# **Gravitational Lensing**

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## OUTLINE

#### • THE DARK UNIVERSE THROUGH GRAVITATIONAL LENSING:

- Dark Matter
  - Cross-sections
  - Neutrino masses and hierarchy
- Dark Energy
- Dark Gravity
  - Modified Newtonian Dynamics, MOND
  - Braneworlds, f(R) etc



### STANDARD COSMOLOGICAL MODEL

 Universe is isotropic and homogeneous, started with Big Bang, is almost flat, governed by Einstein gravity, contains CDM, baryons, photons (++), a Cosmological Constant, and went through a period of rapid inflation, which produced adiabatic, near-gaussian fluctuations.







# ACCELERATION OF THE UNIVERSE



Saul Perlmutter





 Cosmological Constant?

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 $\Omega_M$ 



# EXTRACTING INFORMATION FROM COSMOLOGY

- Global measurements (e.g. of geometry) are few
- Fluctuations are also high in information content
- CMB fluctuations allows robust inference
- Galaxies are easy to measure but their fluctuations may not reflect fluctuations in mass ('galaxy bias')
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### GRAVITATIONAL LENSING

- Coherent distortion of background images by gravity
- Shear, magnification, amplification



Jain & Seljak

- Independent of the *dynamical state* of matter
- Independent of the *nature* of matter
   ROBUST INFERENCE

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## EARLY LENSING



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### LIGHTS ALL ASKEW IN THE HEAVENS

Men of Science More or Less Agog Over Results of Eclipse Observations.

#### **EINSTEIN THEORY TRIUMPHS**

Stars Not Where They Seemed or Were Calculated to be, but Nobody Need Worry.

A BOOK FOR 12 WISE MEN

No More in All the World Could Comprehend It, Said Einstein When. His Daring Publishers Accepted It.

Liebe Mutter! Hente eine glrendige Nach-richt. H. A. Low ty hat min telegraphiert, dass die englischen Expeditionen die Lichtablenkung ander Jour winklich bewilsen habers. Maja selveibt mis leider, dass Du nicht un viel Ichunergere heist, roudere dass For Dir auch work truthe ige danken mache st. Whe ger wirde ich Dir wieder Gesellschaftlistery duss ber nicht dens hedsslichen grihele siberlassen warest! Aber in Weile worde is to doch her bliphen missen and arbedten. tuch mach Hollerud werde ich fits einige tage faturen, um mich chrenfest daukbar gre ereveisen abwohlder Teitverlinst recht sunnylich ast. Och winselve der von Horpen gute lage. Lei munig geguint von Seinen stilbert

## STRONG LENSING



### WEAK LENSING: THE GEORGE W. BUSH YEARS

- 2000 First detections (Bacon et al, Kaiser et al, Wittman et al, van Waerbeke et al)
- 2002+ Weak-lensing selected cluster catalogues (e.g. Miyazake et al, Wittman et al)
- 2003+ Non-parametric masses in clusters (Kneib et al, Clowe et al, Jee et al, Gray++)
- 2003+ Matter power spectrum (Brown et al, Heymans et al, Hoekstra et al, Semboloni++)
- 2004 Bullet cluster challenge to MOND (Clowe et al)
- 2004+ 3D potential reconstruction (Taylor et al, Massey et al)
- 2005+ Evolution of structure (Bacon et al)
- 2006+ 3D analyses (Heavens et al, Kitching et al, Taylor et al)
- 2007 100 sq deg surveys, with small error bars (Benjamin et al, Fu et al)
   2010 COSMOS & CFHTLenS + first surveys designed for lensing (Pan-STARRS 1)
   2019+ Euclid, LSST

### NO DARK MATTER? 'BULLET CLUSTER'

#### Challenges MOND, TeVeS



Hot Gas (X-ray)

Markevitch et al 2002; Clowe et al 2004

Dark Matter (Lensing) Caveat: lensing measures the convergence, K, which is proportional to surface density in GR, but not in general modified gravity models.



### SELF-INTERACTING DARK MATTER?

- Bullet cluster  $\rightarrow \sigma/m < 0.12 \text{ m}^2/\text{kg}$  (Randall et al 2007)
- Limit is about 12 orders of magnitude larger that that of interest for particle physics for typical masses



### WEAK GRAVITATIONAL LENSING

• For small scalar perturbations

$$ds^{2} = \left(1 + \frac{2\Phi}{c^{2}}\right)c^{2}dt^{2} - \left(1 - \frac{2\Psi}{c^{2}}\right)R^{2}(t)\left[dr^{2} + S_{k}^{2}(r)d\psi^{2}\right]$$

• In terms of conformal time  $[d\eta = dt/R(t)]$ 

$$\frac{d^2 \mathbf{x}}{d\eta^2} = -\frac{1}{c^2} \nabla (\Phi + \Psi) \quad [\mathbf{x} = r(\theta_x, \theta_y)]$$

If no anisotropic stress, and GR,  $\Phi = \Psi$ 

$$\frac{d^2 \mathbf{x}}{d\eta^2} = -\frac{2}{c^2} \nabla \Phi$$

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A2218 HST

### AMPLIFICATION, MAGNIFICATION & SHEAR



$$\frac{\partial \beta_i}{\partial \theta_j} = \delta_{ij} - \phi_{,ij} = \begin{pmatrix} 1-\kappa & 0\\ 0 & 1-\kappa \end{pmatrix} + \begin{pmatrix} \gamma_1 & -\gamma_2\\ -\gamma_1 & -\gamma_1 \end{pmatrix}$$

Cosmological lensing potential (GR assumed):

$$\phi(\mathbf{r}) = \frac{2}{c^2} \int_0^r dr' \frac{S_k(r-r')}{S_k(r)S_k(r')} \Phi(\mathbf{r}')$$



Shear, magnification, amplification



Note: dependence is on gravitational potential: lensing probes the mass distribution directly. Galaxy bias is not an issue.

Matter distribution need not be relaxed.



### **3D RECONSTRUCTION: COSMOS FIELD**

COSMOS data (Massey et al 2007)



Beware! poor resolution in z (200 Mpc)



NEUROBIOLOGY Robots that think they're insects PANDEMIC FLU Why the 1918 outbreak was so deadly MOLECULAR MAGNETS An attractive proposition

HEUNSEEN UNIVERSE Dark matter maps reveal cosmic scaffolding

ATUREJOBS leating retirement



### SENSITIVITY TO COSMOLOGY

- Cosmic Shear statistical properties depend on
- a) how clumpy the Universe is, and its *growth rate*, i.e. P(k;t)
- b) the source distances, hence the *distance-redshift relation*, r(z)
- c) The gravity law (e.g. modified Poisson equations)

$$\phi(\vec{r}) = \frac{1}{c^2} \int_0^r \frac{S_k(r-r')}{S_k(r)S_k(r')} \left[\Phi(\vec{r}) + \Psi(\vec{r})\right]$$

$$\phi(\mathbf{r}) = \frac{2}{c^2} \int_0^r dr' \frac{S_k(r-r')}{S_k(r)S_k(r')} \Phi(\mathbf{r}')$$



Euclid proposal

### DARK ENERGY

- Measurable Effects of Dark Energy:
  - Distance-redshift relation

 $p_q = w(a) \ \boldsymbol{\rho}_q c^2$ a(t) = R(t)/R(now)

$$r = \int_0^z dz' \frac{c}{H(z')} \qquad \qquad H(t) = \frac{1}{R} \frac{dR}{dt}$$

where the Hubble parameter is given by

$$H^{2}(a) = H_{0}^{2} \left[ \Omega_{m} a^{-3} + \Omega_{k} a^{-2} + \Omega_{q} \exp\left(3 \int_{1}^{a} \frac{da'}{a'} \left[1 + w(a')\right]\right)^{-1} \right]$$

• Growth rate of perturbations (via H(a))

Assuming DE is smooth,

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$$\ddot{\delta} + 2H\dot{\delta} - 4\pi G\rho_m \delta = 0$$

Assumes GR.  $\delta$  = fractional mass overdensity



### RECENT RESULTS: CFHTLS AND COSMOS

#### New CFHTLS results soon



### DARK ENERGY PROPERTIES: |S W = -|?

#### • CFHTLS: -1.18 < w < -0.88 (95%) $[p=w\rho c^2]$



### DARK MATTER: NEUTRINOS

- Neutrino oscillation experiments give: -
- Since they interact only weakly/ gravitationally, they can exit from small perturbations and partially erase them.
- Free-streaming length is 37 (m\_v/eV)<sup>-1</sup> Mpc; affects the matter power spectrum most sensitive to  $\Sigma$  m<sub>v</sub>
- In principle, one can measure individual masses (de Bernardis et al 2009, but see Jimenez et al 2010)
- Expected error on sum with Planck and Euclid weak lensing: 0.037eV
- (see also Kitching et al 2008, Hannestad et al 2006; also lensing of CMB)

$$\Delta m_{12}^2 = 7.9^{+1.0}_{-0.8} \times 10^{-5} \text{eV}^2$$
$$|\Delta m_{23}^2| = 2.2^{+1.1}_{-0.8} \times 10^{-3} \text{eV}^2$$





### TOMOGRAPHY

- Use knowledge of source redshifts - better
- Bin galaxies according to their estimated redshift ('tomography')
- Powerful check on systematics, and probe for new physics
- COSMOS (Schrabback et al 2010) show expected scaling of lensing signal with redshift:



### SOME CHALLENGES

#### The Forward Process.

Galaxies: Intrinsic galaxy shapes to measured image:





Intrinsic galaxy (shape unknown) Gravitational lensing , causes a **shear (g)** 



Gravitational lensing Atmosphere and telescope causes a *shear (g)* cause a convolution



Detectors measure a pixelated image



Image also contains noise

#### Stars: Point sources to star images:



Intrinsic star (point source)







**Detectors** measure

a pixelated image



Image also contains noise



### DISTANCES ARE ROUGH

- Photometric redshifts (or 'photozs')
- Rough redshift from broad-band colours





### INTRINSIC ALIGNMENTS

G

Image Ellipticity = Source Ellipticity + Shear

 $e = \gamma + e_s$ 

- Dispersion in  $e_s$  is ~0.3; shear is ~0.02
- Two-point statistics (I=foreground 2=background):  $\langle e_1 e_2^* \rangle = \langle \gamma_1 \gamma_2^* \rangle + \langle \gamma_1 e_{s2}^* \rangle + \langle e_{s1} \gamma_2^* \rangle + \langle e_{s1} e_{s2}^* \rangle$ IG
- Previously: II=GI=IG=0.
- Tidal torques (e.g. Heavens et al 2000, Croft & Metzler 2000,...)  $\parallel \neq 0$ . Easily removed by downweighting close pairs
- $GI \neq 0$  (Hirata & Seljak 2004) term is more problematic.
- IG=0? Exercise!









### FUTURE EXPERIMENTS

- Euclid (ESA)
  - Selected for Cosmic Vision ~2019
  - Imaging + spectroscopy
  - 15,000 sq deg, median z=0.9, optical+IR
  - Ideal for Cosmic Shear, also BAOs
  - First space-based experiment designed for lensing







### PROSPECTS FOR DARK ENERGY



### MODIFIED GRAVITY

• It is already necessary to modify Einstein's original equations:

$$G^{\mu\nu} = \frac{8\pi G}{c^4} T^{\mu\nu}$$

 E.g. by including Dark Energy on right hand side, or including Λ on I.h.s. (equivalent)

$$G^{\mu\nu} - \Lambda g^{\mu\nu} = \frac{8\pi G}{c^4} T^{\mu\nu}$$

- Modified gravity: more general modifications to I.h.s.
- e.g. Einstein-Hilbert action changed:

$$R - 2\Lambda \to f(R)$$

### MODIFIED GRAVITY OR DARK ENERGY?

- GEOMETRIC MEASUREMENTS CANNOT TELL.
- Gravity law will give some H(t).



• Friedmann equation plus conservation of energy  $\Rightarrow$  can model any H(t) with GR plus Dark Energy with an equation of state parameter

$$w(t) = -\frac{1}{3} \frac{d}{d \ln R} \ln \left[ \frac{1}{\Omega_m(t)} - 1 \right]$$

• (Exercise. Flatmess assumed)

### BRANEWORLD MODELS

- Extra dimensions... broadly motivated by string theory
- E.g. DGP (Dvali, Gabadadze, Porrati 2000) model

$$S = \frac{1}{r_c} \int_{bulk} d^5 x^{\mu} R^{(5)} + \int_{brane} d^4 x^{\mu} R^{(4)}$$

Leads to a modified Friedmann equation

 $k^2 \Phi_{\mathbf{k}} = -4\pi G a^2 \left(1 + \frac{1}{3\beta}\right) \rho_m \delta_{\mathbf{k}}$ 

 $k^2 \Psi_{\mathbf{k}} = -4\pi G a^2 \left(1 - \frac{1}{3\beta}\right) \rho_m \delta_{\mathbf{k}}$ 

$$H^2 - \frac{H}{r_c} = \frac{8\pi G\rho}{3}$$

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$$\beta = 1 - 2r_c H \left(1 + \frac{\dot{H}}{3H^2}\right)$$

 $r_{c}^{-1} = H_{0}(1 - \Omega_{m})$ 

Growth rate of Newtonian potential is altered

In this case,  $\Phi + \Psi$  obeys the normal Poisson equation



### SUMMARY OF MODIFIED GRAVITY EFFECTS

- Expansion history R(t), [or H(t)] is changed
- Distance-redshift relation r(z) is changed
- Growth rate of matter fluctuations is altered
- Curvature and Newtonian potentials may behave differently
- Different response of photons to density perturbations
- Are these measurable?



### PROSPECTS FOR DARK GRAVITY

Compare GR with Dark Energy with a modified gravity model with the same expansion history.

Compute Expected Bayesian evidence = Probability of model given the data:

 $\sim$  do the data require a modification to GR?

• Euclid + Planck + BAO + SNe:

Probability ratio =  $e^{63}$  (DGP/GR)



Fractional difference in potentials



WMAP +WL (now); Planck ; Planck+Euclid

### 3D LENSING

- Estimated distances for all sources
- Galaxy 'shape' field is a very noisy, radially-smoothed, 3D point process sample of the galaxy shear field. (Heavens 2003)
- Better statistics (Heavens et al 2006; Kitching et al 2007)
- Better control of systematic errors (e.g. Bridle and King 2007)
- Can avoid the highly nonlinear regime where baryon physics is uncertain





# 3D WEAK LENSING FROM CFHTLENS



## SIZE MAGNIFICATION

Bias in size measurement

Looks very promising for space experiments



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• Casaponsa, AFH et al (in preparation)

### CONCLUSIONS

Lensing can probe a variety of phenomena of fundamental interest, such as

- Neutrino masses and their hierarchy
- The properties of the dominant Dark Matter component
- The Dark Energy equation-of-state
- Evidence for modifications to Einstein gravity

CMB and 3D lensing are particularly promising probes, as the physics is well-understood, and they have high sensitivity

Euclid + Planck will test this paradigm very thoroughly

Euclid is a Good Bet!



