

When the dynamics of the bosons is as fast as the particle-boson interaction, we cannot eliminate them perturbatively.

Proof sketch

Many intermediate results:

- Proof of a LR for the harmonic oscillator model with and without long-range interactions
 - Tighter bounds than existing literature
- "Interaction-like" picture removes the influence of the oscillators in the ion's local dynamic.
- Integrate out the bosonic d.o.f. with the propagator.
- Sum up the resulting recurrence only for the spins.
- Feed back to the oscillators for further bounds.

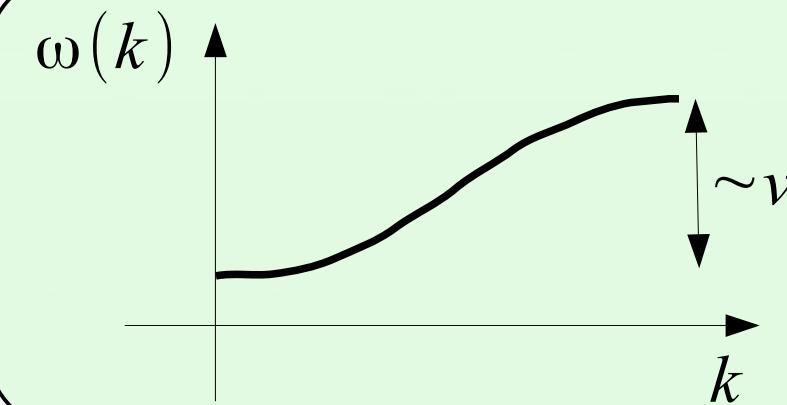
Two theory results

#1 An improved L-R bound for the ion phonons

$$\|[R_i(t), R_j(0)]\| \leq \frac{\exp(vt)}{d_{i,j}^\eta}$$

Long range interaction

$$H_{\text{int}} \sim \sum_{ij} \frac{1}{|x_i - x_j|^3}$$



Two theory results

#2 A resulting L-R bound for the interacting ions / spins

$$\|[\sigma_i(t), \sigma_j(0)]\| \leq \frac{\exp(vt)}{d_{i,j}^\eta} \left(e^{g^2 \frac{t}{v}} - 1 \right)$$

Bosonic LR repeated

$$\| [R_i(t), R_j(0)] \| \leq \frac{\exp(vt)}{d_{i,j}^\eta}$$

Speed of information
exchange

coupling g vs speed v

Bound saturation

#2 A resulting L-R bound for the interacting ions / spins

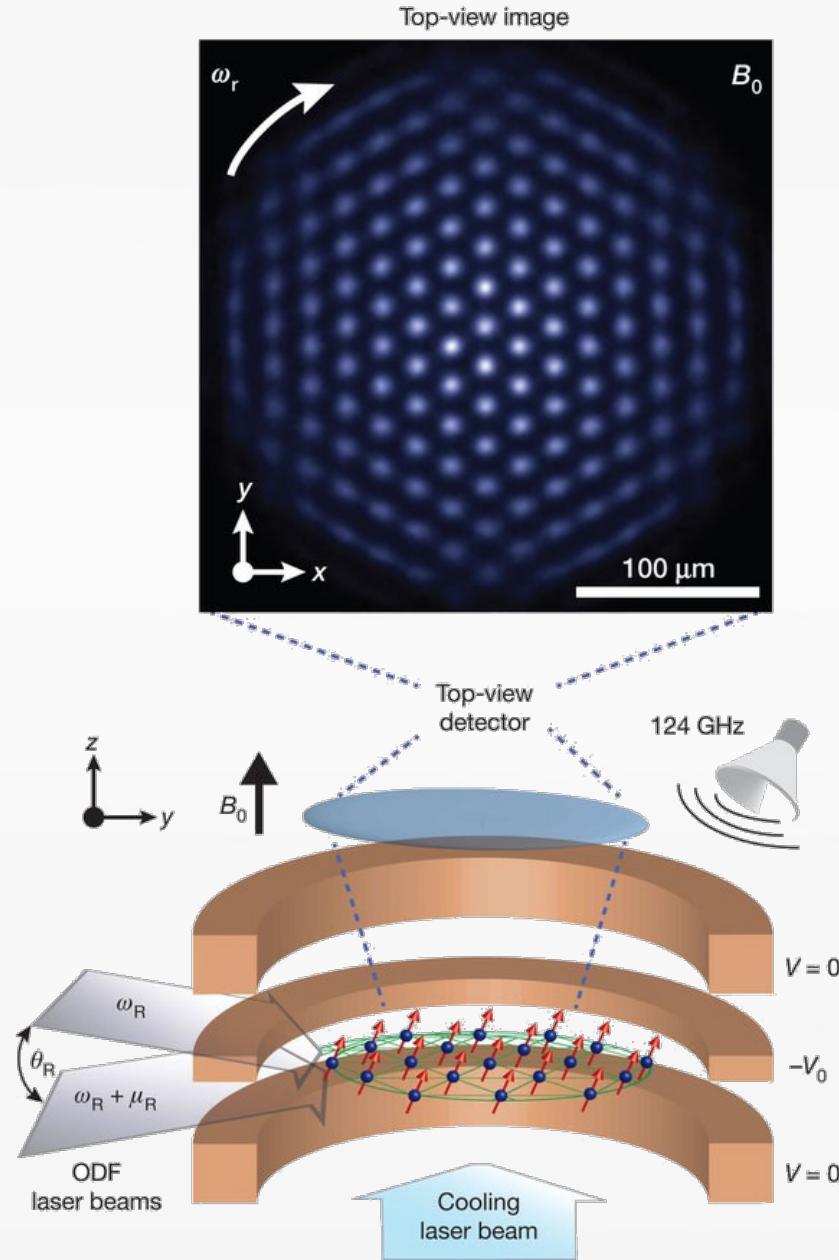
$$\|[\sigma_i(t), \sigma_j(0)]\| \leq \frac{\exp(\nu t)}{d_{i,j}^\eta} \left(e^{\cancel{g^2 \frac{t}{\nu}}} - 1 \right)$$

The bound is saturated in the
impulsive regime

(When the forces act instantaneously)

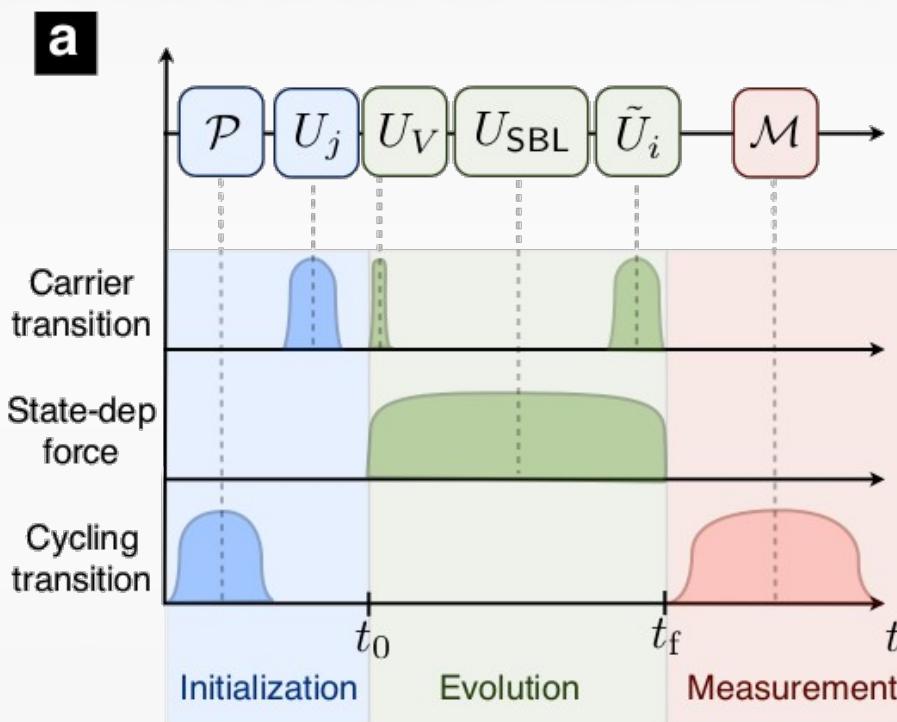
Trapped ions

*J. Britton et al,
Nature 484, 489–492
(2012)*

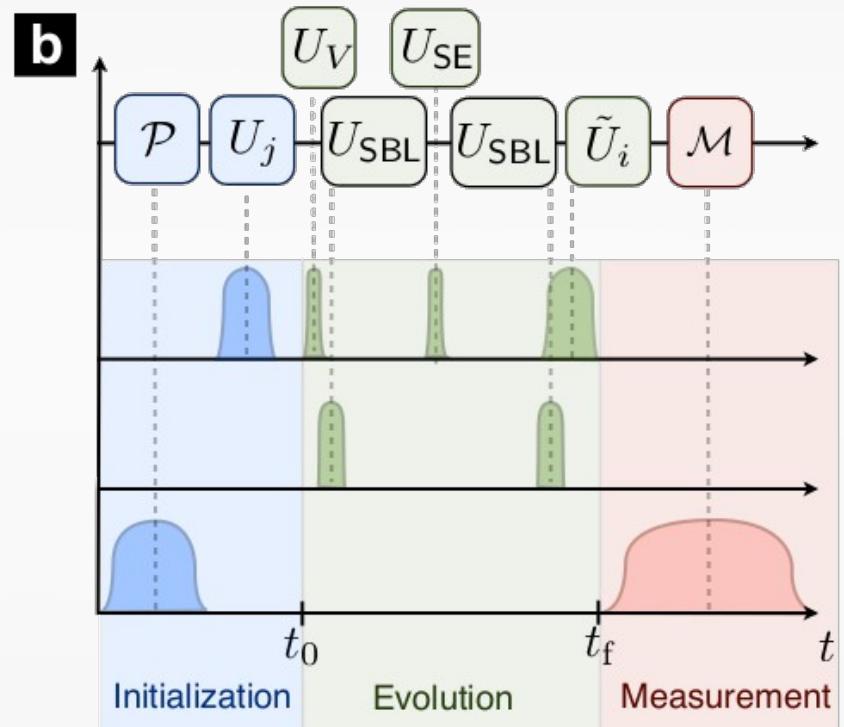


Measurement protocol

Continuous forces

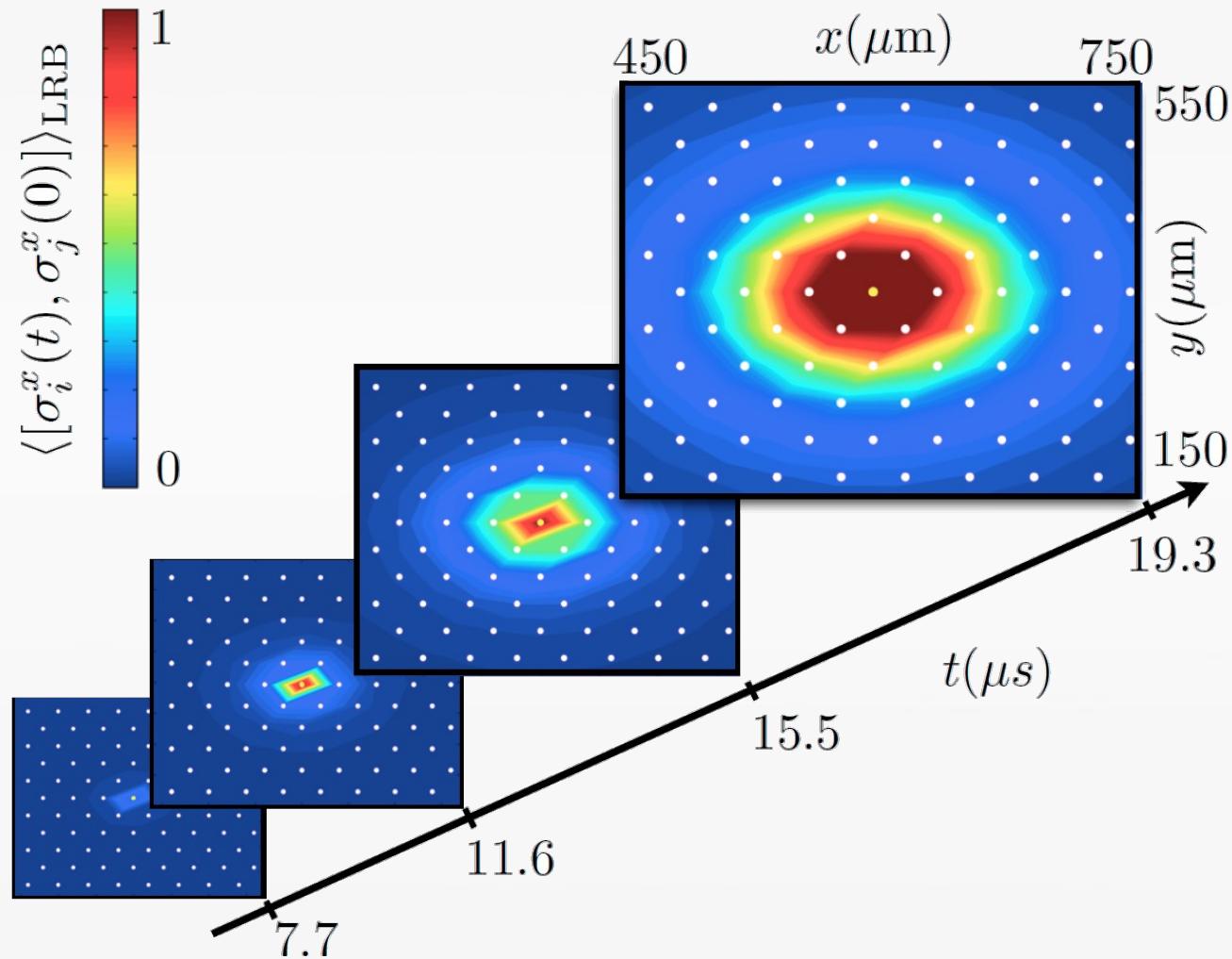


Impulsive forces



$$\langle B(t) \rangle \rightarrow \langle B(t) \rangle + \frac{\varepsilon}{\hbar} \langle [B(t), A(0)] \rangle + O(\varepsilon^2)$$

Correlation spread



Generalized LR bounds

$$\|[A_x(t), B_y(0)]\| \leq c \exp(\nu t - d_{x,y})$$

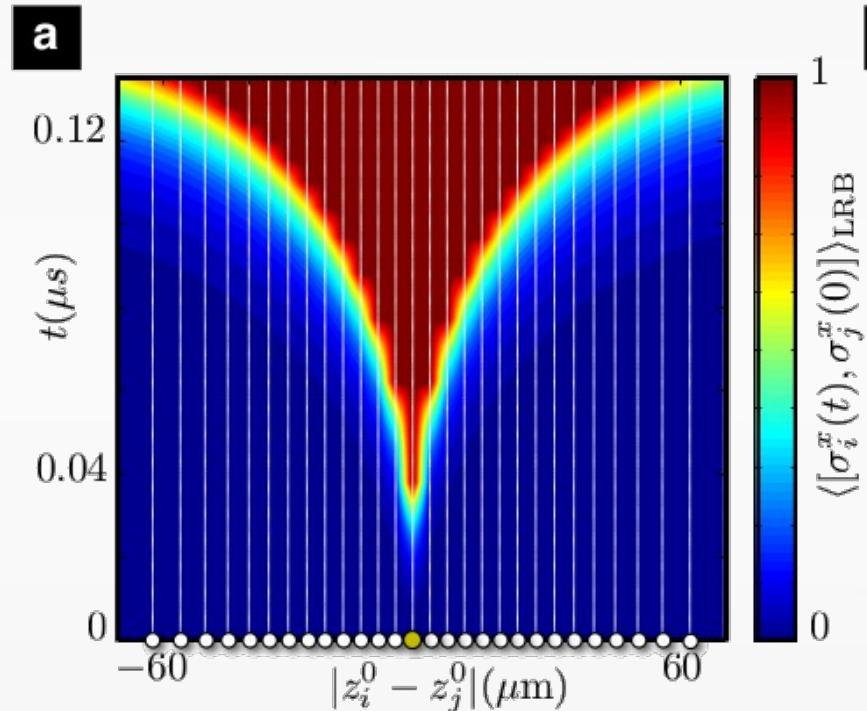
Outside an effective light-cone with speed ν correlations are exponentially attenuated.

$$\|[A_x(t), B_y(0)]\| \leq c \exp(\nu t) \frac{1}{d_{x,y}^\eta}$$

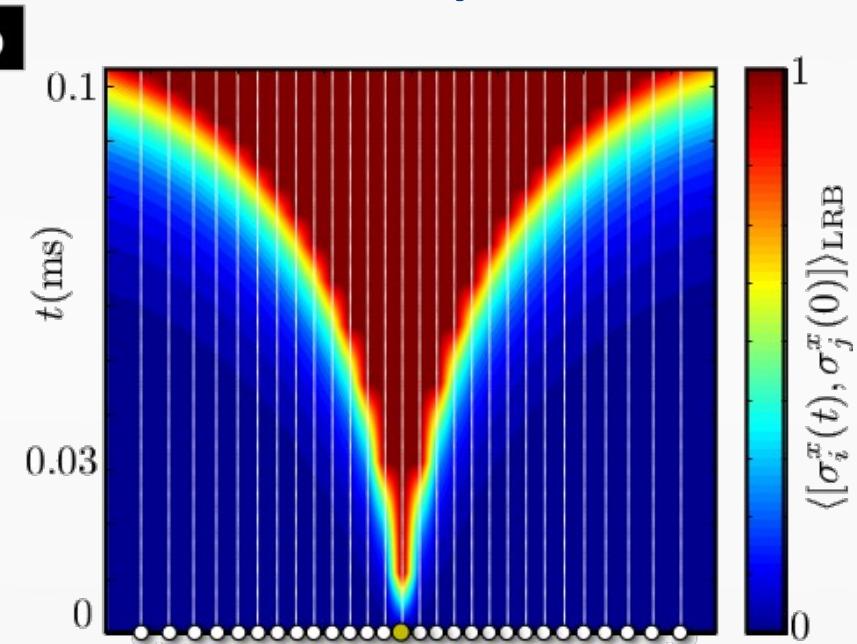
A milder version needed for long range interactions.

Correlation spread

Full model



Effective spins

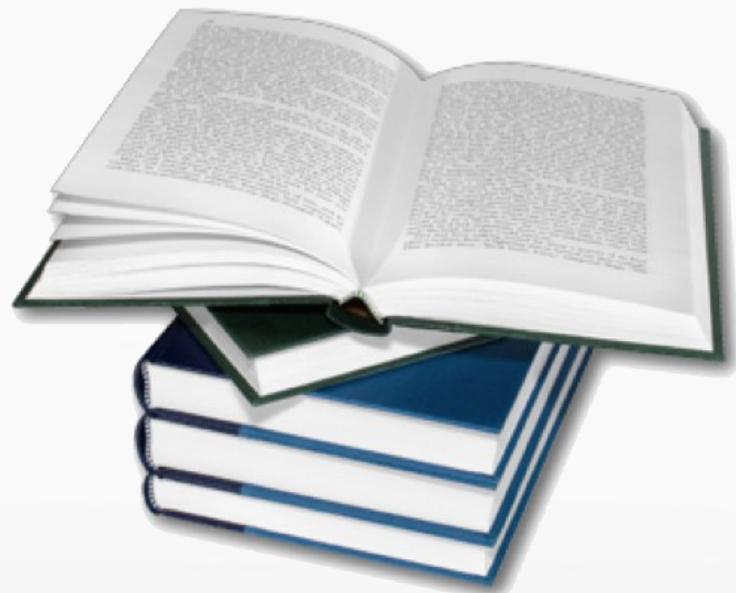


Approximate light cones because of long range interaction.

The full model is faster, because it involves the bosons in a non-perturbative fashion

Summary

- Spin-boson model of interest beyond quantum dissipation.
- New mathematical tools for analyzing correlation spread in these systems.
- Experimental possibility of measuring the propagation of correlations.



Acknowledgements

