

On Slow Motion for Phase Transitions

Calculus of Variations, PDE, and Geometric Measure Theory

Giovanni Leoni

Carnegie Mellon University

September 7, 2015

Papers:

- [G. Dal Maso, I. Fonseca and G.L.](#), 2015, Calc. Var. Partial Differential Equations,

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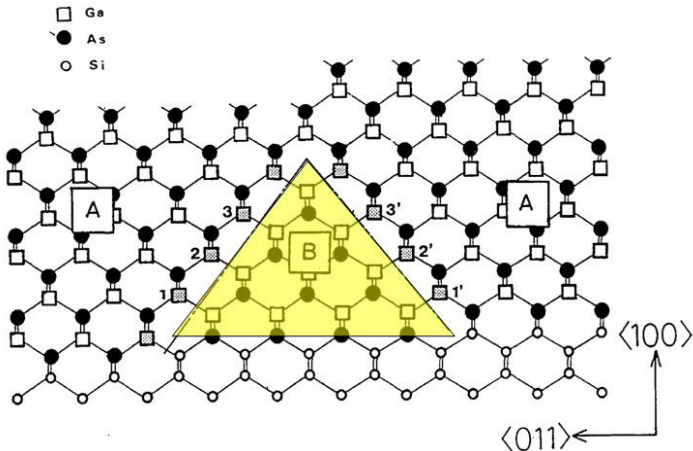
- G. Dal Maso, I. Fonseca and G.L., 2015, Calc. Var. Partial Differential Equations,
- G.L. and R. Murray, 2015, ARMA,
- R. Murray and M. Rinaldi, in preparation,
- I. Fonseca, G.L., G. Hayrapetyan, and M. Rinaldi, in preparation.

Allen and Cahn (1979)

- Motion of a curved antiphase boundary (APB)

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- Motion of a curved **antiphase boundary (APB)**
- APB: Planar defect in polycrystalline materials



- Interfacial motion theory: Smoluchowski (1951), Turnbull (1951)

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- V interfacial velocity,
- μ mobility,
- σ surface free energy per unity area,
- κ mean curvature of the interface.

Allen and Cahn (1979)

$$\begin{cases} u_t = \varepsilon^2 \Delta u - f(u) & \text{in } \Omega \times (0, \infty), \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial\Omega \times (0, \infty), \\ u(x, 0) = u_0(x) & \text{in } \Omega, \end{cases} \quad (\text{ACE})$$

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- $\varepsilon > 0$,
- $f(s) = W'(s)$, W double-well potential.

Prototype

$$W(s) := \frac{1}{4}(s^2 - 1)^2.$$

Allen–Cahn Equation

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$L^2(\Omega)$ gradient flow

$$F_\varepsilon(u) = \int_\Omega \left[\frac{1}{4}(u^2 - 1)^2 + \varepsilon^2 |\nabla u|^2 \right] dx.$$

- -1 and 1 two phases of the crystal,

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- $\varepsilon^2 |\nabla u|^2$ penalizes interfaces.

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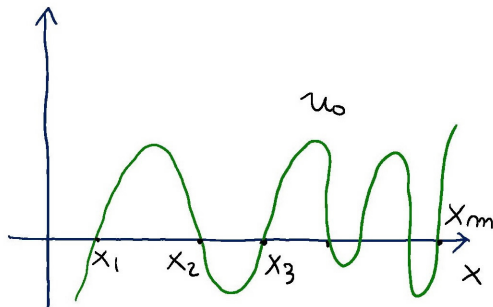
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- Mugnai and Röger (2008): Gamma convergence by gradient flows Sandier and Serfaty (2004).

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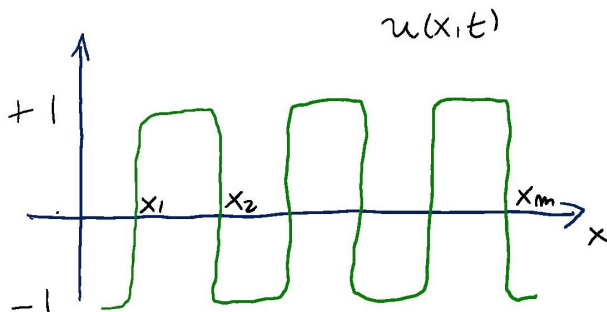
- u_0 smooth with m simple zeros



Phase I: Formation of Interfaces

$$u_t = \varepsilon^2 u_{xx} - u^3 + u$$

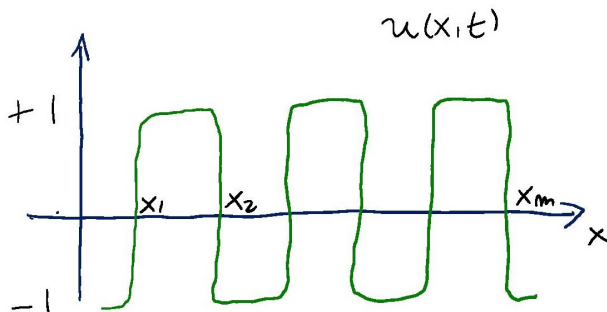
- $\varepsilon \ll 1$ for small times $t \ll 1$, $\varepsilon^2 u_{xx}$ negligible



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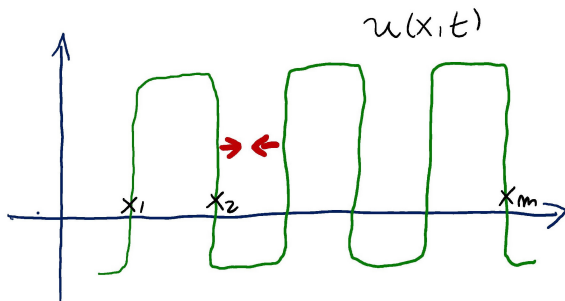
- $\varepsilon \ll 1$ for small times $t \ll 1$, $\varepsilon^2 u_{xx}$ negligible
- $u_t \simeq -u^3 + u$



Phase II: Slow Motion

$$u_t = \varepsilon^2 u_{xx} - u^3 + u$$

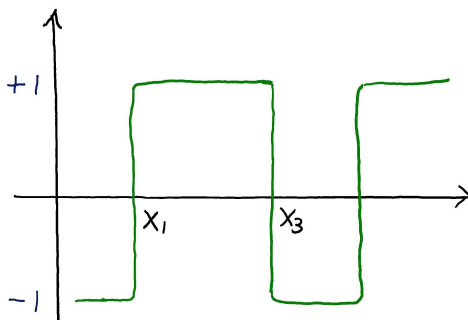
- $\varepsilon^2 u_{xx}$ no longer negligible



Phase III: Disappearance of Interfaces

$$u_t = \varepsilon^2 u_{xx} - u^3 + u$$

- when interfaces are close enough to a or b or to each other, they disappear quickly



Slow Motion: The Dynamical Approach

$$\begin{cases} u_t = \varepsilon^2 u_{xx} - u^3 + u & (x, t) \in (a, b) \times (0, \infty), \\ u_x(a, t) = u_x(b, t) = 0 & t \in (0, \infty), \\ u(x, 0) = u_0(x) & x \in (a, b), \end{cases}$$

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- Carr and Pego (1989), speed of order $ce^{-c/\varepsilon}$:

$$\frac{dx_n(t)}{dt} = \frac{64}{\varepsilon} \left[e^{-\sqrt{2}(x_n(t) - x_{n-1}(t))/\varepsilon} - e^{-\sqrt{2}(x_{n+1}(t) - x_n(t))/\varepsilon} \right]$$

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- Chen (1992), $N \geq 1$.

Slow Motion: The Energy Approach

$$G_\varepsilon(u) = \int_a^b \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |u'|^2 \right] dx.$$

- If $u_\varepsilon \in H^1(a, b)$ and $u_\varepsilon \rightarrow u_0$ in $L^1(a, b)$,
 $u_0 \in BV((a, b); \{-1, 1\})$,

$$\liminf_{\varepsilon \rightarrow 0^+} G_\varepsilon(u_\varepsilon) \geq c_0 m$$

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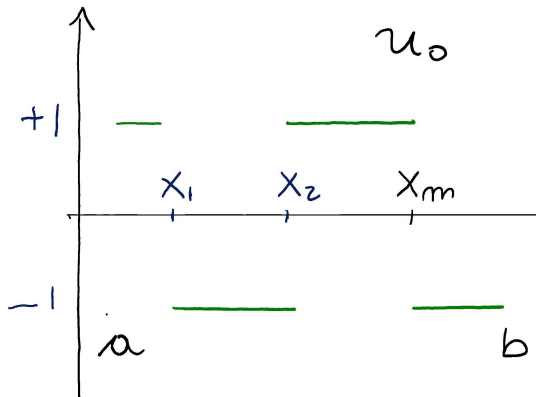
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Sternberg (1988) for $N \geq 1$,
- Modica and Mortola (1979) $W(s) = \sin^2(\pi s)$,

Previous Results: The 1-Dimensional Case

Slow Motion: The Energy Approach

If $u_\varepsilon \in H^1(a, b)$ and $u_\varepsilon \rightarrow u_0$ in $L^1(a, b)$,

$$\liminf_{\varepsilon \rightarrow 0^+} G_\varepsilon(u_\varepsilon) \geq c_0 m$$



Slow Motion: The Energy Approach

Improved Energy Estimate, I

$$G_\varepsilon(u) = \int_a^b \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |u'|^2 \right] dx.$$

Theorem (Bronsard and Kohn (1990))

For every $k \in \mathbb{N}$ there exist $\delta_k > 0$, $C_k > 0$ such that if $u \in H^1(a, b)$ and $\|u - u_0\|_{L^1} \leq \delta_k$,

$$G_\varepsilon(u) \geq c_0 m - C_k \varepsilon^k.$$

Slow Motion: The Energy Approach

Algebraic Slow Motion

$$\begin{cases} u_t = \varepsilon^2 u_{xx} - u^3 + u & (x, t) \in (a, b) \times (0, \infty), \\ u_x(a, t) = u_x(b, t) = 0 & t \in (0, \infty), \\ u(x, 0) = u_{0,\varepsilon}(x) & x \in (a, b). \end{cases} \quad (\text{AC})$$

Theorem (Bronsard and Kohn (1990))

For every $k \in \mathbb{N}$, if $u_{0,\varepsilon} \in H^1(a, b)$, $u_{0,\varepsilon} \rightarrow u_0$ in $L^1(a, b)$ and

$$G_\varepsilon(u_{0,\varepsilon}) \leq c_0 m + \varepsilon^k,$$

then solutions u_ε of (AC) satisfy for every $n > 0$,

$$\lim_{\varepsilon \rightarrow 0^+} \sup_{0 \leq t \leq n\varepsilon^{-k}} \int_a^b |u_\varepsilon(x, t) - u_0(x)| dx = 0.$$

Slow Motion: The Energy Approach

Improved Energy Estimate, II

$$G_\varepsilon(u) = \int_a^b \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |u'|^2 \right] dx.$$

Theorem (Grant (1995))

For every $L > 0$ there exist $\delta > 0$, $C > 0$ such that if $u \in H^1(a, b)$ and $\|u - u_0\|_{L^1} \leq \delta$,

$$G_\varepsilon(u) \geq c_0 m - C e^{-L/\varepsilon}.$$

Slow Motion: The Energy Approach

Exponential Slow Motion

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Theorem (Grant (1995))

For every $L > 0$, if $u_{0,\varepsilon} \in H^1(a, b)$, $u_{0,\varepsilon} \rightarrow u_0$ in $L^1(a, b)$ and

$$G_\varepsilon(u_{0,\varepsilon}) \leq c_0 m + e^{-L/\varepsilon},$$

then solutions u_ε of (AC) satisfy

$$\lim_{\varepsilon \rightarrow 0^+} \sup_{0 \leq t \leq e^{-L/\varepsilon}} \int_a^b |u_\varepsilon(x, t) - u_0(x)| dx = 0.$$

Slow Motion: The Energy Approach

Sharp Energy Estimate, III

$$G_\varepsilon(u) = \int_{\mathbb{T}} \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |u'|^2 \right] dx.$$

Theorem (Bellettini, Noyam, Novaga (2013))

For every $\{u_\varepsilon\} \subset H^1(\mathbb{T})$ with $u_\varepsilon \rightarrow u_0$ in $L^1(\mathbb{T})$, there exist $d_{n,\varepsilon} \rightarrow x_{n+1} - x_n$, $n = 1, \dots, m$, such that

$$G_\varepsilon(u_\varepsilon) \geq c_0 m - 16\sqrt{2} \sum_{n=1}^m e^{-\sqrt{2}d_{n,\varepsilon}/\varepsilon} + o\left(e^{-3\sqrt{2}d_{n,\varepsilon}/(2\varepsilon)}\right),$$

- \mathbb{T} one-dimensional torus.

Slow Motion: The Energy Approach

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- \mathbb{T} one-dimensional torus.
- This inequality is sharp.

Slow Motion: The Energy Approach

The **gradient flow** of

$$G_0(x_1, \dots, x_m) := c_0 m - 16\sqrt{2} \sum_{n=1}^m e^{-\sqrt{2}(x_{n+1}-x_n)/\varepsilon}$$

is given by

$$\begin{aligned} \frac{dx_n(t)}{dt} &= \frac{\partial G_0}{\partial x_n}(x_1, \dots, x_m) \\ &= \frac{64}{\varepsilon} \left[e^{-\sqrt{2}(x_n(t)-x_{n-1}(t))/\varepsilon} - e^{-\sqrt{2}(x_{n+1}(t)-x_n(t))/\varepsilon} \right]. \end{aligned}$$

- recover the results of Carr and Pego (1989).

Slow Motion: The N -Dimensional Case

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Dynamical Approach:

- [Kowalczyk \(1997\)](#), dynamics of a straight interface on a special domain,

Energy Approach:

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Energy Approach:

- Bronsard and Kohn (1991), radial case.

Slow Motion: The Energy Approach

Van Der Waals– Cahn–Hilliard Theory for Phase Transitions

- Free energy

$$F_\varepsilon(u) = \int_\Omega \left[\frac{1}{4}(u^2 - 1)^2 + \varepsilon^2 |\nabla u|^2 \right] dx.$$

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- $\int_\Omega u \, dx = m$, where m **total mass** of the fluid,
- **Van Der Waals (1893), Cahn and Hilliard (1958), Gurtin (1987)**

Slow Motion: The Energy Approach

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$$\liminf_{\varepsilon \rightarrow 0^+} G_\varepsilon(u_\varepsilon) \geq c_0 \operatorname{Per}_\Omega(\{u = 1\}),$$

- $\operatorname{Per}_\Omega(\{u = 1\})$ surface area of $\partial\{u = 1\}$ in Ω ,

Slow Motion: The Energy Approach

Van Der Waals– Cahn–Hilliard Theory for Phase Transitions

$$G_\varepsilon(u) = \int_\Omega \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |\nabla u|^2 \right] dx.$$

- If $u_\varepsilon \in H^1(\Omega)$, $\int_\Omega u_\varepsilon dx = m$, and $u_\varepsilon \rightarrow u$ in $L^1(\Omega)$, $u \in BV(\Omega; \{-1, 1\})$,

$$\liminf_{\varepsilon \rightarrow 0^+} G_\varepsilon(u_\varepsilon) \geq c_0 \operatorname{Per}_\Omega(\{u = 1\}),$$

- $\operatorname{Per}_\Omega(\{u = 1\})$ surface area of $\partial\{u = 1\}$ in Ω ,
- Carr, Gurtin, and Slemrod (1984) for $N = 1$, Modica and Mortola (1979), Modica (1987), Sternberg (1988) for $N \geq 1$.

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Global minimizers u_0 of

$$G_0(u) := c_0 \operatorname{Per}_\Omega(\{u = 1\}), \quad u \in BV(\Omega; \{-1, 1\}), \quad \int_\Omega u \, dx = m,$$

are of the form

$$u_0 = -1\chi_{\Omega \setminus E_0} + 1\chi_{E_0},$$

where $E_0 \subset \Omega$ minimizer of

$$\min \left\{ \operatorname{Per}_\Omega(E) : E \subset \Omega, E \text{ measurable, } \operatorname{meas}(E) = \frac{\operatorname{meas}(\Omega) - m}{2} \right\}.$$

- E_0 is regular for $N \leq 7$, has constant mean curvature κ ,

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- Gonzalez, Massari and Tamanini (1983), Grüter (1987)

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Sharp Energy Estimate

$$G_\varepsilon(u) = \int_\Omega \left[\frac{1}{4\varepsilon} (u^2 - 1)^2 + \varepsilon |\nabla u|^2 \right] dx.$$

Theorem (G.L. and Murray)

Let $\Omega \subset \mathbb{R}^N$ be open, bounded, of class $C^{2,\alpha}$, $\alpha > 0$. Then for \mathcal{L}^1 a.e. mass m and for every $\{u_\varepsilon\} \subset H^1(\Omega)$ with $\int_\Omega u_\varepsilon dx = m$ and $u_\varepsilon \rightarrow u_0$ in $L^1(\Omega)$,

$$G_\varepsilon(u_\varepsilon) \geq c_0 \operatorname{Per}_\Omega(E_0) - \frac{(N-1)^2}{9} \kappa^2 \varepsilon + o(\varepsilon).$$

- This inequality is sharp.

Slow Motion: The Energy Approach

$$G_\varepsilon(u) = \int_\Omega \left[\frac{1}{2\beta\varepsilon} |u^2 - 1|^\beta + \varepsilon |\nabla u|^2 \right] dx, \quad 1 < \beta < 2.$$

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- $E_0 \Rightarrow$ Ball compactly contained in Ω .

Why for a.e. Mass?

- Isoperimetric function

$$I_{\Omega}(s) := \min \{ \text{Per}_{\Omega}(E) : E \subset \Omega \text{ measurable,} \\ \text{meas}(E) = s \}.$$

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- [Sternberg and Zumbrun \(1998, 1999\)](#)

Slow Motion: Non Local Allen–Cahn Equation

Murray and Rinaldi

$$\begin{cases} u_t = \varepsilon^2 \Delta u - u^3 + u + \lambda & \text{in } \Omega \times (0, \infty), \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial\Omega \times (0, \infty), \\ u(x, 0) = u_{0,\varepsilon}(x) & \text{in } \Omega, \end{cases}$$

- $\lambda(t) = \frac{1}{\text{meas}(\Omega)} \int_{\Omega} (u^3(x, t) - u(x, t)) \, dx$ Lagrange multiplier accounting for $\int_{\Omega} u(x, t) \, dx = m$,

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- proposed by Rubinstein and Sternberg (1992).

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$$G_{\varepsilon}(u_{0,\varepsilon}) \leq c_0 \text{Per}_{\Omega}(E_0) + \varepsilon,$$

then for every $M > 0$ solutions u_{ε} of (NAC) satisfy

$$\lim_{\varepsilon \rightarrow 0^+} \sup_{0 \leq t \leq M/\varepsilon} \int_{\Omega} |u_{\varepsilon}(x, t) - u_0(x)| dx = 0.$$

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- $E_0 = \text{Ball} \Subset \Omega$ Done

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- $E_0 = \text{Ball} \Subset \Omega$ **Done**
- E_0 **isolated local minimizer** **In preparation.**

Slow Motion: Swift–Hohenberg Equation

Fonseca, G.L., Harapetyan and Rinaldi

$$u_t + \varepsilon^2 \gamma \Delta^2 u - \Delta u + \frac{1}{\varepsilon^2} (u - u^3) = 0.$$

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 - [Cicalese, Spadaro and Zeppieri \(2011\)](#), [Chermisi, Dal Maso, Fonseca, and G.L. \(2011\)](#), [Fonseca, Hayrapetyan, G. L. and Zwicknagl \(2014\)](#).

Gamma-Asymptotic Developments

Definition

X metric space, $F_\varepsilon^{(0)} : X \rightarrow (-\infty, \infty]$ has a **Γ -asymptotic development of order k** ,

$$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \dots + \varepsilon^k F^{(k)} + o(\varepsilon^k),$$

if

- $F_\varepsilon^{(0)} \xrightarrow{\Gamma} F^{(0)}$

- Anzellotti and Baldo (1993)

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- $F_\varepsilon^{(i)} := \frac{F_\varepsilon^{(i-1)} - \inf_{x \in X} F^{(i-1)}}{\varepsilon} \xrightarrow{\Gamma} F^{(i)}$ for $i = 1, \dots, k$.

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Gamma-Asymptotic Developments: First Example

Anzellotti and Baldo (1993)

$$F_\varepsilon^{(0)}(u) := \int_{\Omega} (W(u) + \varepsilon^2 |\nabla u|^2) dx$$

if $u \in H^1(\Omega)$, $u = g$ on $\partial\Omega$, $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$.

- $\Omega \subset \mathbb{R}^N$ open, bounded, connected, C^2 boundary,

Theorem (Anzellotti & Baldo)

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$$F^{(1)}(u) := 2c_0 \text{Per}_\Omega(u^{-1}([c, d])) + 2 \int_{\partial\Omega} |\Phi(u) - \Phi(g)| d\mathcal{H}^{N-1}$$

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- $\Phi(t) := \int_b^t \sqrt{W(s)} ds$, $c_0 := \Phi(c) = \int_b^c \sqrt{W(s)} ds$

2-Order Gamma-Asymptotic Development

- $N = 1$, $\Omega = (0, 1)$, $g(0) := u_0 \in (a, b)$,
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2-Order Gamma-Asymptotic Development

- $N = 1$, $\Omega = (0, 1)$, $g(0) := u_0 \in (a, b)$,
 $g(1) := u_1 \in (c, d)$,

$$F_\varepsilon^{(0)}(u) := \int_0^1 (W(u) + \varepsilon^2 |u''|^2) dx$$

if $u \in H^1(\Omega)$, $u = g$ on $\partial\Omega$, $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$.

Theorem (Anzellotti & Baldo)

$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \varepsilon^2 F^{(2)} + o(\varepsilon^2)$, where

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- and $F^{(2)}(u) := \infty$ otherwise in $L^1(\Omega)$.

Gamma-Asymptotic Developments: Second Example

Anzellotti, Baldo, and Orlandi (1996)

$$F_\varepsilon^{(0)}(u) := \int_{\Omega} (u^2 + \varepsilon^2 |\nabla u|^2) dx$$

if $u \in H^1(\Omega)$ and $u = g$ on $\partial\Omega$, and $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\mathbb{R}^N)$.

- $\Omega \subset \mathbb{R}^N$ open, bounded, connected, C^3 boundary,

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- $\Omega \subset \mathbb{R}^N$ open, bounded, connected, C^3 boundary,
- $g \in C^2(\partial\Omega)$, $g > 0$

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if $u \in H^1(\Omega)$ and $u = g$ on $\partial\Omega$, and $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$.

Theorem (Anzellotti, Baldo, & Orlandi)

$$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \varepsilon^2 F^{(2)} + o(\varepsilon^2),$$

where

- $F^{(0)}(u) := \int_{\Omega} u^2 dx$ if $u \in L^2(\Omega)$ and $F^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$,

$$F_\varepsilon^{(0)}(u) := \int_{\Omega} (u^2 + \varepsilon^2 |\nabla u|^2) dx$$

if $u \in H^1(\Omega)$ and $u = g$ on $\partial\Omega$, and $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$.

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Theorem (Anzellotti, Baldo, & Orlandi)

$$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \varepsilon^2 F^{(2)} + o(\varepsilon^2),$$

where

- $F^{(2)}(u) := -\frac{1}{2} \int_{\partial\Omega} g^2 K_1 d\mathcal{H}^{N-1}$ if $u = 0$ a.e. in Ω and $F^{(2)}(u) := \infty$ otherwise in $L^1(\Omega)$,

$$F_\varepsilon^{(0)}(u) := \int_{\Omega} (u^2 + \varepsilon^2 |\nabla u|^2) \, dx$$

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- $F^{(2)}(u) := -\frac{1}{2} \int_{\partial\Omega} g^2 K_1 \, d\mathcal{H}^{N-1}$ if $u = 0$ a.e. in Ω and $F^{(2)}(u) := \infty$ otherwise in $L^1(\Omega)$,
- for $x \in \partial\Omega$, $K_1(x)$ is the symmetric curvature of order 1 at x .

Gamma-Asymptotic Developments: Cahn-Hilliard

$$F_\varepsilon^{(0)}(u) := \int_{\Omega} ((u^2 - 1)^2 + \varepsilon^2 |\nabla u|^2) dx$$

if $u \in H^1(\Omega)$, $\int_{\Omega} u dx = m$ and $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L^1(\Omega)$.

$$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \varepsilon^2 F^{(2)} + o(\varepsilon^2).$$

Anzellotti, Baldo, and Orlandi (1996)

Problem

Is $F^{(2)}$ a functional depending on the curvature of the jump surface?

$$F_\varepsilon^{(0)}(u) := \int_{\mathbb{R}^2} ((u^2 - 1)^2 + \varepsilon^2 |\nabla u|^2) dx$$

if $u \in H_{\text{loc}}^1(\mathbb{R}^2)$, $\int_{\mathbb{R}^2} (u(x) + 1) dx = 2$ and $F_\varepsilon^{(0)}(u) := \infty$ otherwise in $L_{\text{loc}}^1(\mathbb{R}^2)$.

Theorem (Dal Maso, Fonseca, Focardi, & G.L.)

$$F_\varepsilon^{(0)} \stackrel{\Gamma}{=} F^{(0)} + \varepsilon F^{(1)} + \varepsilon^2 F^{(2)} + o(\varepsilon^2),$$

where

$$F^{(2)} = 0.$$