

A certain necessary condition of possible blow up for Navier-Stokes equations

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Cauchy Problem

- ▶ Navier-Stokes equations

$$\begin{aligned}\partial_t v(x, t) + v(x, t) \cdot \nabla v(x, t) - \Delta v(x, t) + \nabla q(x, t) &= 0 \\ \operatorname{div} v(x, t) &= 0\end{aligned}$$

for $(x, t) \in \mathbb{R}^3 \times]0, \infty[$

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for $(x, t) \in \mathbb{R}^3 \times]0, \infty[$

- ▶ initial condition

$$v(\cdot, 0) = v_0(\cdot) \in C_{0,0}^\infty(\mathbb{R}^3)$$

Blowup Time

- ▶ T is a blowup time if

$$M(t) = \sup_{x \in \mathbb{R}^3} |v(x, t)| \rightarrow \infty$$

as $t \rightarrow T - 0$

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- ▶ **Question** is whether or not $T = \infty$

Leray's Necessary Conditions

$$\|v(\cdot, t)\|_\infty \geq \frac{C}{\sqrt{T-t}}$$

Other norms blow up:

$$\|v(\cdot, t)\|_p = \left(\int_{\mathbb{R}^3} |v(x, t)|^p dx \right)^{\frac{1}{p}} \geq \frac{C(p)}{(T-t)^{\frac{p-3}{2p}}}$$

$$\|\nabla v(\cdot, t)\|_2 \geq \frac{C}{(T-t)^{\frac{1}{4}}}$$

Limit Case $p = 3$

$$\exists t_k \rightarrow T : \|v(\cdot, t_k)\|_3 \rightarrow \infty$$

Question:

$$\|v(\cdot, t)\|_3 \rightarrow \infty \quad \text{as } t \rightarrow T?$$



if there exists $t_k \rightarrow T$ s.t. $\sup_k \|v(\cdot, t_k)\|_3 =: M < \infty$, then T is not a blow up time

Negative Issue

Assume that there exists a function f such that $f(\varepsilon) \rightarrow \infty$ as $\varepsilon \rightarrow 0$ and

$$\|v(\cdot, t)\|_3 \geq f(T - t)$$

Indeed, if

$$\|v(\cdot, 0)\|_3 \geq f(T),$$

then for $v^\lambda(y, s) = \lambda v(\lambda y, \lambda^2 s)$

$$\|v(\cdot, 0)\|_3 = \|v^\lambda(\cdot, 0)\|_3 \geq f(T/\lambda^2) \rightarrow \infty$$

as $\lambda \rightarrow \infty$

Outline

- ▶ I. Scaling with a parameter proportional to the distance (in time) to blowup

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- ▶ IV. Decay at spatial infinity by ε -regularity
- ▶ V. Backward uniqueness

1. Blowup L_3 -Profile

$$M := \sup_k \|v(\cdot, t_k)\|_3 < \infty$$

\Downarrow

$$\|v(\cdot, T)\|_3 < \infty$$

12. Scaling

$$u^{(k)}(y, s) = \lambda_k v(x, t), \quad p^{(k)}(y, s) = \lambda_k^2 q(x, t),$$

where $(y, s) \in \mathbb{R}^3 \times]-\lambda_k^{-2}T, 0[$, $x = \lambda_k y$, $t = T + \lambda_k^2 s$,

$$\lambda_k = \sqrt{\frac{T - t_k}{S}}$$

and a positive parameter $S < 10$ will be defined later;

$$\sup_{k \in \mathbb{N}} \|u^{(k)}(\cdot, -S)\|_3 = M < \infty$$

II1. Decomposition $u^{(k)} = v^{(k)} + w^{(k)}$

$$\partial_t w^{(k)} - \Delta w^{(k)} = -\nabla r^{(k)}, \quad \operatorname{div} w^{(k)} = 0 \quad \text{in } \mathbb{R}^3 \times]-S, 0[,$$

$$w^{(k)}(\cdot, -S) = u^{(k)}(\cdot, -S)$$

and

$$\sup_k \{ \|w^{(k)}\|_{L^5(\mathbb{R}^3 \times]-S, 0[} + \|w^{(k)}\|_{L^{3,\infty}(\mathbb{R}^3 \times]-S, 0[} \} \leq c(M) < \infty$$

II2. Decomposition $u^{(k)} = v^{(k)} + w^{(k)}$ II

$$\partial_t v^{(k)} + \operatorname{div}(v^{(k)} + w^{(k)}) \otimes (v^{(k)} + w^{(k)}) - \Delta v^{(k)} = -\nabla p^{(k)}$$

$$\operatorname{div} v^{(k)} = 0 \quad \text{in } \mathbb{R}^3 \times]-S, 0[$$

$$v^{(k)}(\cdot, -S) = 0$$

II.3. Pressure

$$p^{(k)}(x, t) = -\frac{1}{3}|u^{(k)}(x, t)|^2 + \frac{1}{4\pi} \int_{\mathbb{R}^3} K(x-y) : u^{(k)}(y, t) \otimes u^{(k)}(y, t) dy$$

II4. Pressure Decomposition

For $x_0 \in \mathbb{R}^3$ and for $x \in B(x_0, 3/2)$, we let

$$p_{x_0}^{(k)}(x, t) \equiv p^{(k)}(x, t) - c_{x_0}^{(k)}(t) = p_{x_0}^{1(k)}(x, t) + p_{x_0}^{2(k)}(x, t)$$

where

$$p_{x_0}^{1(k)}(x, t) = -\frac{1}{3}|u^{(k)}(x, t)|^2 + \frac{1}{4\pi} \int_{B(x_0, 2)} K(x-y) : u^{(k)}(y, t) \otimes u^{(k)}(y, t) dy$$

$$p_{x_0}^{2(k)}(x, t) = \frac{1}{4\pi} \int_{\mathbb{R}^3 \setminus B(x_0, 2)} (K(x-y) - K(x_0-y)) : u^{(k)}(y, t) \otimes u^{(k)}(y, t) dy$$

$$c_{x_0}^{(k)}(t) = \frac{1}{4\pi} \int_{\mathbb{R}^3 \setminus B(x_0, 2)} K(x_0 - y) : u^{(k)}(y, t) \otimes u^{(k)}(y, t) dy$$

II.5. Local Energy

$$\alpha(s) = \alpha(s; k, S) = \|v^{(k)}(\cdot, s)\|_{2, \text{unif}}^2$$

$$\beta(s) = \beta(s; k, S) = \sup_{x \in \mathbb{R}^3} \int_{-S}^s \int_{B(x, 1)} |\nabla v^{(k)}|^2 dy d\tau, \quad \text{where}$$

$$\|g\|_{L_2, \text{unif}} = \sup_{x_0 \in \mathbb{R}^3} \|g\|_{L_2(B(x_0, 1))}.$$

More quantities:

$$\gamma(s) = \gamma(s; k, S) = \sup_{x \in \mathbb{R}^3} \int_{-S}^s \int_{B(x, 1)} |v^{(k)}(y, \tau)|^3 dy d\tau$$

$$\delta(s) = \delta(s; k, S) = \sup_{x \in \mathbb{R}^3} \int_{-S}^s \int_{B(x, 3/2)} |p^{(k)}(y, \tau) - c_x^{(k)}(\tau)|^{\frac{3}{2}} dy d\tau$$

II6. Pressure estimates

$$\|p_{x_0}^{1(k)}(\cdot, t)\|_{L_{\frac{3}{2}}(B(x_0, 3/2))} \leq c(M)(\|v^{(k)}(\cdot, t)\|_{L_3(B(x_0, 2))}^2 + 1)$$

and

$$\sup_{B(x_0, 3/2)} |p_{x_0}^{2(k)}(x, t)| \leq c(M)(\|v^{(k)}(\cdot, t)\|_{L_{2, \text{unif}}}^2 + 1)$$

II7. Energy Estimate and Choice of S

$$\alpha(s) + \beta(s) \leq c(M) \left[(s + S)^{\frac{1}{5}} + \int_{-S}^s \left(\alpha(\tau) (1 + \|w^{(k)}(\cdot, \tau)\|_{L^5, \text{unif}}^5) + \alpha^3(\tau) \right) d\tau \right]$$

$\forall s \in [-S, 0[$ and $c(M)$ is independent of k , s , and S .

There is a positive constant $S(M)$:

$$\alpha(s) \leq \frac{1}{10}$$

and

$$\alpha(s) \leq c(M)(s + S)^{\frac{1}{5}}$$

for any $s \in]-S(M), 0[$

III.1. Limiting Procedure

$$w^{(k)} \rightarrow w$$

$$w(x, t) = \frac{1}{(4\pi(s + S))^{\frac{3}{2}}} \int_{\mathbb{R}^3} \exp\left(-\frac{|x - y|^2}{4(s + S)}\right) w_0(y) dy$$

where w_0 is the weak $L_3(\mathbb{R}^3)$ -limit of the sequence $u^{(k)}(\cdot, -S)$;

$$\sup_{-S < s < 0} \sup_{x_0 \in \mathbb{R}^3} \|w(\cdot, s)\|_{L_2(B(x_0, 1))}^2 +$$

$$+ \sup_{x_0 \in \mathbb{R}^3} \int_{-S}^0 \int_{B(x_0, 1)} |\nabla w(y, s)|^2 dy ds \leq c(M) < \infty$$

$$w \in C([-S, 0]; L_3(\mathbb{R}^3)) \cap L_5(\mathbb{R}^3 \times]-S, 0])$$

III.2. Limit function u is a local energy solution

$$u^{(k)} \rightarrow u, \quad p^{(k)} \rightarrow p$$

$$\sup_{-S < s < 0} \sup_{x_0 \in \mathbb{R}^3} \|u(\cdot, s)\|_{L_2(B(x_0, 1))}^2 + \sup_{x_0 \in \mathbb{R}^3} \int_{-S}^0 \int_{B(x_0, 1)} |\nabla u(y, s)|^2 dy ds < \infty;$$

$$p \in L_{\frac{3}{2}}(-S, 0; L_{\frac{3}{2}, \text{loc}}(\mathbb{R}^3));$$

the function

$$s \mapsto \int_{\mathbb{R}^3} u(y, s) \cdot w(y) dy$$

is continuous on $[-S, 0]$ for any compactly supported $w \in L_2(\mathbb{R}^3)$;

III.3. Limit function u is a local energy solution

$$\partial_t u + u \cdot \nabla u - \Delta u = -\nabla p, \quad \operatorname{div} u = 0$$

in $\mathbb{R}^3 \times]-S, 0[$ in the sense of distributions;

for any $x_0 \in \mathbb{R}^3$, there exists a function $c_{x_0} \in L^3_{\frac{3}{2}}(-S, 0)$ such that

$$p(x, s) - c_{x_0}(s) = p_{x_0}^1(x, s) + p_{x_0}^2(x, s)$$

for all $x \in B(x_0, 3/2)$, where

$$p_{x_0}^1(x, s) = -\frac{1}{3}|u(x, s)|^2 + \frac{1}{4\pi} \int_{B(x_0, 2)} K(x-y) : u(y, s) \otimes u(y, s) dy,$$

$$p_{x_0}^2(x, s) = \frac{1}{4\pi} \int_{\mathbb{R}^3 \setminus B(x_0, 2)} (K(x-y) - K(x_0-y)) : u(y, s) \otimes u(y, s) dy;$$

III4. Limit function u is a local energy solution

for any $s \in]-S, 0[$ and for $\varphi \in C_0^\infty(\mathbb{R}^3 \times]-S, S[)$,

$$\begin{aligned} & \int_{\mathbb{R}^3} \varphi^2(y, s) |u(y, s)|^2 dy + 2 \int_{-S}^s \int_{\mathbb{R}^3} \varphi^2 |\nabla u|^2 dy d\tau \leq \\ & \leq \int_{-S}^s \int_{\mathbb{R}^3} \left(|u|^2 (\Delta \varphi^2 + \partial \varphi^2) + u \cdot \nabla \varphi^2 (|u|^2 + 2p) \right) dy d\tau \end{aligned}$$

III5. Limit function u is a local energy solution

$$\sup_{x_0 \in \mathbb{R}^3} \|v(\cdot, s)\|_{L_2(B(x_0, 1))}^2 \leq c(M)(s + S)^{\frac{1}{5}}$$

for all $s \in [-S, 0]$.

\Downarrow

$$v(\cdot, s) \rightarrow 0 \quad \text{in } L_{2, \text{loc}}(\mathbb{R}^3)$$

as $s \downarrow -S$.

\Downarrow

$$u(\cdot, s) \rightarrow w_0 \quad \text{in } L_{2, \text{loc}}(\mathbb{R}^3).$$

as $s \downarrow -S$.

IV1 Limit function is not trivial

$$\frac{1}{a^2} \int_{Q(a)} (|u^{(k)}|^3 + |\tilde{p}^{(k)}|^{\frac{3}{2}}) dy ds = \frac{1}{(a\lambda_k)^2} \int_{Q(z_T, a\lambda_k)} (|v|^3 + |q - b^{(k)}|^{\frac{3}{2}}) dx dt$$

for all $0 < a < a_* = \inf\{1, \sqrt{S/10}, \sqrt{T/10}\}$ and for all $\lambda_k \leq 1$. Here, $z_T = (0, T)$, $\tilde{p}^{(k)} \equiv \tilde{p}_2^{(k)}$, and $b^{(k)}(t) = \lambda_k^{-2} c_2^{(k)}(s)$. Since the pair v and $q - b^{(k)}$ is a suitable weak solution to the Navier-Stokes equations in $Q(z_T, \lambda_k a_*)$, we find

$$\frac{1}{a^2} \int_{Q(a)} (|u^{(k)}|^3 + |\tilde{p}^{(k)}|^{\frac{3}{2}}) dy ds > \varepsilon \quad (1)$$

for all $0 < a < a_*$ with a positive universal constant ε

IV2 Limit function is not trivial

$$\frac{1}{a^2} \int_{Q(a)} |u^{(k)}|^3 dy ds \rightarrow \frac{1}{a^2} \int_{Q(a)} |u|^3 dy ds$$

for all $0 < a < a_*$ and

$$\sup_{k \in \mathbb{N}} \frac{1}{a_*^2} \int_{Q(a_*)} (|u^{(k)}|^3 + |\tilde{p}^{(k)}|^{\frac{3}{2}}) dy ds = M_1 < \infty$$

IV3 Limit function is not trivial

$$\tilde{p}^{(k)} = p_1^{(k)} + p_2^{(k)}$$

with

$$\|p_1^{(k)}(\cdot, s)\|_{\frac{3}{2}, B(a_*)} \leq c \|u^{(k)}(\cdot, s)\|_{3, B(a_*)}^2$$

$p_2^{(k)}$ is harmonic

$$\begin{aligned} \varepsilon &\leq \frac{1}{a^2} \int_{Q(a)} (|\tilde{p}^{(k)}|^{\frac{3}{2}} + |u^{(k)}|^3) dy ds \leq \\ &\leq c \frac{1}{a^2} \int_{Q(a_*)} |u^{(k)}|^3 dy ds + cM_1 a a_*^2 \end{aligned}$$

for all $0 < a < a_*/2$

IV4. Spatial Decay

$$w_0 \in L_3$$

\Downarrow

$$\|w_0\|_{2, B(x_0, 1)} \rightarrow 0$$

as $|x_0| \rightarrow \infty$

\Downarrow

There are positive numbers R , $T \in]a_*, S[$, and c_k with $k = 0, 1, \dots$ such that

$$|\nabla^k u(x, t)| \leq c_k$$

for any $x \in \mathbb{R}^3 \setminus B(R/2)$ and for any $t \in]-T, 0[$

V1. Backward Uniqueness

$$v(\cdot, T) \in L_3(\mathbb{R}^3) \Rightarrow u(\cdot, 0) = 0$$

\Downarrow

$$\omega = \nabla \wedge u = 0$$

at $t = 0$.

\Downarrow

Backward Uniqueness

$$\omega = 0$$

in $(\mathbb{R}^3 \setminus B(R/2)) \times]-T, 0[$

V2. Backward Uniqueness

technical part is to show

$$\omega = 0$$

in $\mathbb{R}^3 \times]-T, 0[$



$u(\cdot, t)$ is a bounded harmonic function in \mathbb{R}^3 with

$$\|u(\cdot, t)\|_{L_2(B(x_0, 1))} \rightarrow 0$$

as $|x_0| \rightarrow \infty$



$$u = 0$$

in $\mathbb{R}^3 \times]-T, 0[$